

Theory perspective

Disclaimer: my own theory perspective

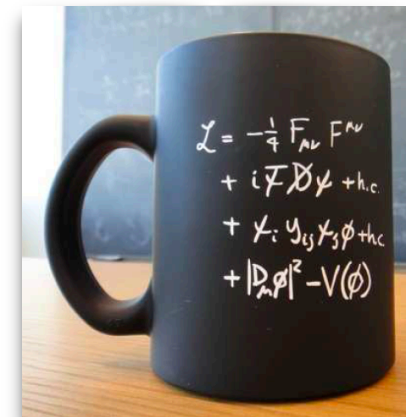
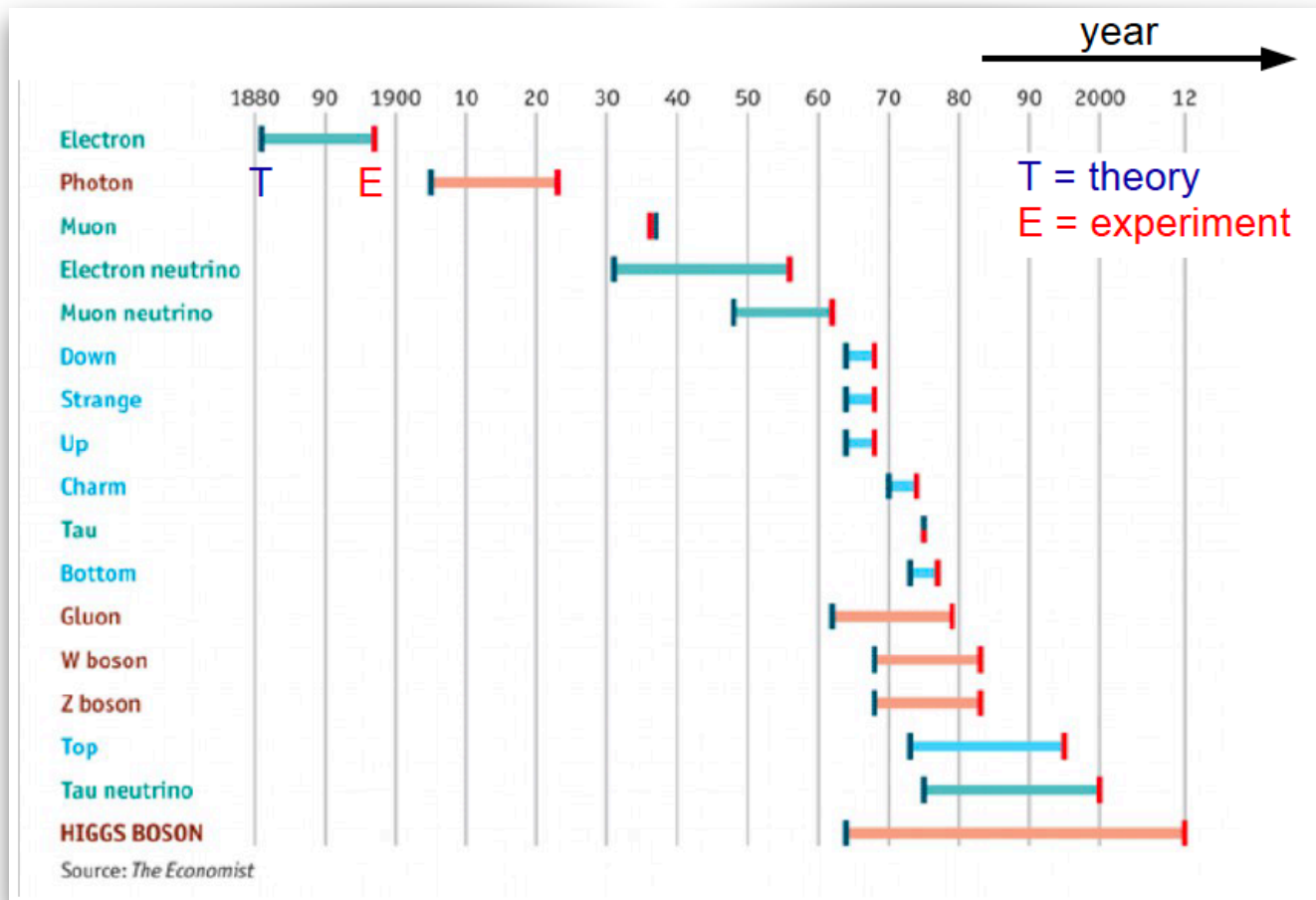
Stefania Gori
UC Santa Cruz



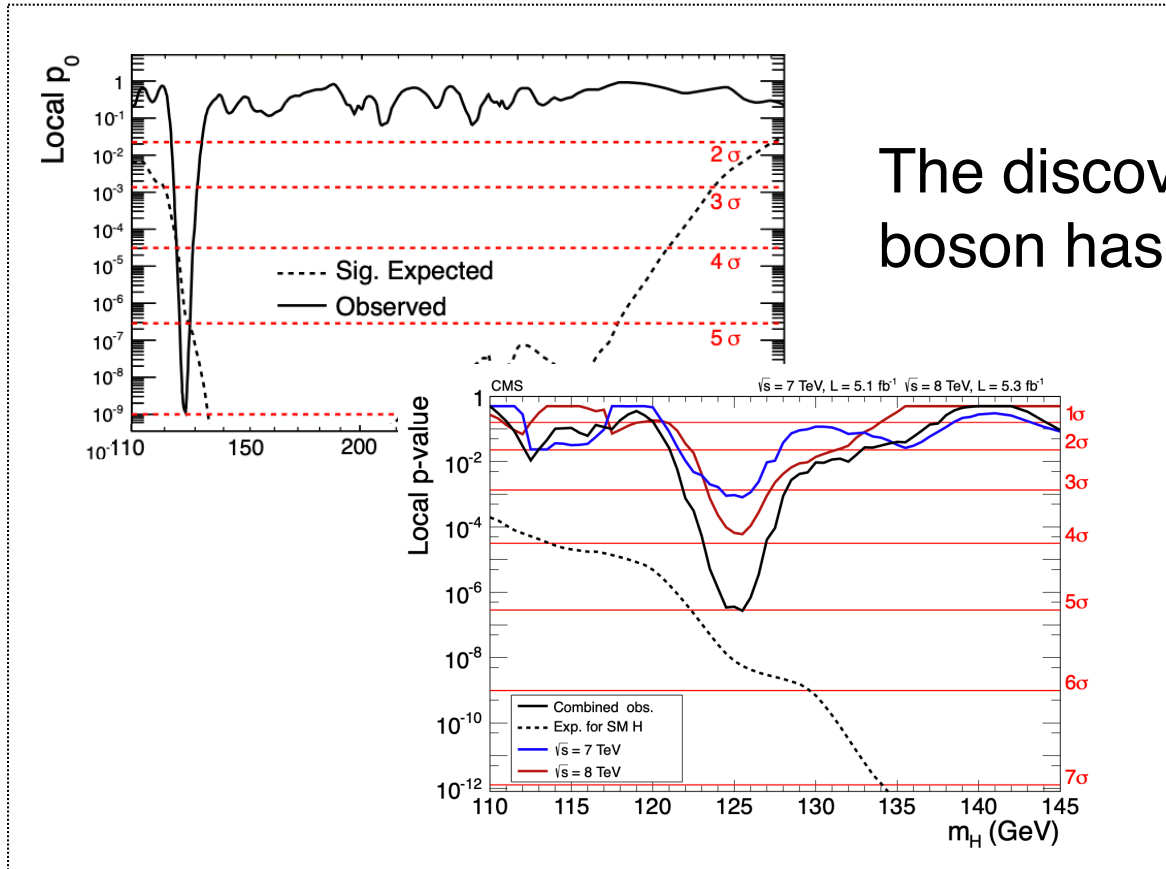
11th edition of the Large Hadron Collider Physics Conference (LHCp)
Belgrade,
May 26, 2023

Particle physics: past, present & future

Big success of HEP in the past 50+ years. HEP has enjoyed the remarkable achievement of uninterrupted fundamental discoveries!

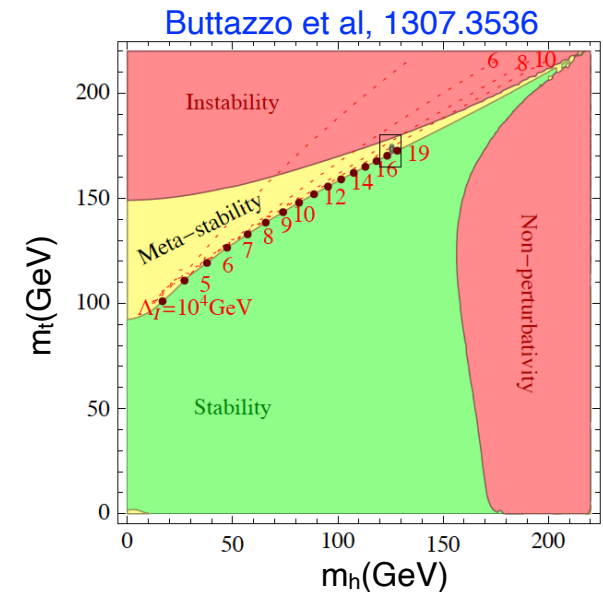


The LHC: a discovery machine



The discovery of the Higgs boson has been a turning point.

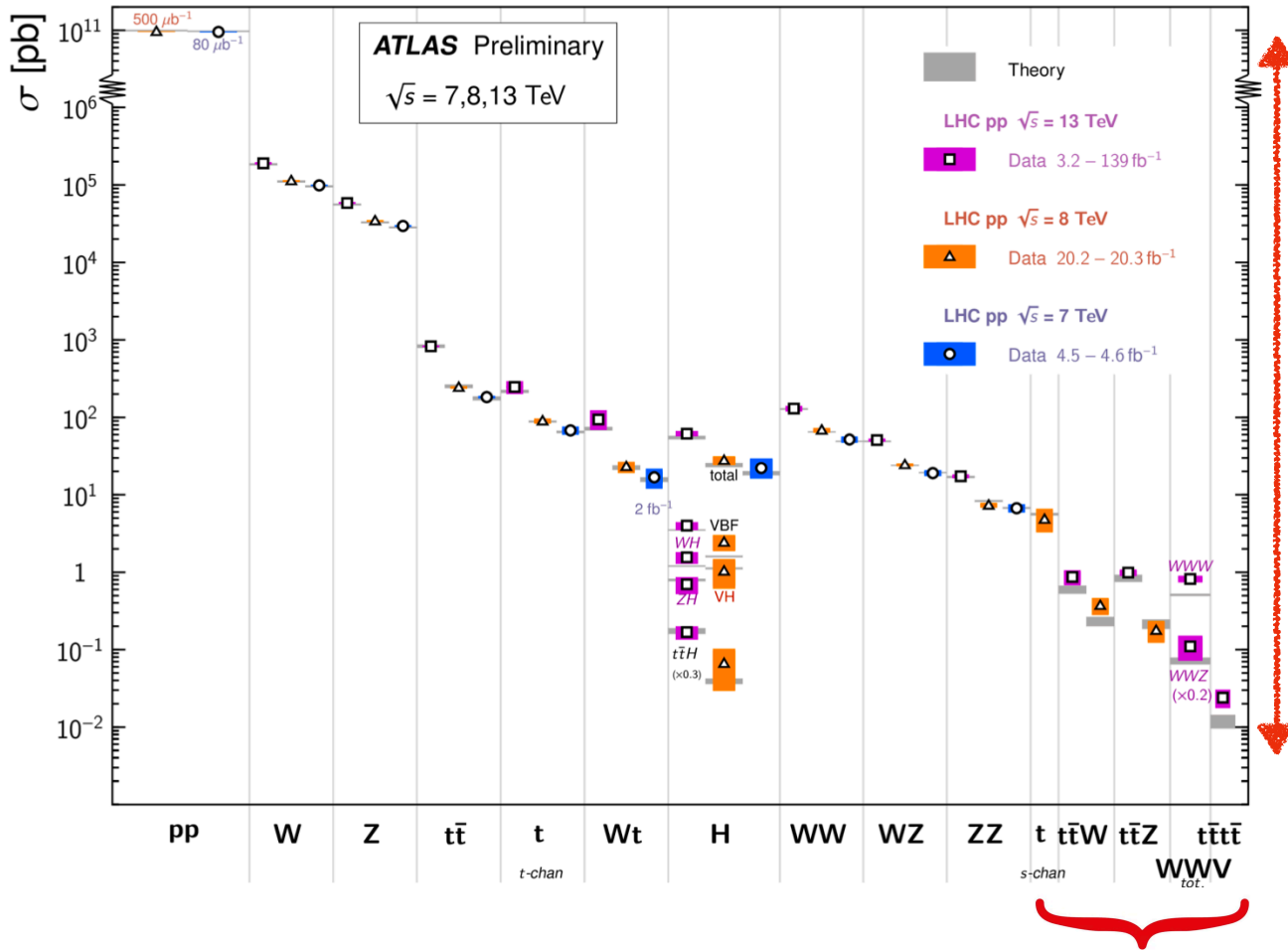
With the discovery of the Higgs at 125 GeV, for the first time in our history, **we have a self-consistent theory that can be extrapolated to exponentially higher energies.**



The LHC: a precision machine

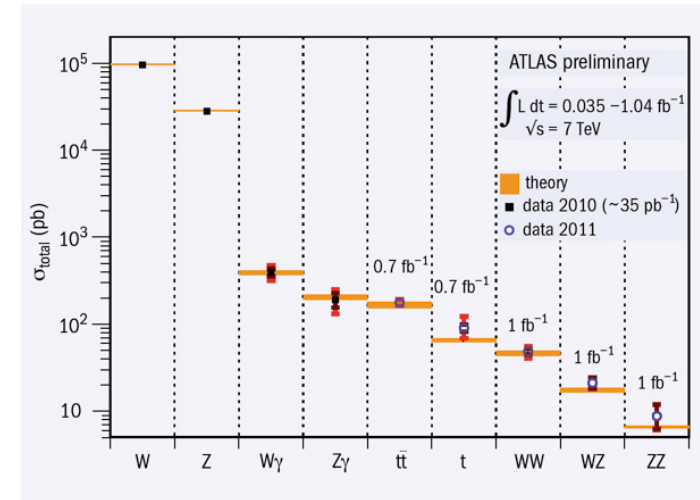
Standard Model Total Production Cross Section Measurements

Status: February 2022



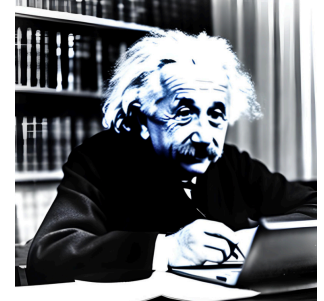
13 orders of magnitude

As at the end of 2011:

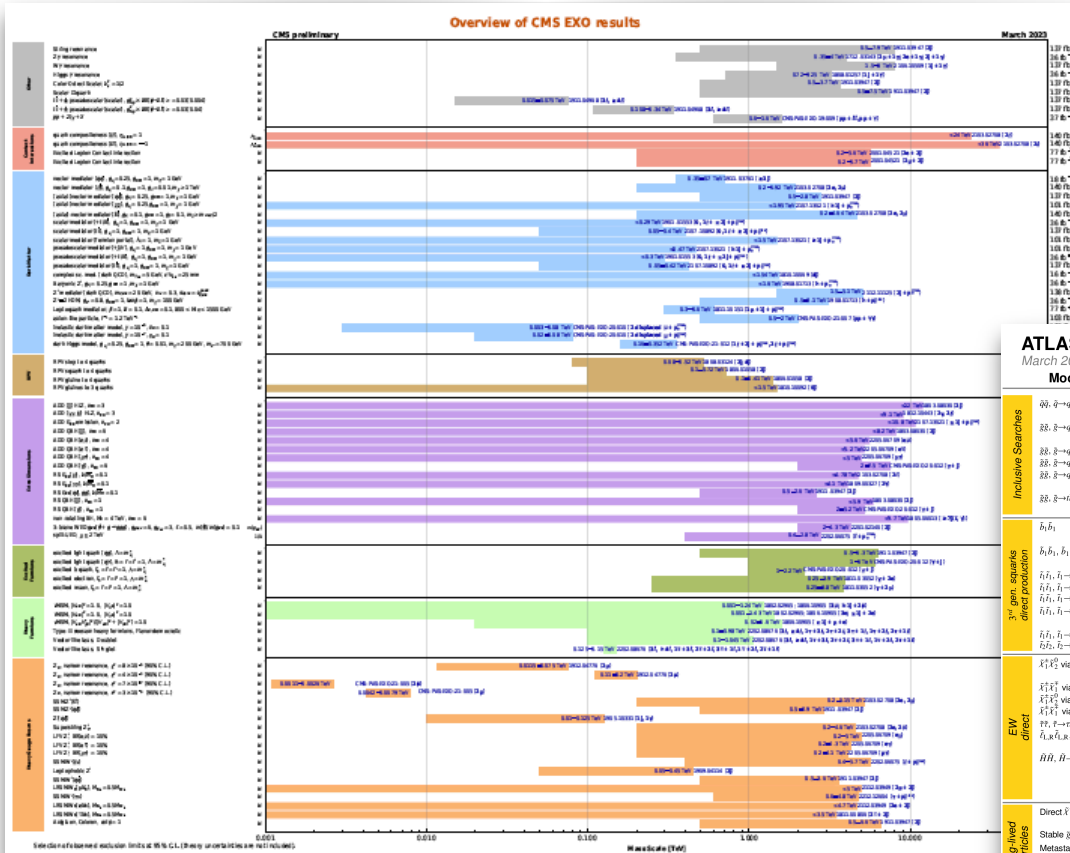


LHC Run 1 & 2: Experimental and theoretical triumph

The LHC: a machine that challenges us theorists!



(AI generated)



ATLAS SUSY Searches* - 95% CL Lower Limits
March 2023

Model	Signature	$\mathcal{L} dt [fb^{-1}]$	Mass limit	Reference		
Inclusive Searches	$\tilde{g}, \tilde{q} \rightarrow \tilde{q} \tilde{g}$	0 e, μ	2-6 jets E_{T}^{miss}	139 \tilde{g} [16, 18, 20, 21]	$m(\tilde{t}_1) < 400$ GeV 2010.14293	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	monojet	1-3 jets E_{T}^{miss}	139 \tilde{g} [16, 18, 20, 21]	$m(\tilde{t}_1) < 500$ GeV 2102.10874	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	0 e, μ	2-6 jets E_{T}^{miss}	139 \tilde{g}	Forbidden 2010.14293	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	1 e, μ	2-6 jets E_{T}^{miss}	139 \tilde{g}	$m(\tilde{t}_1) < 600$ GeV 2101.01629	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	ee, $\mu\mu$	2 jets E_{T}^{miss}	139 \tilde{g}	$m(\tilde{t}_1) < 700$ GeV 2204.13072	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	0 e, μ	7-11 jets E_{T}^{miss}	139 \tilde{g}	$m(\tilde{t}_1) < 600$ GeV 2008.06032	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	SS e, μ	6 jets	139 \tilde{g}	$m(\tilde{t}_1) < 200$ GeV 1909.08457	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	0 e, μ	3 b jets E_{T}^{miss}	139 \tilde{g}	$m(\tilde{t}_1) < 500$ GeV 2211.08028	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	SS e, μ	6 jets	139 \tilde{g}	$m(\tilde{t}_1) < 300$ GeV 1909.08457	
	$\tilde{g}, \tilde{g} \rightarrow \tilde{g} \tilde{g}$	0 e, μ	2 b jets E_{T}^{miss}	139 \tilde{g}	$m(\tilde{t}_1) < 400$ GeV 2101.12527	
3 rd gen. squarks direct production	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	0 e, μ	6 b jets E_{T}^{miss}	139 \tilde{b}_1	2101.12527	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	2 τ	2 b jets E_{T}^{miss}	139 \tilde{b}_1	2101.12527	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	0 e, μ	6 b jets E_{T}^{miss}	139 \tilde{b}_1	Forbidden 1908.03122	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	0 e, μ	2 b jets E_{T}^{miss}	139 \tilde{b}_1	Forbidden 2103.08189	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	0 e, μ	2 b jets E_{T}^{miss}	139 \tilde{b}_1	Forbidden 2012.03799	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	1 e, μ	3 jets 1 b E_{T}^{miss}	139 \tilde{b}_1	$m(\tilde{t}_1) < 500$ GeV 2108.07665	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	1-2 τ	2 jets 1 b E_{T}^{miss}	139 \tilde{b}_1	Forbidden 1805.01649	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	0 e, μ	2 c jets E_{T}^{miss}	361 \tilde{b}_1	$m(\tilde{t}_1) < 600$ GeV 2102.10874	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	0 e, μ	monojet E_{T}^{miss}	139 \tilde{b}_1	$m(\tilde{t}_1) < 500$ GeV 2006.05880	
	$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	1-2 e, μ	1-4 b jets E_{T}^{miss}	139 \tilde{b}_1	Forbidden 2006.05880	
$\tilde{b}_1, \tilde{b}_1 \rightarrow \tilde{b}_1 \tilde{b}_1^* \rightarrow \text{hadrons}$	3 e, μ	1 b jets E_{T}^{miss}	139 \tilde{b}_1	Forbidden 2006.05880		
EW direct	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	Multiple ℓ jets	≥ 1 jet E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0$, wino-bino 2106.0676, 2108.07586	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	ee, $\mu\mu$	≥ 1 jet E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 5$ GeV, wino-bino 1911.12606	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	2 e, μ	0 jets E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0$, wino-bino 1908.08215	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	Multiple ℓ jets	Multiple ℓ jets E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 70$ GeV, wino-bino 2004.10084, 2108.07586	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	2 e, μ	0 jets E_{T}^{miss}	139 \tilde{t}_1	Forbidden 1908.08215	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	2 e, μ	0 jets E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0$ 1911.06660	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	2 τ	0 jets E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0$ 1908.08215	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	2 e, μ	≥ 1 jet E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 10$ GeV 1911.12606	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	2 τ	0 jets E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0$ 1908.08215	
	$\tilde{t}_1 \tilde{t}_1^* \rightarrow WZ$	2 e, μ	0 jets E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0$ 1911.12606	
Long-lived particles	Direct $\tilde{t}_1 \tilde{t}_1^*$ prod., long-lived \tilde{t}_1	Disapp. trk	1 jet E_{T}^{miss}	139 \tilde{t}_1	Pure Wino 2001.02472	
	Stable \tilde{t}_1 R-hadron	pixel dEdx	E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 100$ GeV 2005.06013	
	Metastable \tilde{t}_1 R-hadron, $\tilde{t}_1 \rightarrow \tau \tilde{t}_1$	pixel dEdx	E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0.1$ ns 2005.06013	
	$\tilde{t}_1, \tilde{t}_1 \rightarrow \tau \tilde{t}_1$	Diagn. lep	E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 0.1$ ns 2011.07812	
	$\tilde{t}_1, \tilde{t}_1 \rightarrow \tau \tilde{t}_1$	pixel dEdx	E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 10$ ns 2005.06013	
	RPV	$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$	3 e, μ	0 jets E_{T}^{miss}	139 \tilde{t}_1	Pure Wino 2011.0543
		$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$	4 e, μ	0 jets E_{T}^{miss}	139 \tilde{t}_1	$m(\tilde{t}_1) < 200$ GeV 2103.11684
		$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$	Multiple	4-5 large jets E_{T}^{miss}	361 \tilde{t}_1	Large \tilde{t}_1 1804.05080
		$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$	Multiple	Multiple jets E_{T}^{miss}	361 \tilde{t}_1	$m(\tilde{t}_1) < 200$ GeV, bino-bino ATLAS-COUP-2018-003
		$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$	$\geq 4b$	Multiple jets E_{T}^{miss}	361 \tilde{t}_1	Forbidden 2010.01015
$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$		2 jets + 2 b	2 jets + 2 b E_{T}^{miss}	361 \tilde{t}_1	$m(\tilde{t}_1) < 500$ GeV 1710.07711	
$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$		2 e, μ	2 b E_{T}^{miss}	361 \tilde{t}_1	$m(\tilde{t}_1) < 100$ GeV, bino-bino 1710.05544	
$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$		1 μ	DV E_{T}^{miss}	136 \tilde{t}_1	$m(\tilde{t}_1) < 100$ GeV, $\cos(\theta) < 1$ 2003.11956	
$\tilde{t}_1 \tilde{t}_1^* \rightarrow \tau \tilde{t}_1 \tau \tilde{t}_1^* \rightarrow \tau \tau \tau \tau$		1-2 e, μ	≥ 6 jets E_{T}^{miss}	139 \tilde{t}_1	Pure Higgsino 2106.09609	

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Where do we go from here?

Big unanswered questions

Origin of the electroweak scale

Nature of Dark Matter (DM)

Strong CP problem

Matter-antimatter asymmetry

Origin of neutrino masses

Flavor problem

the known
unknown

Old questions but **remarkable progress** has been done
in the past ~decade

Big unanswered questions

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Strong CP problem

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Origin of neutrino masses

Flavor problem

the known
unknown

Overarching question:
what is the unknown?

the unknown
unknown

Old questions but **remarkable progress** has been done
in the past ~decade

A diversification of the field

HEP has dramatically broadened in the past 10 years



HEP is closer than ever to other fields in physics:
gravitational waves, condensed matter, atomic physics, ...
Stronger and stronger complementarity.



What can we discover next?

Discoveries in particle physics

Particle physics is **not only about discovery new particles.**

It's about the laws of nature, which include the interactions and properties of the particles that we have already discovered.

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In this sense, we had several “Higgs discoveries” after the 2012 discovery:

- * Latest Higgs discovery:

Higgs interacting with **bottom quarks**

([August 2018](#), ATLAS: Phys. Lett. B 786 (2018) 59, CMS: PRL 121 (2018) 121801)

- * Latest Higgs evidence for:

$H \rightarrow \mu\mu$ ([July 2020](#), ATLAS: Phys. Lett. B 812 (2021) 135980, CMS: JHEP 01 (2021) 14)

$H \rightarrow ll \gamma\gamma$ ([March 2021](#), ATLAS: 2103.10322)

Discoveries in particle physics

Part 2
of this talk

Particle physics is **not only about discovery new particles.**

It's about the laws of nature, which include the interactions and properties of the particles that we have already discovered.

Part 1
of this talk

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Part 1 of this talk

Particle physics is about the laws of nature, which include the interactions and properties of the particles that we have already discovered.



The future Higgs discoveries



“The Higgs is SM-like” but...

We need to understand if the Higgs

- * interacts with the 2nd generation
- * interacts with itself
- * is CP violating
- * interacts with DM/a dark sector
- * interacts with new Higgs bosons

...

Unanswered questions

- the flavor puzzle
- the matter antimatter asymmetry, the origin of the electroweak scale
- the matter antimatter asymmetry
- the origin of DM, the strong CP problem(?)
- the origin of the electroweak scale

The future Higgs discoveries



“The Higgs is SM-like” but...

We need to understand **if the Higgs interacts with the 2nd generation**

Before the Higgs discover, no evidence for Yukawa force between fundamental particles

Now, we have established it and we are **eagerly awaiting for the discovery of the muon yukawa!**
(the first coupling to 2nd generations!)

The future Higgs discoveries



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~2-3 σ evidence at Run II

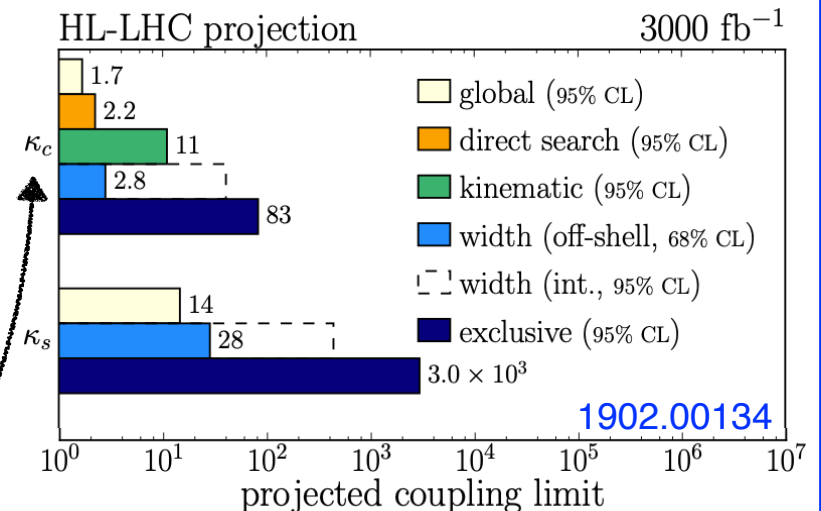
Muon: Expected discovery at Run III.

~5% level measurement at the HL-LHC

Charm:

$|k_{cl}| < 8.5(12.4)$ ATLAS: [2201.11428](#)

$1.1 < |k_{cl}| < 5.5 (< 3.4)$ CMS: [2205.05550](#)



Lot of theory effort proposing new methods to explore this Yukawa

Models that ameliorate the **flavor puzzle** can predict an enhancement of second generation couplings (and all other couplings SM-like), e.g. flavorful 2HDM [1507.07927](#), [1508.01501](#), [1908.11376](#)

The future Higgs discoveries

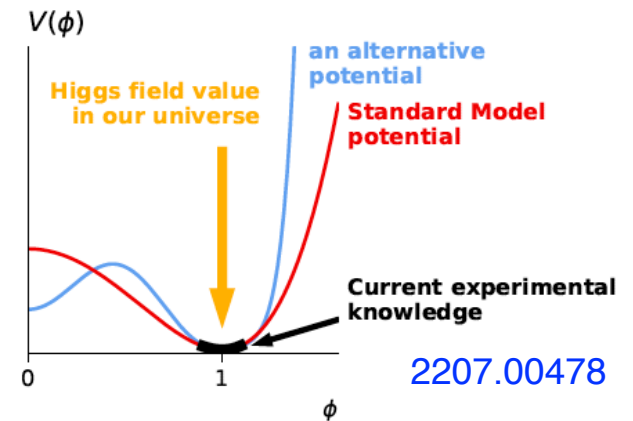
“The Higgs is SM-like” but...

We need to understand **if the Higgs interacts with itself**

In the SM, the Higgs self-interactions are fully determined:

$$V(h) = \frac{m_h^2}{2}h^2 + \frac{m_h^2}{2v}h^3 + \frac{m_h^2}{8v^2}h^4$$

First self-interacting fundamental particle ever seen in Nature?



The future Higgs discoveries

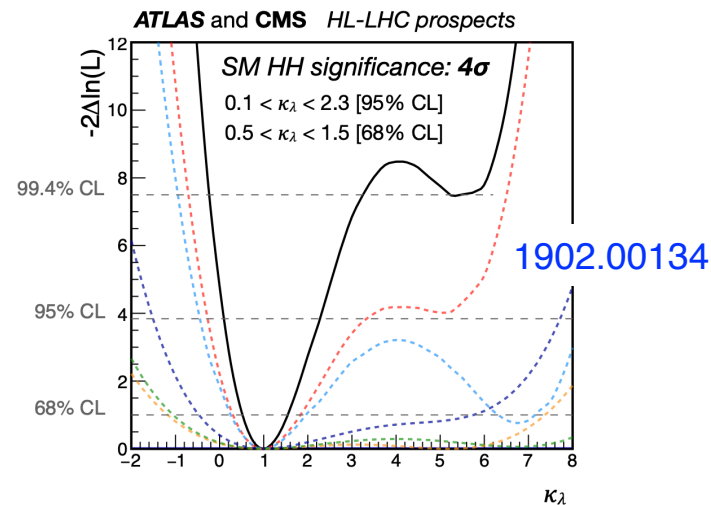
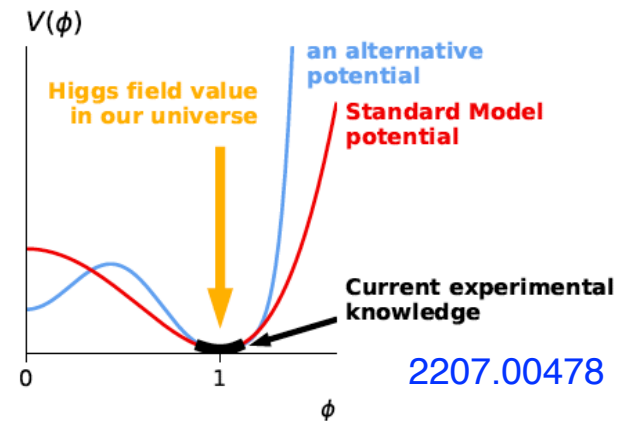
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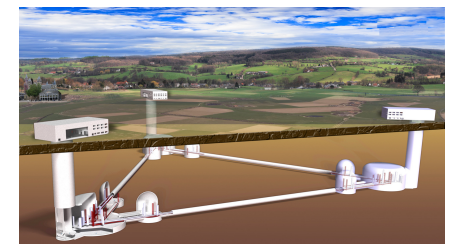
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Hope for a discovery?

Deep implications on the fate of the universe

If the phase transition that led to electroweak symmetry breaking is different than in the SM, chances to have observation in gravitational wave detectors



The Higgs precision program (theory)

“The Higgs is SM-like” but...**Do we really know?**

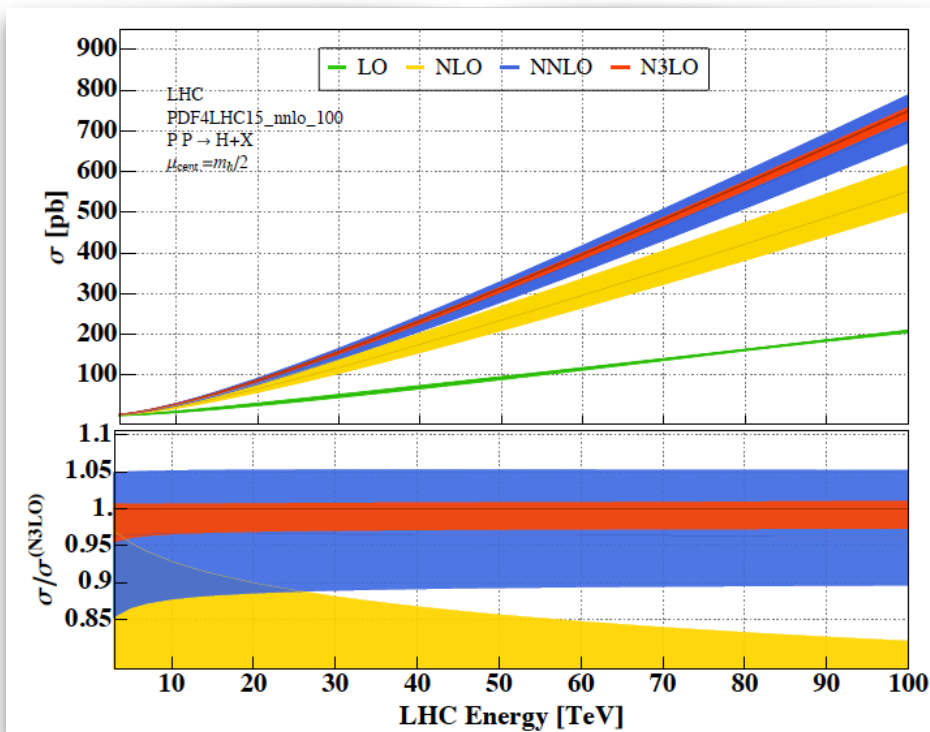
The measurement of the Higgs boson productions and decays represent a concrete deliverable for present and future collider projects.

The Higgs precision program (theory)

“The Higgs is SM-like” but...Do we really know?

Theory uncertainties have decreased significantly over the years.

2209.06138



2015

2002

1991

1978

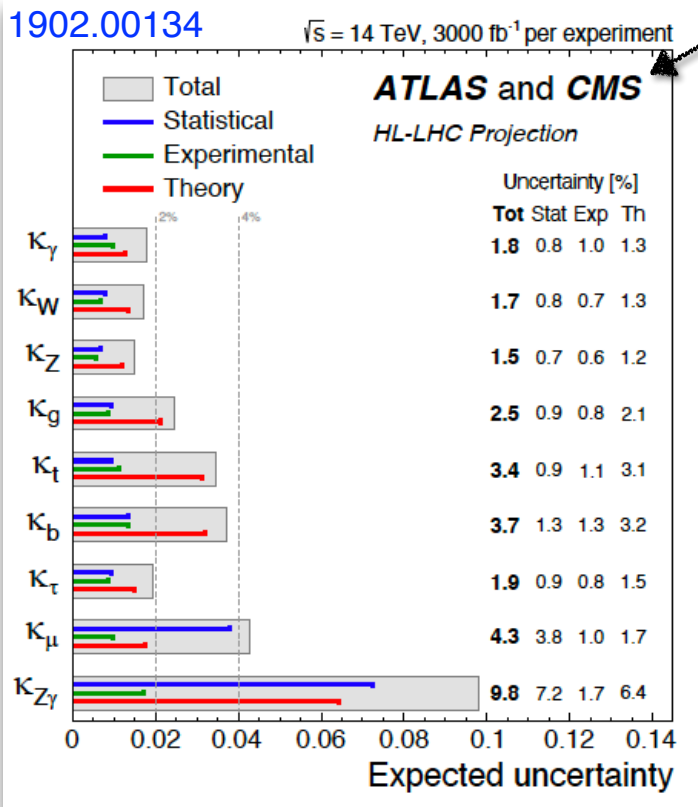
Unprecedented accuracy is needed to interpret future data.

The measurement of the Higgs boson productions and decays represent a concrete deliverable for present and future collider projects.

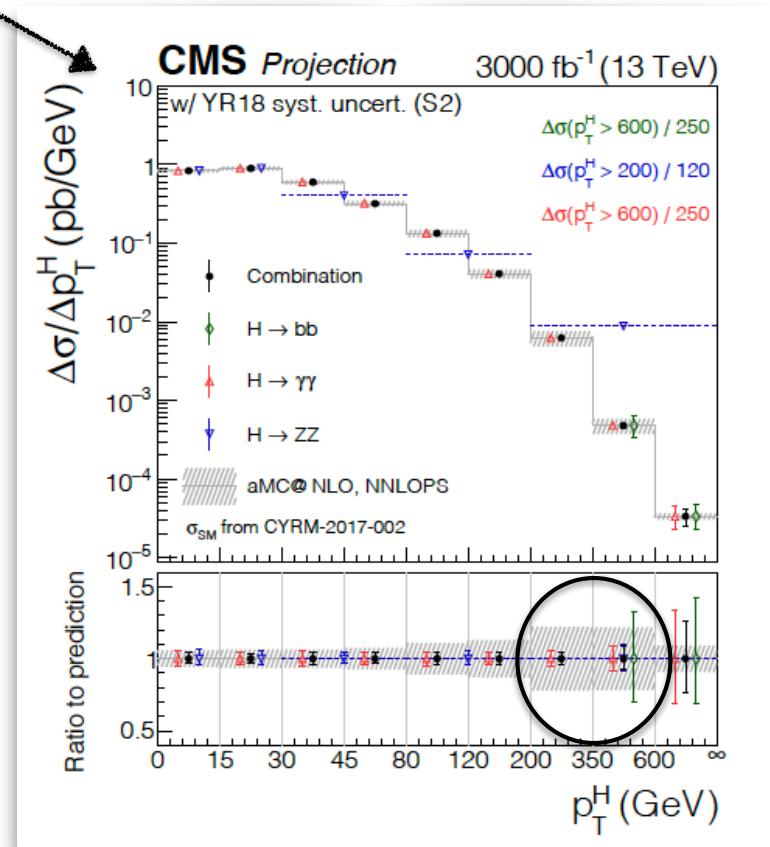
The Higgs precision program (theory+exp)

We need to go further! Both in measurements and in theory predictions

Kappa framework and differential distributions



Projections assume theory uncertainties are halved!



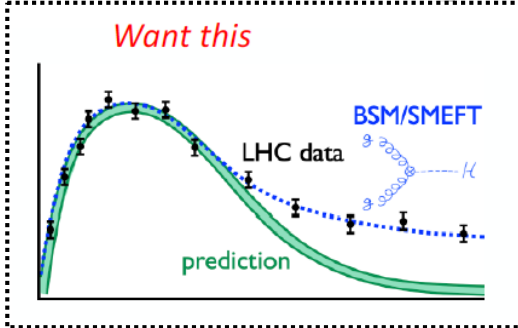
1902.00134

Learning about the Higgs at its best

Higgs precision measurements and EFTs

$$\text{Experiment} = \text{Theory}_{\text{SM}} + \sum_i \frac{C_i^6}{\Lambda^2} + \dots$$

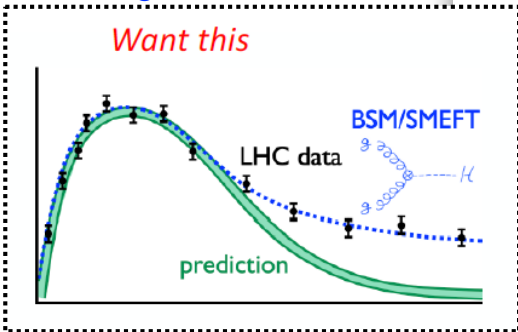
Zanderighi



Higgs precision measurements and EFTs

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Zanderighi

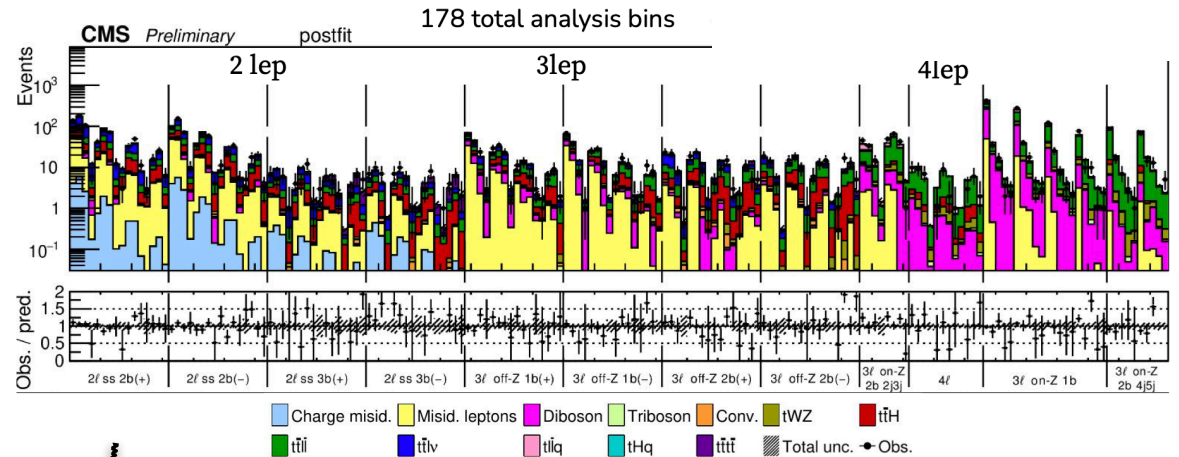


Theorists are at the forefront of SMEFT fits.
 Lot of work still needed for better understanding uncertainties.

It's not just the Higgs anymore.

- * Higgs data,
 - * top quark data,
 - * gauge boson pair production,
 - * EW precision observables
- all contribute to SMEFT predictions.

CMS PAS TOP-22-006



Operator category	WCs
Two heavy quarks	$c_{tq}, c_{\bar{q}q}, c_{\varphi Q}^3, c_{\varphi t}, c_{\varphi tb}, c_{tW}, c_{tZ}, c_{bW}, c_{tG}$
Two heavy quarks two leptons	$c_{Q\ell}^{3(\ell)}, c_{Q\ell}^{-3(\ell)}, c_{Q\ell}^{(\ell)}, c_{t\ell}^{(\ell)}, c_{t\ell}^{S(\ell)}, c_t^{T(\ell)}$
Two light quarks two heavy quarks	$c_{Qq}^{31}, c_{Qq}^{38}, c_{Qq}^{11}, c_{Qq}^{18}, c_{tq}^1, c_{tq}^8$
Four heavy quarks	$c_{Qq}^1, c_{Qq}^8, c_{tq}^1, c_{tt}^1$

The flavor precision program

Incredible success of the LHC flavor physics program.

Many discoveries

For example:

discovery of the $B_s \rightarrow \mu\mu$ decay

first observation of CP violation in charged B meson and B_s meson decays,

first observation of CP violation in charm meson decays,

first observation of charm mixing at a single experiment,

discovery of many new hadrons

world's best measurement of CP asymmetries in both B_d mixing and B_s mixing

multifaceted determination of $b \rightarrow sll$ transitions

Theory enabling experiments and experiments enabling theory.

For example:

* many new lattice calculations (B and D-meson form factors, quark masses, ...)

Continuous updates of the Flavor Lattice Average Group (FLAG) website,

<http://flag.unibe.ch>

* development of very sophisticated EFT fits for $b \rightarrow sll$ transitions

The flavor precision program

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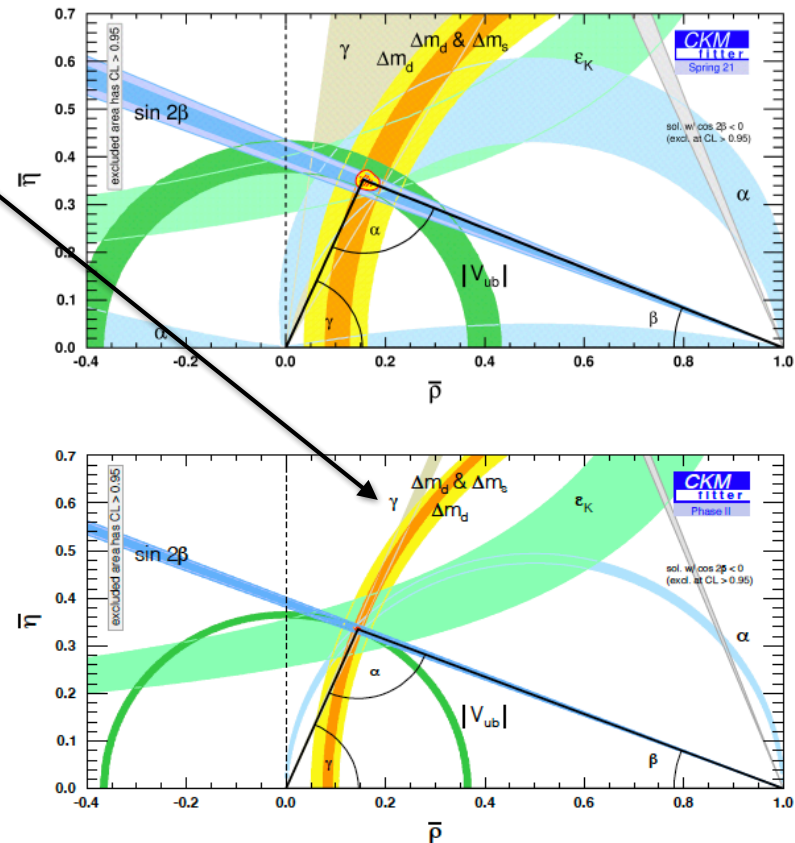
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Observable	Current LHCb	Upgrade II
EW Penguins		
$R_K (1 < q^2 < 6 \text{ GeV}^2 c^4)$	0.044 [6]	0.007
$R_{K^*} (1 < q^2 < 6 \text{ GeV}^2 c^4)$	0.12 [19]	0.009
CKM tests		
γ	4° [5]	0.35°
$\sin 2\beta$, with $B^0 \rightarrow J/\psi K_S^0$	0.04 [20]	0.003
ϕ_s , with $B_s^0 \rightarrow J/\psi \phi$	32 mrad [21]	4 mrad
ϕ_s , with $B_s^0 \rightarrow D_s^+ D_s^-$	170 mrad [23]	9 mrad
ϕ_s^{ss} , with $B_s^0 \rightarrow \phi \phi$	154 mrad [24]	11 mrad
a_{sl}^s	33×10^{-4} [26]	3×10^{-4}
$ V_{ub} / V_{cb} $	6% [27]	1%
$B_s^0, B^0 \rightarrow \mu^+ \mu^-$		
$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)/\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)$	69% [4, 28]	11%
$\tau_{B_s^0 \rightarrow \mu^+ \mu^-}$	14% [4, 28]	2%
$S_{\mu\mu}$	-	0.2
$b \rightarrow c \ell^- \bar{\nu}_\ell$ LUV studies		
$R(D^*)$	0.026 [30, 31]	0.002
$R(J/\psi)$	0.24 [32]	0.02
Charm		
$\Delta A_{CP}(KK - \pi\pi)$	29×10^{-5} [7]	3.3×10^{-5}
$A_\Gamma (\approx x \sin \phi)$	11×10^{-5} [33]	1.2×10^{-5}
$\Delta x (D^0 \rightarrow K_S^0 \pi^+ \pi^-)$	18×10^{-5} [34]	1.6×10^{-5}



1808.08865



Future insights on the flavor structure of Nature

Precision flavor physics and New Physics

Historically, measuring flavor transitions led to big discoveries in particle physics

<p>* Measurement of the tiny branching ratio of $K_L \rightarrow \mu^+ \mu^-$</p> <p>➡ prediction of the charm quark (Glashow, Iliopoulos, Maiani, 1970)</p> 	<p>* Observation of CP violation in kaon anti-kaon oscillations</p> <p>➡ prediction of the 3rd generation quarks (Kobayashi, Maskawa, 1973)</p> 
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50th Anniversary
last February


<https://www-conf.kek.jp/KM50/>

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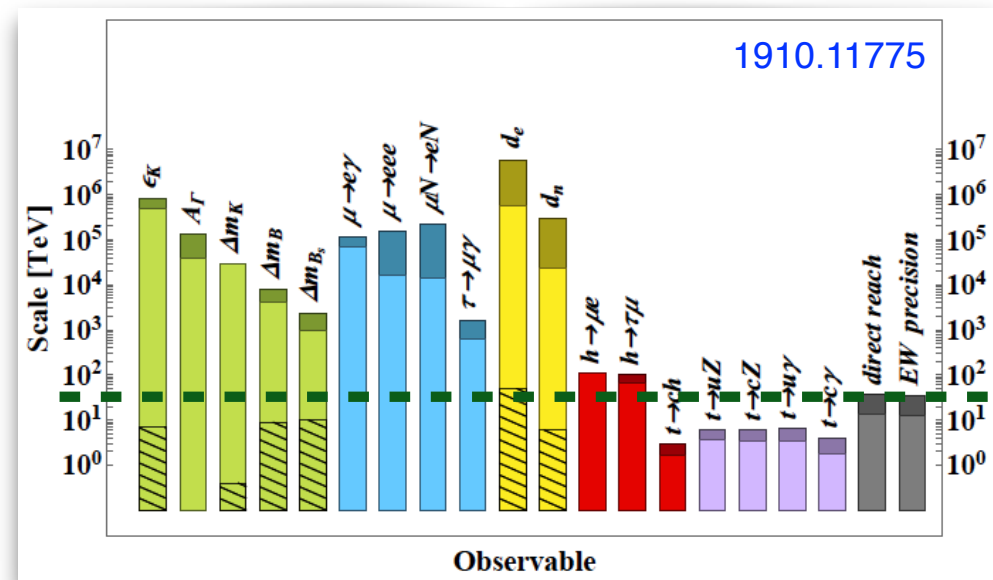
* Observation of CP violation in kaon anti-kaon oscillations

➔ prediction of the 3rd generation quarks (Kobayashi, Maskawa, 1973)



Flavor transitions: access to very high New physics scales, not directly accessible at collider experiments.

➔ Future indirect discovery of a high new physics scale?



Discovery through precision measurements

We have discussed the importance of the Higgs and the flavor precision programs.

More in general,

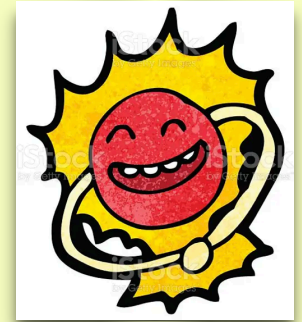
*“Precision measurements are not only a way of testing and consolidating known theories, but also an **extremely powerful tool for detecting hints of new phenomena** in a way that is complementary to and — in some cases — more far-reaching than direct exploration.”*

A roadmap for the future

CERN Director-General Fabiola Gianotti and Gian Giudice, Head of CERN's Theory Department, comment on the scientific vision and priorities for the field laid out in the recently updated European Strategy for Particle Physics

2 OCTOBER, 2020

[Nature Physics volume 16, pages 997–998 \(2020\)](#)



The direct quest for new particles

the known
unknown

Probing broad ideas motivated
by big open problems

(1. origin of EW scale, 2. Dark Matter,
3. strong CP problem)

theoretical
prior

the unknown
unknown

Searching for exotics.
Model independent searches.

no
theoretical
prior

1.

Probing broad ideas: the origin of the EW scale

Several classes of models have been proposed

* Solutions based on symmetries
(SUSY, composite Higgs, neutral naturalness, ...)



Signals for the LHC:
weak scale dynamics,
new Higgs bosons

* Solutions based on symmetries + low
cutoff (relaxion, [1504.07551](#))



Higgs portal, $\xi |H|^2 |\phi|^2$
[1610.02025](#)

* ...

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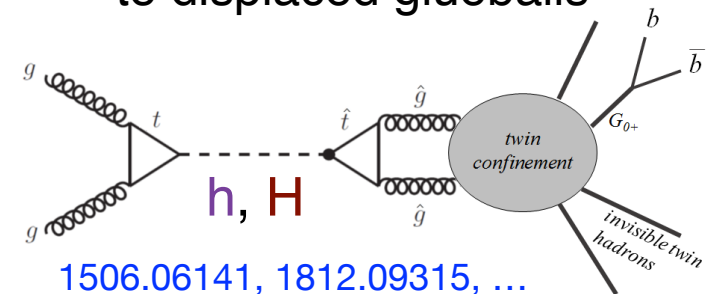
→ Higgs portal, $\xi |H|^2 |\phi|^2$
[1610.02025](#)

Lots of info from future Higgs measurements.

Motivation for new searches for **additional Higgs bosons**.

Example:

Higgs exotic decays +
Heavy Higgs decays
to displaced glueballs

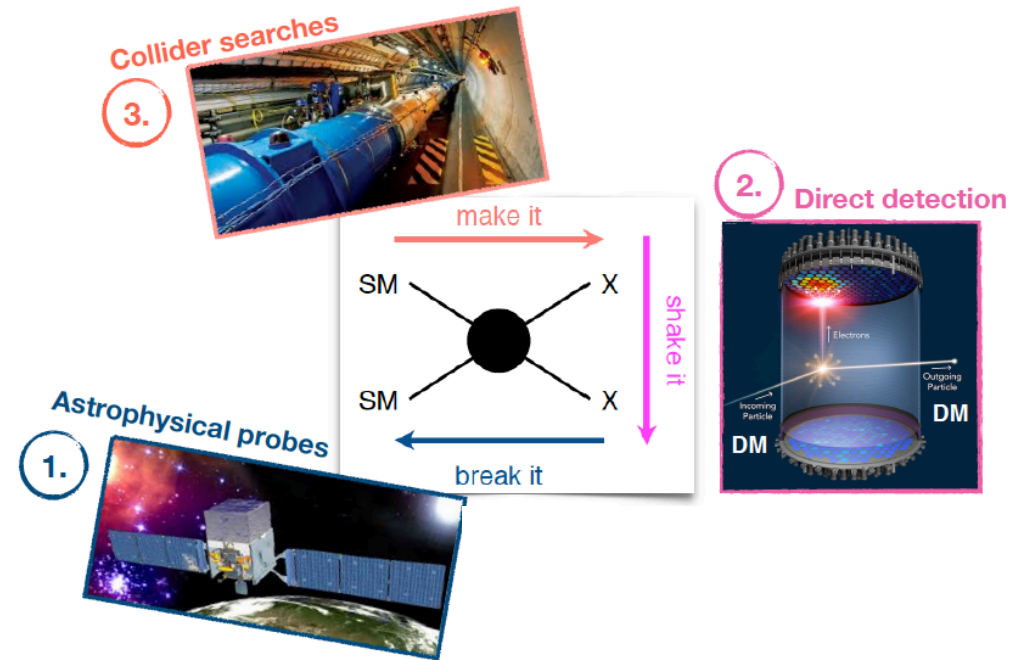


[1506.06141](#), [1812.09315](#), ...

Probing broad ideas: WIMP (present)

Idea of WIMP DM has motivated a large experimental endeavor

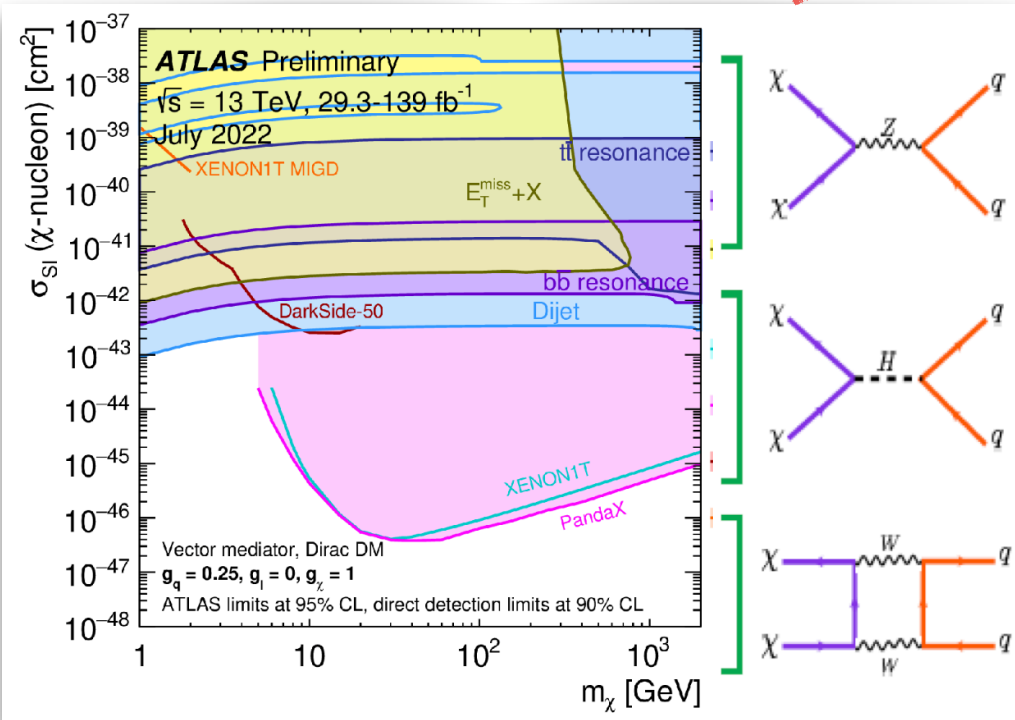
Thermal freeze-out DM at around the electroweak scale is a very predictive framework



Probing broad ideas: WIMP (present)

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Thermal freeze-out DM at around the electroweak scale is a very predictive framework



(interpretations on this plane are strongly model dependent)

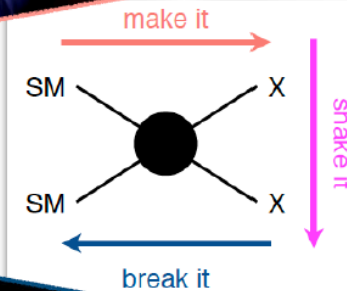
Collider searches

3.

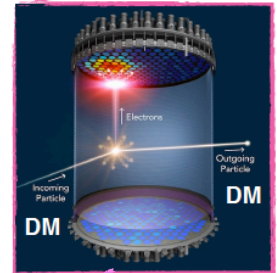


Astrophysical probes

1.



2. Direct detection



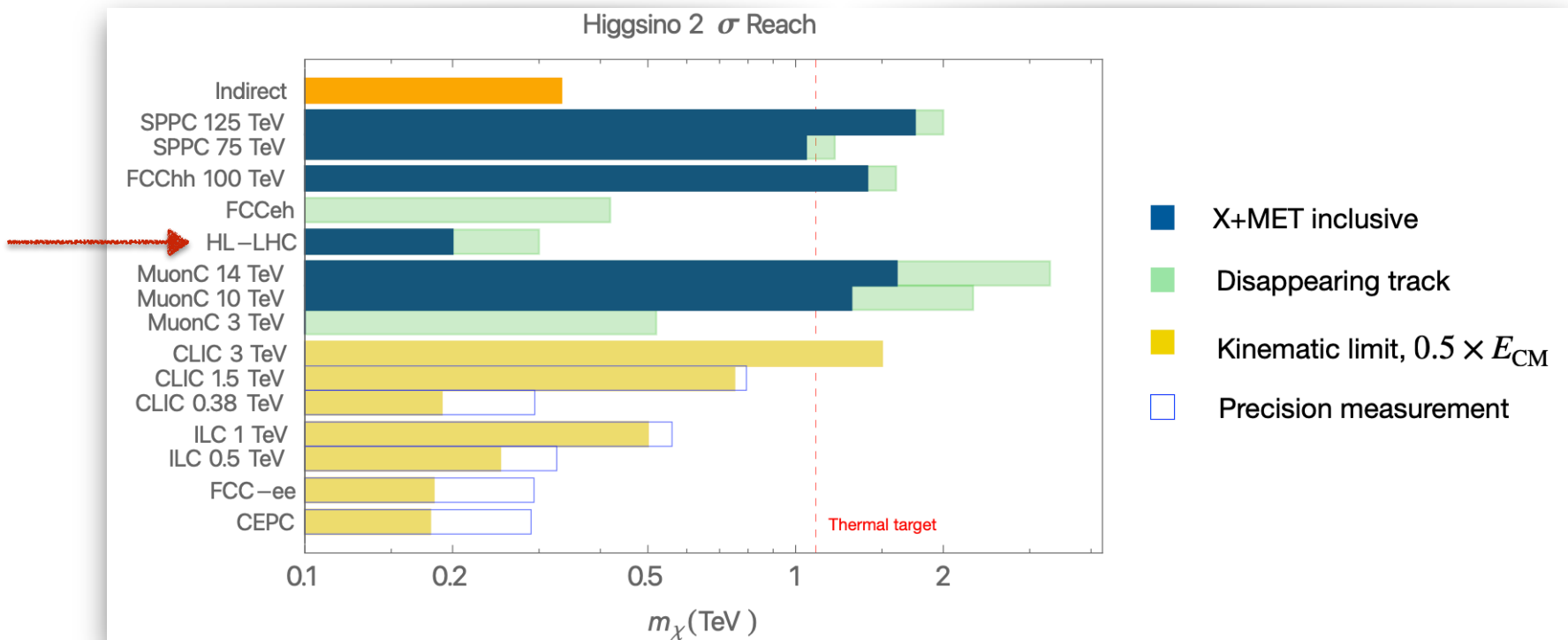
Milestone:
vanilla WIMPs have been probed!

Some models are more challenging: Higgsino DM

Probing broad ideas: WIMP (future)

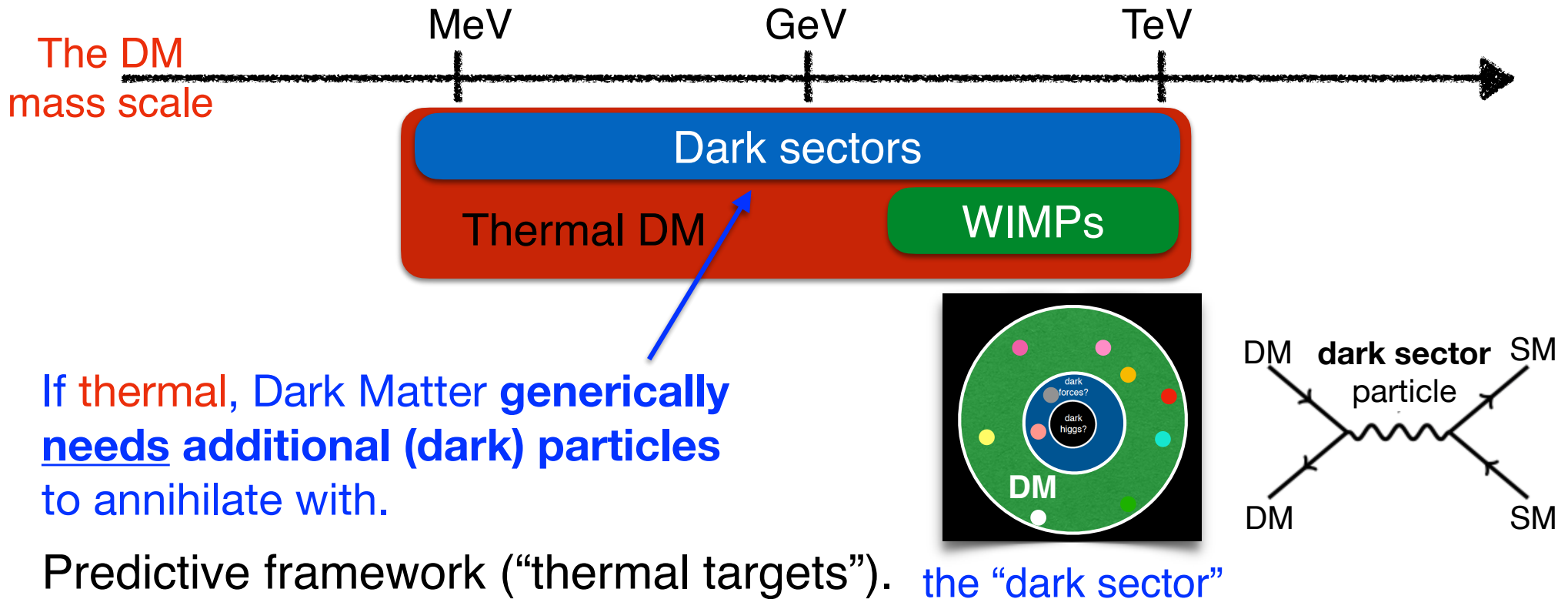
Towards a more comprehensive coverage...

Higgsinos at future colliders

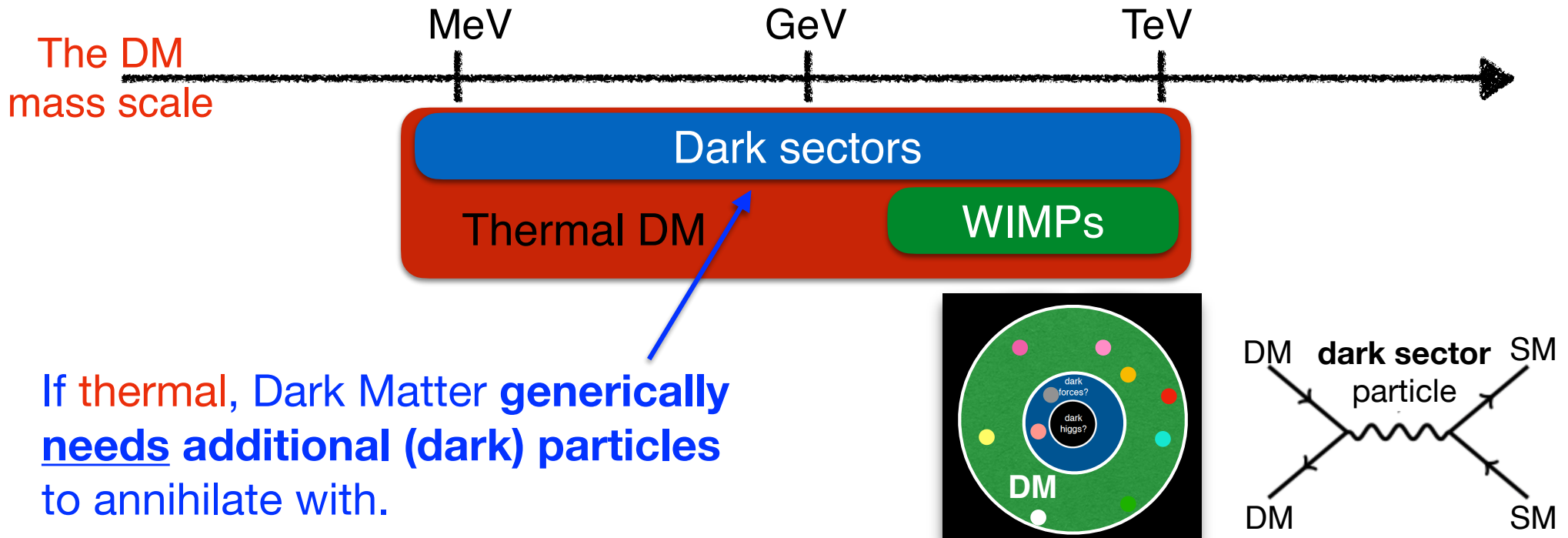


Snowmass, 2211.11084

A broader framework for thermal DM



A broader framework for thermal DM



Predictive framework (“thermal targets”). the “dark sector”

Several LHC search strategies to have access to the dark sector:

- * Higgs exotic decays
- * Long lived particles
- * Low mass resonances
- * ...

+ **New search strategies for DM direct detection experiments.**

Collaboration with condensed matter physicists.

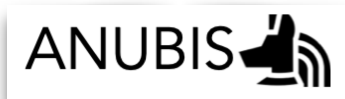
The dark sector: a guideline for new experiments

The physics of dark sectors is vast and motivated by many open problems in particle physics (not only DM, but also the strong CP problem, neutrino masses, ...).

This physics has motivated the flourishing of proposals for new experiments and detectors.

The intense collaboration between theorists and experimentalists enabled these proposals to happen.

Work at the interface between theory and experiment should be encouraged.



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Work at the interface between theory and experiment should be encouraged.

Physics beyond colliders initiative (<https://pbc.web.cern.ch>) & Feebly Interacting Particles Physics Centre.

Complementarity with international efforts

(B-factories in Japan, neutrino experiments & fixed target experiments in the US, ...)

The dark sector: a guideline for new experiments

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Work at the interface between theory and experiment should be encouraged.

In the coming decade, we will have the opportunity to broadly probe the thermal freeze-out framework (similarly to what we have done in the past decade with WIMPs)

Probing broad ideas: axions and axion-like-particles

Strong CP problem:

why is the QCD θ parameter so small? $\mathcal{L}_{\text{QCD}} \supset \bar{\theta} \frac{g^2}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu}$

QCD axion: elegant way to address this problem.

Dynamical solution to achieve: $\bar{\theta} \lesssim 10^{-10}$

in agreement with EDM constraints

Axions are a popular solution, but not much for the LHC to say.

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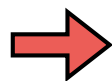
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Axions are a popular solution, but not much for the LHC to say.

Non minimal models with heavier axion-like-particles (ALPs) have been proposed to address the strong CP problem. (e.g., [1710.04213](#); [2208.10504](#); ...)

The ALP EFT at dim-5:

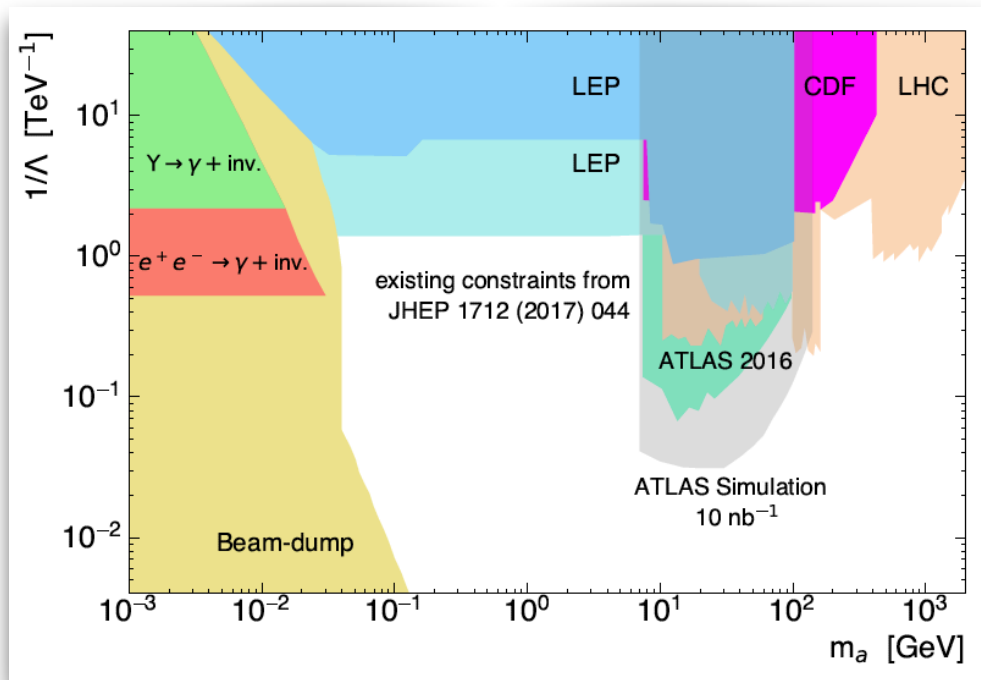
$$\mathcal{L} \supset -\frac{g_{ag}}{4} a G_{\mu\nu}^a \tilde{G}^{a\mu\nu} - \frac{g_{aW}}{4} a W_{\mu\nu}^a \tilde{W}^{a\mu\nu} - \frac{g_{aB}}{4} a B_{\mu\nu} \tilde{B}^{\mu\nu} + ig_{af} (\partial_\mu a) (\bar{f} \gamma^\mu \gamma_5 f)$$



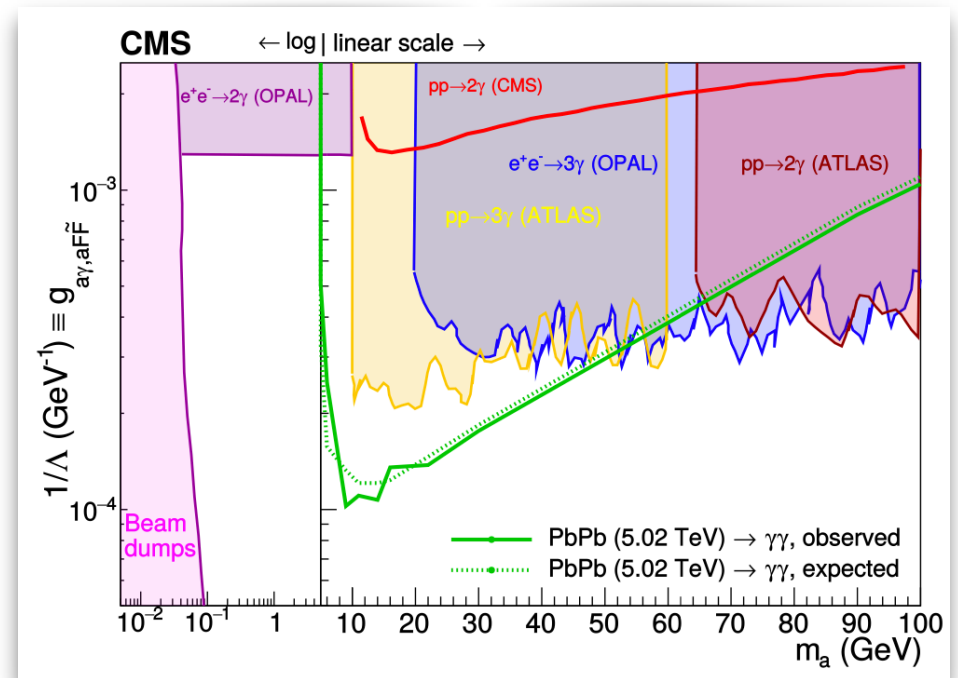
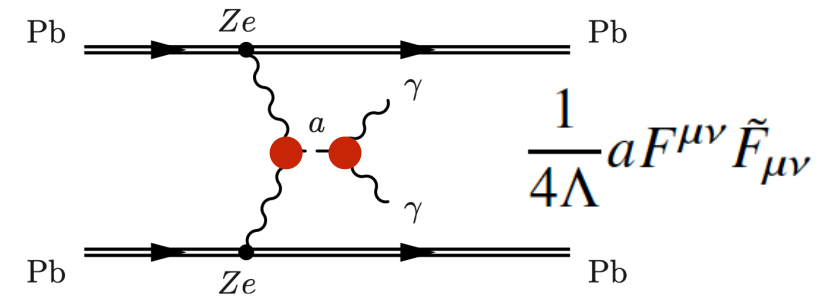
Plenty of new opportunities for the LHC

Interplay between heavy ion physics and axion physics

Axion production in ultra-peripheral collisions of heavy ions



ATL-PHYS-PUB-2018-018



1810.04602

Theory paper: [1607.06083](#)

1.

2.

3.

Long-lived particles

Long-lived-particles generically arise in NP models.

NP particles can have a width that is suppressed by

- small mass splitting
- multi-body final states
- symmetry
- high NP scale

Examples

Models to explain the origin of the EW scale

SUSY: pure wino states

$$\delta m_\chi = M_{\chi^\pm} - M_{\chi^0} \simeq 166 \text{ MeV}$$

$$\tau_{\chi^\pm}^{-1} \propto \frac{G_F^2}{\pi} f_\pi^2 \delta m_\chi^3 \sim \frac{1}{6 \text{ cm}}$$

Twin Higgs: glueballs

Models to explain the origin of neutrino masses

Sterile neutrinos

1.

2.

3.

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Appreciable gaps

- * Singly produced LLPs
- * Low-mass (< 20 GeV)
- * LLP from Higgs decays
- * High multiplicities
- * displaced taus
- * small displacements
- * ...

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Sterile neutrinos

A flourishing effort in unconventional/exotic signals

We should leave no stone unturned

We do not know what is the ultimate theory of Nature

the unknown
unknown

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Just in the past year...

the unknown
unknown

Periodic resonances

[2305.10894](#)

Forward proton scattering
in association with
light-by-light scattering
mediated by an axion-like
particle

[2304.10953](#)



Heavy long-lived multi-
charged particles
producing anomalously
high ionization

[2303.13613](#)

Fractionally charged particles with
energy loss in the tracking detector

[CMS PAS EXO-19-006](#)

Long-lived particles using trackless and
out-of-time jet information

[2212.06695](#)

Creativity is important!

Let the data teach us

No clear evidence of New Physics from direct searches

→ Fully model independent searches are more and more motivated

Looking at Nature with
no theoretical prior.

Let the data teach us

No clear evidence of New Physics from direct searches

➔ Fully model independent searches are more and more motivated

1. Generic searches:

Phase-space not tailored to a specific model.
Comparison with Standard Model predictions.

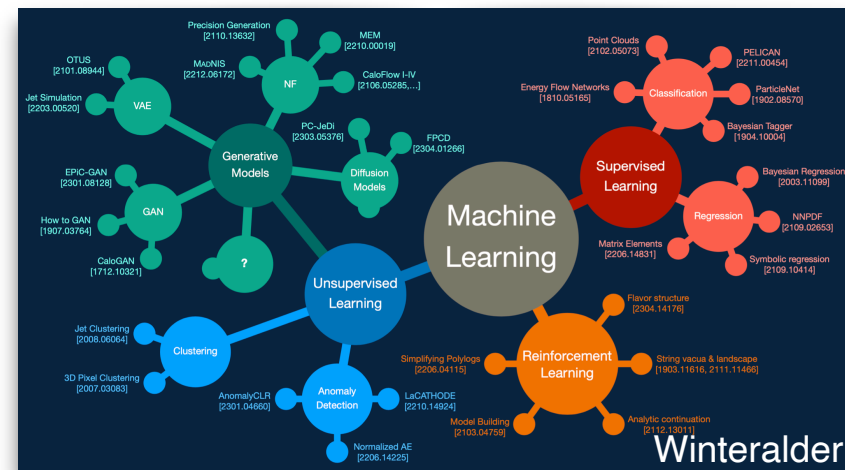
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2. Data itself can teach us!

Deep learning revolution since a few years

Machine learning (ML) for HEP

e.g. unsupervised learning for anomaly detection



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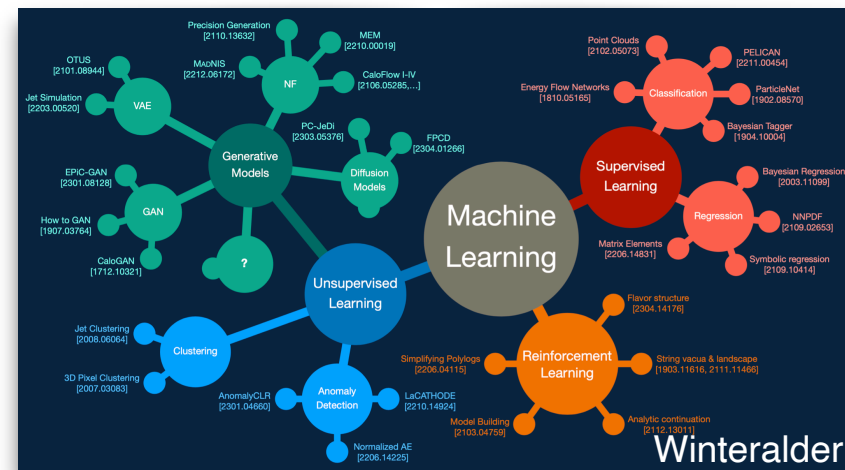
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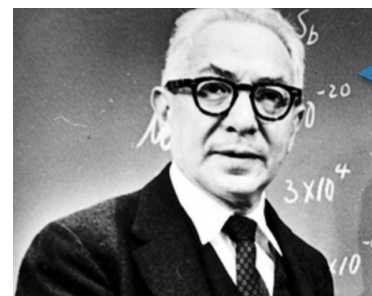
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e.g. unsupervised learning for anomaly detection



We should be ready for the
(theoretically) unexpected



μ : who
ordered
that?

Exploring the unknown unknown

*“Humanity’s thirst for knowledge, curiosity and spirit of exploration have always been the engines that drive particle physics. Unsurprisingly, the more we **dive into uncharted territory** the more difficult it becomes to predict what future experimental endeavours could find. **This is the very essence of research: if we knew for certain what future experiments will discover, we would not need to build them.**”*

A roadmap for the future

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Outlook & take home messages

Particle physics is **not only about discovery new particles.**

It's about the laws of nature, which include the interactions and properties of the particles that we have already discovered.

The breadth of HEP is getting larger and larger.

Learning about Nature at the most fundamental level is a global effort.

The LHC and future colliders will play an essential role in this endeavor.

Looking forward to the future of HEP!

