

Status and Outlook for B -mixing parameters on the lattice

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Outline

- ① Introduction
- ② Challenges in b -physics on the lattice
- ③ Status: Recent results
- ④ Outlook: Ongoing Work by RBC/UKQCD/JLQCD
- ⑤ Conclusion

Flavour Physics: The Cabibbo-Kobayashi-Maskawa matrix

experiment \approx CKM-factors \times non-perturbative inputs \times known terms

CKM Matrix

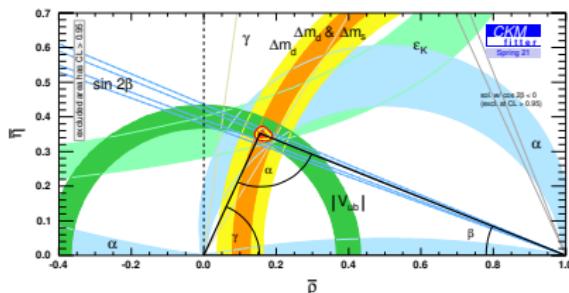
- parameterises transitions up-type \leftrightarrow down-type

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

- fundamental SM parameters
- unitary matrix in the SM:

$$VV^\dagger = \mathbb{1}_3$$

Unitarity Triangle

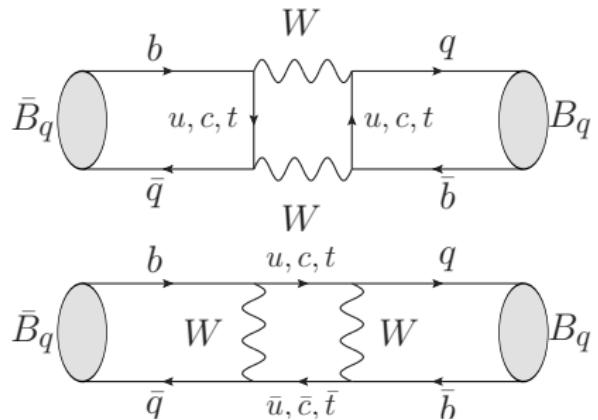


yellow & orange bands = exp. measurements of $\Delta m_{d(s)}$ + SM predictions for mixing parameters
⇒ Tests CKM unitarity!

Non-unitarity of CKM \Leftrightarrow New Physics Beyond the SM

Neutral $B_{(s)}$ meson mixing - background

Neutral mesons oscillate:



where $q = d, s$

mass eigenstate \neq flavour eigenstate

$$|B_{L,H}\rangle = p |B_q^0\rangle \pm q |\bar{B}_q^0\rangle$$

\Rightarrow **splittings** in mass eigenstates:

- mass splitting $\Delta m_q \equiv m_H - m_L$
- width splitting $\Delta\Gamma_q \equiv \Gamma_L - \Gamma_H$

Time dependence:

$$|B_q^0(t)\rangle = g_+(t) |B_q^0\rangle + \frac{q}{p} g_-(t) |\bar{B}_q^0\rangle$$

$$|\bar{B}_q^0(t)\rangle = \frac{p}{q} g_-(t) |B_q^0\rangle + g_+(t) |\bar{B}_q^0\rangle$$

Occurs at loop level in SM \Rightarrow **sensitive probe of new physics!**

Neutral $B_{(s)}$ meson mixing - experiment

$$|g_{\pm}(t)|^2 = \frac{e^{-\Gamma_q t}}{2} \left[\cosh \left(\frac{\Delta\Gamma_q}{2} t \right) \pm \cos(\Delta m_q t) \right]$$

Δm experimentally accessible as a frequency!

B_d^0 : Many results

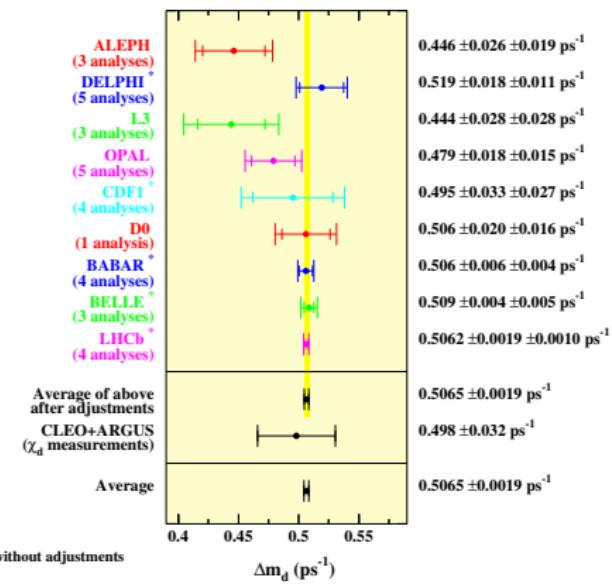
B_s^0 : "Only" CDF, CMS and LHCb

$$\Delta m_d = 0.5065(19)\text{ps}^{-1}$$

$$\Delta m_s = 17.765(06)\text{ps}^{-1}$$

Well below per cent level!

[HFLAV 2206.07501]



Neutral $B_{(s)}$ meson mixing - theory

Short distance dominated \Rightarrow described by $\mathcal{H}^{\Delta b=2}$ eff. weak Hamiltonian.
OPE factorises this into

- Perturbative model-dependent Wilson coefficients $C_i(\mu)$
- Non-perturbative model-independent matrix elements

$$\left\langle B_{(s)}^0 \right| \mathcal{H}^{\Delta b=2} \left| \bar{B}_{(s)}^0 \right\rangle = \sum_i C_i(\mu) \left\langle B_{(s)}^0 \right| \mathcal{O}_i^{\Delta b=2}(\mu) \left| \bar{B}_{(s)}^0 \right\rangle$$

- 5 independent (parity even) operators \mathcal{O}_i , only \mathcal{O}_1 relevant for Δm :

$$\mathcal{O}_1 = (\bar{b}_a \gamma_\mu (1 - \gamma_5) q_a) (\bar{b}_b \gamma_\mu (1 - \gamma_5) q_b) = \mathcal{O}_{VV+AA}$$

- Define bag parameters: $\hat{B}_{B_q}^{(i)} = \langle \bar{B}_q^0 | \mathcal{O}_i | B_q^0 \rangle / \langle \bar{B}_q^0 | \mathcal{O}_i | B_q^0 \rangle_{VSA}$

$$\Delta m_q = |V_{tb}^* V_{tq}|^2 \times f_{B_q}^2 \hat{B}_{B_q}^{(1)} \times m_{B_q} \mathcal{K}$$

\Rightarrow **Non-perturbative matrix elements calculable on the lattice**

Extracting CKM matrix elements

We write Δm_q in terms of the Renormalisation Group Independent (RGI) bag parameter \hat{B}_{B_q} :

$$\Delta m_q = |V_{tb}^* V_{tq}|^2 \frac{G_F^2 m_W^2 m_{B_q}^2}{6\pi^2} S_0(x_t) \eta_{2B} f_{B_q}^2 \hat{B}_{B_q}^{(1)}$$

- Δm_d and Δm_s are known experimentally to $\ll 1\%$ **accuracy**
- Combined other inputs are known to **0.4%**
- Typical precision for $f_{B_q} \sqrt{\hat{B}_{B_q}^{(1)}}$ is currently **a few percent**.
- Part of statistic and systematic errors cancel in $SU(3)$ breaking ratios:

$$\xi^2 \equiv \frac{f_{B_s}^2 \hat{B}_{B_s}^{(1)}}{f_{B_d}^2 \hat{B}_{B_d}^{(1)}} = \left| \frac{V_{td}}{V_{ts}} \right|^2 \frac{\Delta m_s}{\Delta m_d} \frac{m_B}{m_{B_s}} \quad \Rightarrow \text{access to } |V_{td}/V_{ts}|$$

Lattice QCD in a nutshell

Based on the **Path Integral** formulation.

$$\langle \mathcal{O} \rangle_M = \frac{1}{\mathcal{Z}} \int \mathcal{D}[\psi, \bar{\psi}, U] \mathcal{O}[\psi, \bar{\psi}, U] e^{iS[\psi, \bar{\psi}, U]}$$

Minkowski: Highly oscillatory, infinite dimensional integral. X

⇒ Wick rotate to Euclidean (i.e. imaginary) time ($t \rightarrow i\tau$).

$$\langle \mathcal{O} \rangle_E = \frac{1}{\mathcal{Z}} \int \mathcal{D}[\psi, \bar{\psi}, U] \mathcal{O}_E[\psi, \bar{\psi}, U] e^{-S_E[\psi, \bar{\psi}, U]}$$

Euclidean: Exponentially decaying, infinite dimensional integral. X

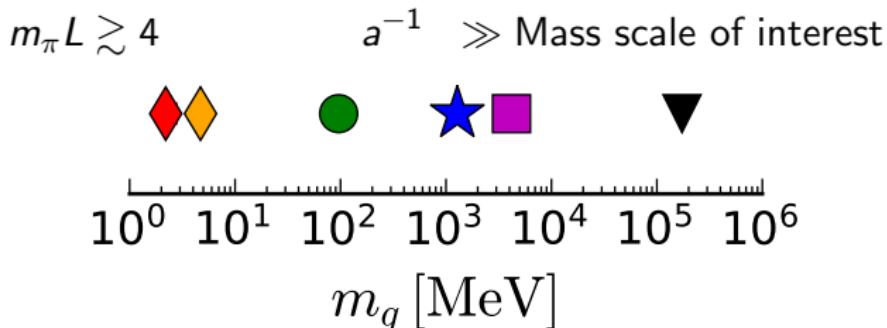
⇒ Discretise space-time and interpret as a probability distribution.

- Lattice spacing a (UV regulator)
- Box of length L (IR regulator)
- $\int \rightarrow \sum$, $\partial \rightarrow$ finite differences
- Evaluate stochastically

Lattice: Exponentially decaying and finite dimensional ✓

Multiple scale problem on the lattice: back of the envelope

Control effects of IR (finite volume) and UV (discretisation) regulators:



For $m_\pi = m_\pi^{\text{phys}} \sim 140 \text{ MeV}$ and $\bar{m}_b(m_b) \approx 4.2 \text{ GeV}$:

$$L \gtrsim 5.6 \text{ fm} \quad a^{-1} \gg 4.2 \text{ GeV} \approx (0.05 \text{ fm})^{-1}$$

Requires $N \equiv L/a \gg 120 \Rightarrow N^3 \times (2N) \gg 4 \times 10^8$ lattice sites.

VERY EXPENSIVE to satisfy both constraints simultaneously...

... needs to be repeated for different values of a .

How to simulate the b -quark?

For now choose between:

Effective action for b

- Can tune to m_b
- comes with systematic errors which are hard to estimate/reduce

Relativistic action for b

- Theoretically cleaner and systematically improvable
- Need to control extrapolation in heavy quark mass

Different properties:

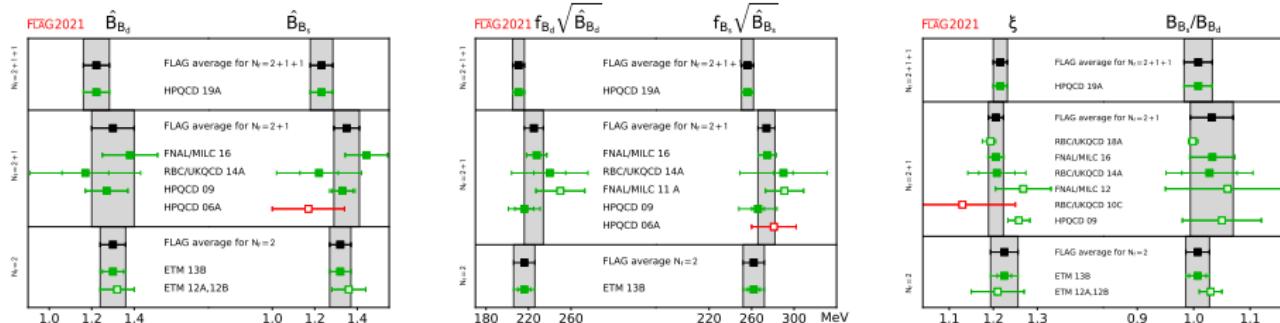
- | | | |
|----------------------|---------------------|-------------------|
| • computational cost | • tuning errors | • cut off effects |
| • chirality | • systematic errors | • renormalisation |

BUT SOON:

Huge efforts in the community to produce **very fine lattice spacings**:

⇒ **Direct simulation of $\approx m_b^{\text{phys}}$ will become possible!**

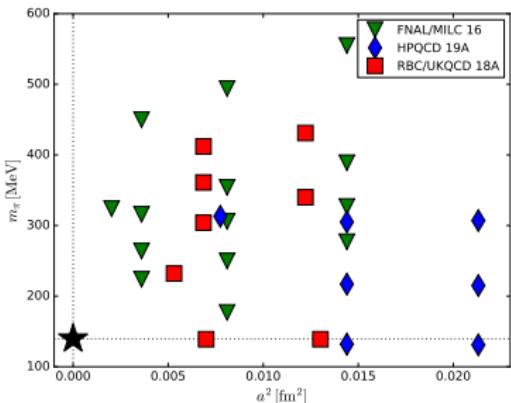
$B - \bar{B}$ mixing results: FLAG 2021 [2111.09849]



Fewer results than in the light sector,
but very complementary results from

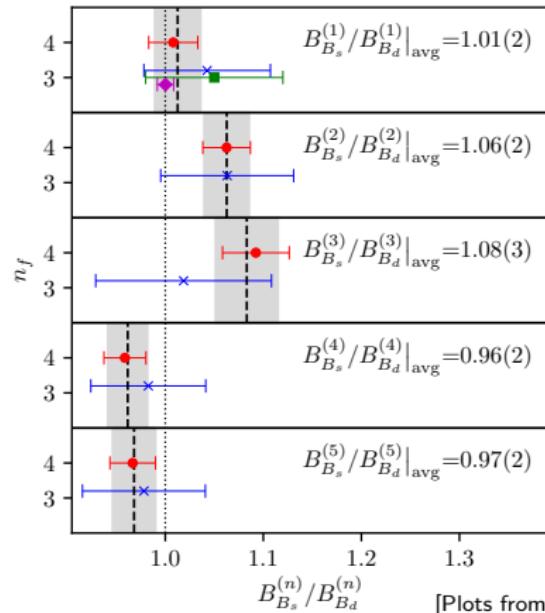
- different gauge field configurations
- different valence light actions
- different valence heavy actions
- different methodologies

Includes computations with m_π^{phys} !

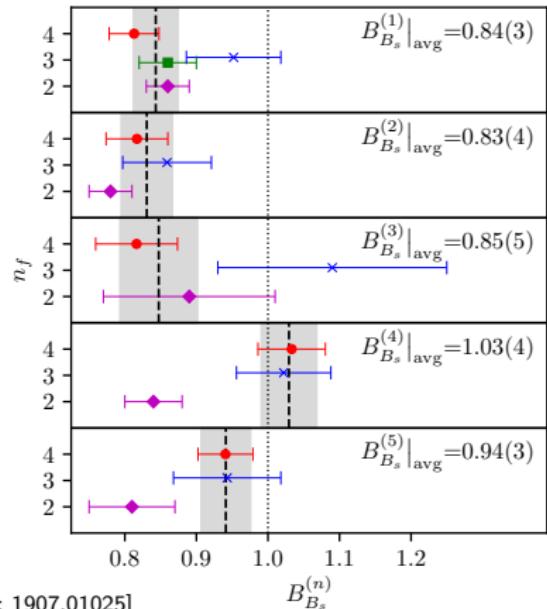


Going beyond $\hat{B}_{B_q}^{(1)}$

HPQCD19, FNAL/MILC16, HPQCD09, RBC/UKQCD18



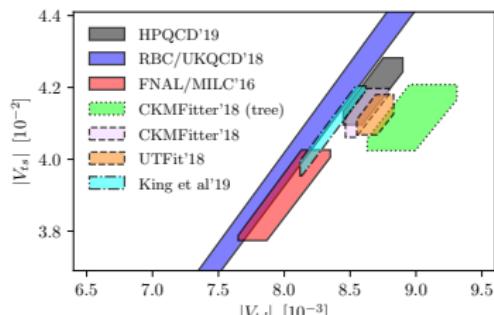
HPQCD19, FNAL/MILC16, HPQCD09, ETM13



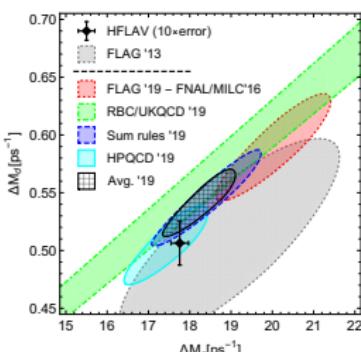
Mostly in agreement, but some spread in $B_{B_s}^{(i)}$ parameters.

CKM matrix elements

	$f_{B_d}(\hat{B}_{B_d})^{1/2}$ [GeV]	$f_{B_s}(\hat{B}_{B_s})^{1/2}$ [GeV]	ξ	$ V_{td} $	$ V_{ts} $	$ V_{ts}/V_{td} $
HPQCD19	0.2106(55)	0.2561(57)	1.216(16)	0.00867(23)	0.04189(93)	0.2071(27)
FNAL/MILC16	0.2277(98)	0.2746(88)	1.206(18)	0.00800(35)	0.0390(13)	0.2052(33)
RBC/UKQCD18			1.194(⁺¹² ₋₁₉)			0.2033 (⁺¹⁶ ₋₃₀)



[Plots from ↑ 1907.01025, ↓ 1909.11087]

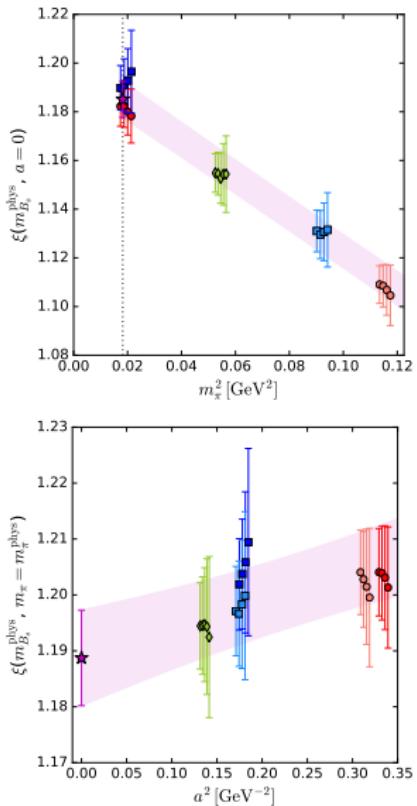
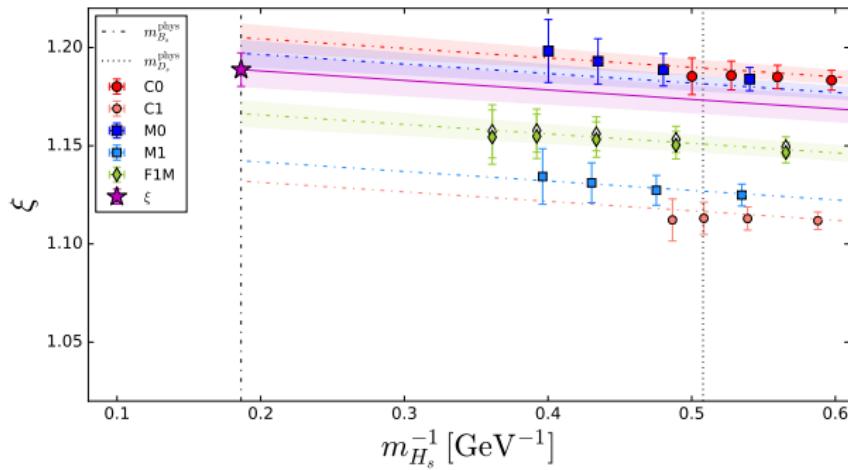


- Reasonable agreement between lattice results, but some spread
 - Tree-only fit somewhat differs
 - Uncertainty dominated by theory
- ⇒ **Further theory work required!**

Current lattice uncertainties dominated by

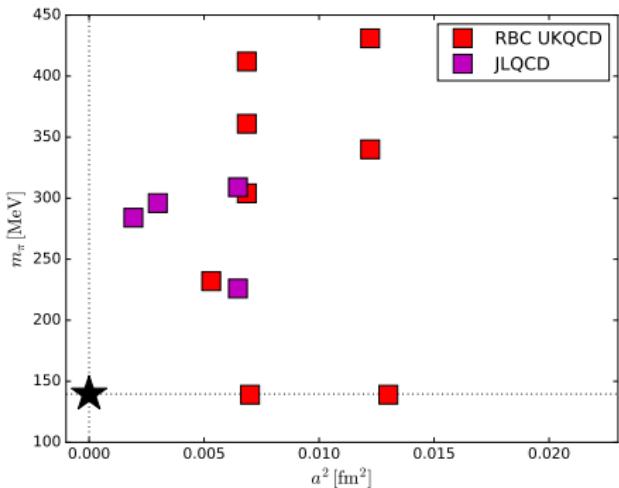
- Renormalisation/matching
 - Heavy quarks (discretisation/extrapolation)
- ⇒ we are working on improving both!

- Computation of $SU(3)$ breaking ratios f_{D_s}/f_D , f_{B_s}/f_B , B_{B_s}/B_{B_d} and ξ
- Renormalisation constants cancel
- Extrapolation from $m_h < m_b$ to m_b
- Simultaneous fit to m_π^2 , a^2 , $\frac{1}{m_H}$



RBC/UKQCD'18 dominated by

- chiral-continuum fit
- ⇒ heavy quark extrapolation
- ⇒ estimates of higher order $1/m_H$ terms



- Supplement existing dataset with finer JLQCD ensembles
⇒ **reduce extrapolation** in m_H significantly
- all domain wall fermion set-up
⇒ only **physical block-diagonal renormalisation pattern**
- all 5 operators $\hat{B}_{B_d}^{(i)}$ and $\hat{B}_{B_s}^{(i)}$
- (3+3) lattice spacings, 2 ensembles with physical pion mass
⇒ good control over all required limits
- new correlator fitting strategy

RBC/UKQCD/JLQCD: analysis strategy

Huge data-set

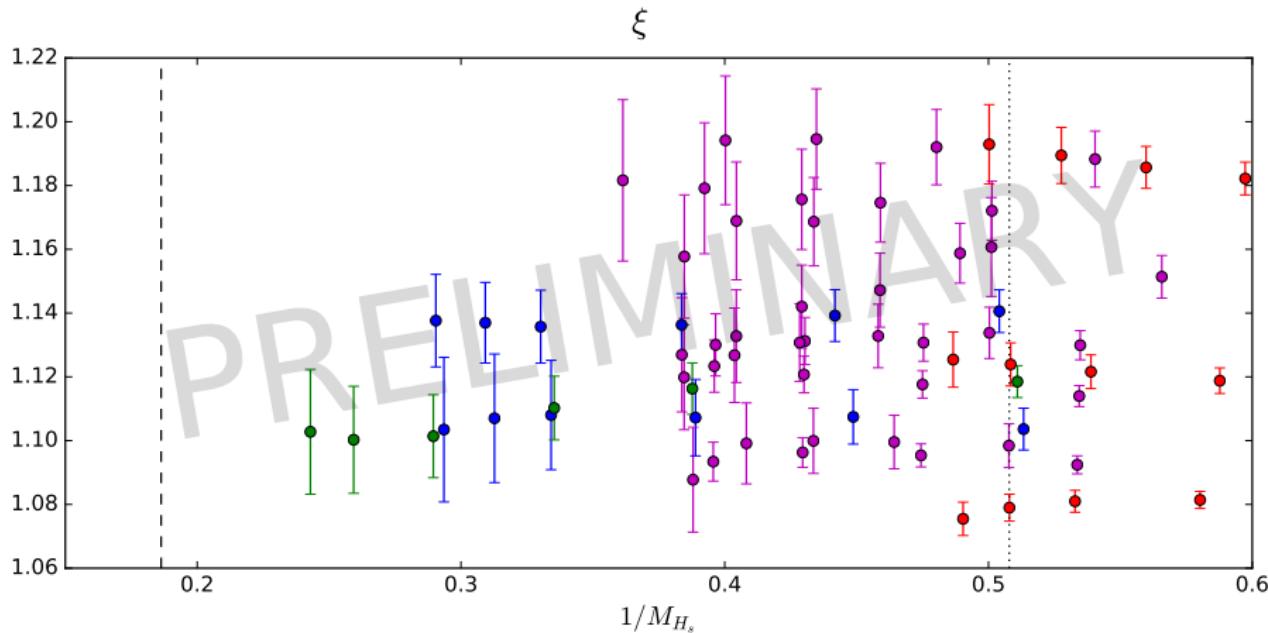
- 15 ensembles with varying $a, L, m_l (\rightarrow M_\pi), m_s$.
- wide range of heavy-quark masses (4-6) per ensemble $m_{H_s}^{\max} \sim 75\% m_{B_s}^{\text{phys}}$
- 6 lattice spacings (2 different discretisations)
- 2 m_π^{phys} ensembles.
- 2 pairs of ensembles which only differ in m_s .
- 1 pair which differs in V .

Fitting strategy & extrapolations

- Two independent correlation function analyses. (✓)
- RI/sMOM renormalisation analysis. (✓, preliminary)
- Fit to perform (ongoing)
 - $a \rightarrow 0, L \rightarrow \infty$
 - $M_\pi \rightarrow M_\pi^{\text{phys}}, m_s \rightarrow m_s^{\text{phys}}$
 - $M_{H_s} \rightarrow M_{B_s}$
- Assess robustness of fit and all systematic uncertainties.
for $\hat{B}_{B_q}^{(i)}, f_{B_q}, f_{B_q} \sqrt{\hat{B}_{B_q}^{(i)}}, \xi, \text{ ratios.}$

GOAL: Controlled prediction of all observables and their correlations

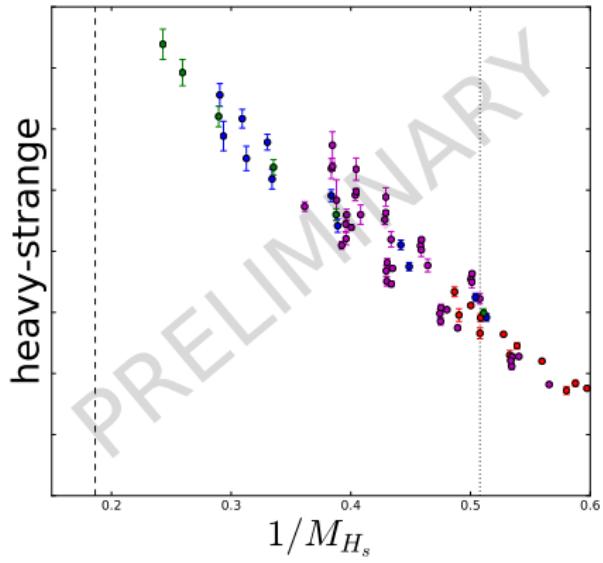
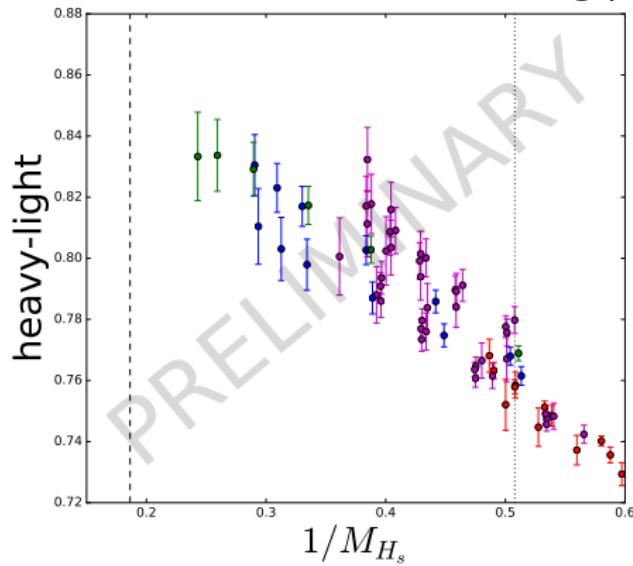
RBC/UKQCD/JLQCD: preliminary data – ratios



- Mild behaviour with $1/M_{H_s}$
- Data a lot closer to physical b -quark mass
- Precision on individual datapoints $\lesssim O(1\%)$.
- Starting to do global fits

RBC/UKQCD/JLQCD: preliminary data – bag parameters

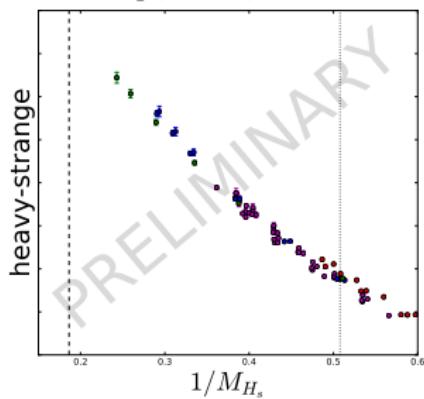
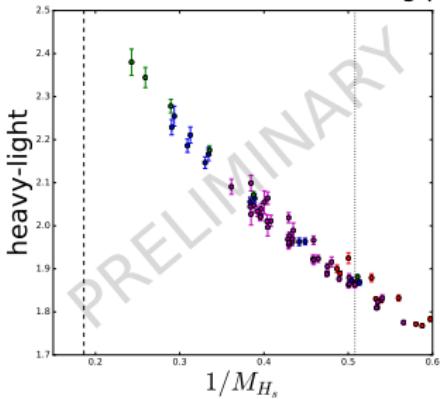
renormalised bag parameters B_1 (VV+AA)



- a^{-1}, M_π, m_s, V effects not taken into account yet....
- ... but seem to be mild
- strong $1/M_{H_s}$ behaviour but similar for all ensembles

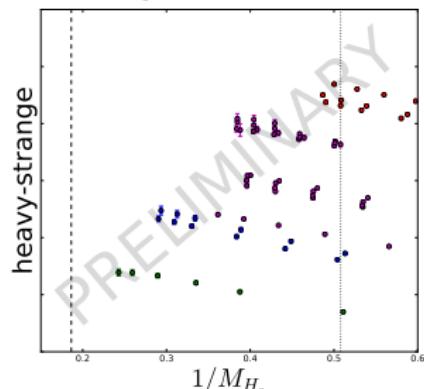
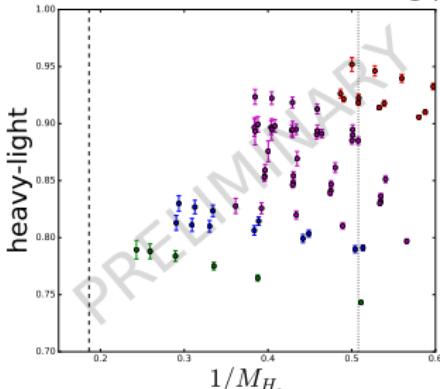
RBC/UKQCD/JLQCD: preliminary data – bag parameters

renormalised bag parameters B_2 (VV-AA)



- Mild discretisation effects
- Seemingly mild behaviour with a , M_π , m_s , V
- strong $1/M_{H_s}$ dependence

renormalised bag parameters B_3 (SS-PP)



- Large discretisation effects
- More noticeable chiral effects
- mild $1/M_{H_s}$ dependence

Summary and Outlook

Literature Summary

- Complementary lattice results:
 - ensembles
 - light quark actions
 - heavy quark treatment
 - renormalisation
- Physical pion mass results
- Results for full operator basis
- First result (ratios only) w/o effective action for the b -quark

RBC/UKQCD/JLQCD

- Huge data set (15 ensembles)
- Target all 5+5 bag parameters, decay constants, ratios
- Good HQ reach \Rightarrow no effective action prescription for b -quark
- χ -symmetry \Rightarrow continuum-like block-diagonal renormalisation
- Control all systematics (data)

CKM Status and Future

- $|V_{td}|$ and $|V_{ts}|$ known at $\approx 2.5\%$, $|V_{td}/V_{ts}|$ at $\approx 1.5\%$
- Uncertainty theory dominated - work is ongoing
- RBC/UKQCD/JLQCD: Aim to reduce uncertainties with **all DWF** approach (data generated using Grid and Hadrons [Grid, Hadrons])

ADDITIONAL SLIDES

Neutral $B_{(s)}$ meson mixing - theory

$$B_q^0 \xrightarrow{\mathcal{H}^{\Delta b=2}} \bar{B}_q^0$$

$$B_q^0 \xrightarrow{\mathcal{H}^{\Delta b=1}} P_n \xrightarrow{\mathcal{H}^{\Delta b=1}} \bar{B}_q^0$$

$$\langle B_q^0 | \mathcal{H}_{\text{eff}} | \bar{B}_q^0 \rangle \propto \underbrace{\langle B_q^0 | \mathcal{H}^{\Delta b=2} | \bar{B}_q^0 \rangle}_{\text{Short distance}} + \underbrace{\sum_n \frac{\langle B_q^0 | \mathcal{H}^{\Delta b=1} | n \rangle \langle n | \mathcal{H}^{\Delta b=1} | \bar{B}_q^0 \rangle}{E_n - M_{B_q}}}_{\text{Long distance}}$$

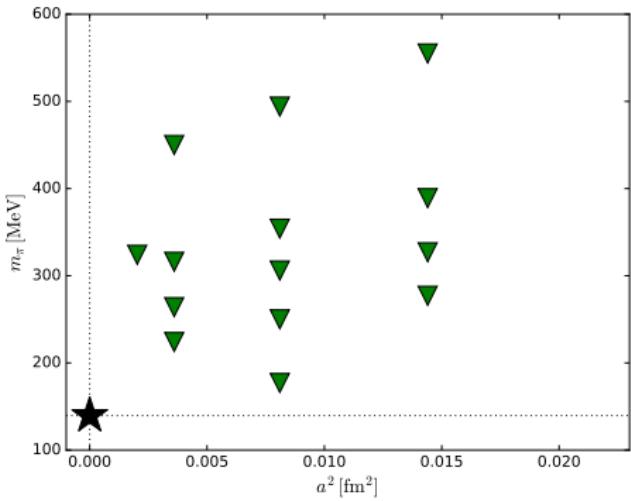
$$\text{short distance} \propto \left| \sum_{q'=u,c,t} \frac{m_{q'}^2}{M_W^2} V_{q'b} V_{q'q}^* \right|^2 \approx \frac{m_t^4}{M_W^4} |V_{tb} V_{tq}^*|^2$$

SD: Top enhanced: $m_t^2 V_{tb} V_{tq}^* \gg m_c^2 V_{cb} V_{cq}^* \gg m_u^2 V_{ub} V_{uq}^*$

LD: Only m_c, m_u in intermediate states: no top + CKM suppressed

⇒ Short distance dominated.

- $N_f = 2 + 1$ flavours of staggered quarks (asqtad) in the sea
- 4 lattice spacings, pion masses from 177 – 555 MeV
- valence light & strange: asqtad
- Fermilab method for the b -quark
- *mostly non-perturbative* 1-loop lattice perturbation theory
- Computation of $f_{B_q} \sqrt{\hat{B}_{B_q}}$ and ξ
- f_{B_q} taken from the PDG average to access to \hat{B}_{B_q}
- all 5 operators for B_d and B_s



$$f_{B_d} \sqrt{\hat{B}_{B_d}} = 227.7(9.5) \text{ MeV}$$

$$f_{B_s} \sqrt{\hat{B}_{B_s}} = 274.6(8.4) \text{ MeV}$$

$$\xi = 1.206(18)$$

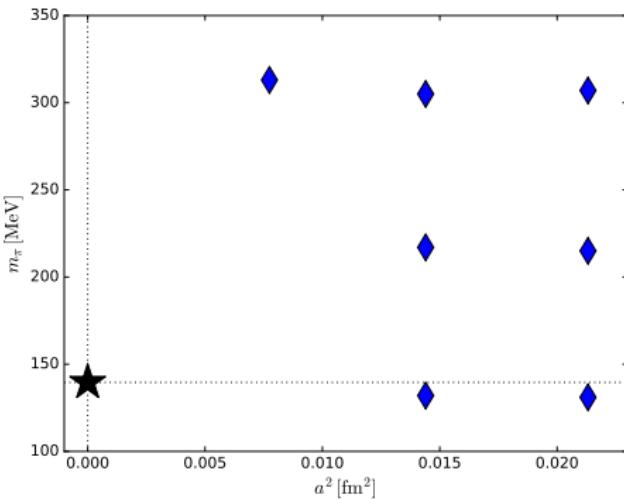
FNAL/MILC: error budget [%] [1602.03560]

	statistics	inputs	κ tuning	matching	chiral	LQ disc	HQ disc	fit total
$\langle \mathcal{O}_1^d \rangle$	4.2	0.4	2.1	3.2	2.3	0.6	4.6	7.7
$\langle \mathcal{O}_2^d \rangle$	4.6	0.3	1.1	3.7	2.6	0.6	4.6	8.0
$\langle \mathcal{O}_3^d \rangle$	8.7	0.2	2.1	12.6	4.8	1.2	9.9	19.0
$\langle \mathcal{O}_4^d \rangle$	3.7	0.4	1.7	2.2	1.9	0.5	3.9	6.4
$\langle \mathcal{O}_5^d \rangle$	4.7	0.5	2.5	4.7	2.7	0.8	4.9	9.1
$\langle \mathcal{O}_1^s \rangle$	2.9	0.4	1.5	2.1	1.6	0.4	3.2	5.4
$\langle \mathcal{O}_2^s \rangle$	3.1	0.3	0.8	2.5	1.6	0.4	3.1	5.5
$\langle \mathcal{O}_3^s \rangle$	5.9	0.3	1.4	8.6	3.0	0.7	6.9	13.0
$\langle \mathcal{O}_4^s \rangle$	2.7	0.4	1.2	1.6	1.3	0.3	2.9	4.8
$\langle \mathcal{O}_5^s \rangle$	3.4	0.4	1.8	3.4	1.9	0.5	3.6	6.7
ξ	0.8	0.4	0.3	0.5	0.4	0.1	0.7	1.4

Uncertainty dominated by chiral-continuum limit fit, in particular

- statistical
- heavy quark discretisation errors
- matching

- $N_f = 2 + 1 + 1$ flavours of staggered quarks (HISQ) in sea
- light quarks using HISQ in the valence
- 3 lattice spacings, 2 physical pion mass ensembles
- improved nonrelativistic QCD action for the b
- all 5 operators for B_d and B_s
- blinded analysis
- Computation of the $\hat{B}_{B_q}^{(i)}$
- ξ and $f_{B_q} \sqrt{\hat{B}_{B_q}}$ accessed by using decay constants taken from a different computation



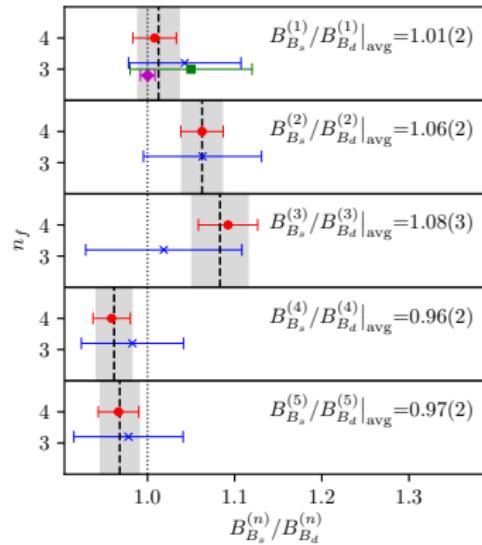
$$\hat{B}_{B_d}^{(1)} = 1.222(61)$$

$$\hat{B}_{B_s}^{(1)} = 1.232(53)$$

$$\hat{B}_{B_s}^{(1)} / \hat{B}_{B_d}^{(1)} = 1.008(25)$$

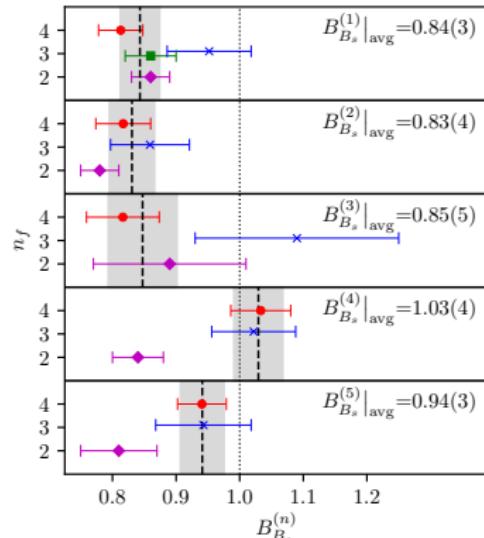
HPQCD: error budget and results [1907.01025]

HPQCD19, FNAL/MILC16, HPQCD09, RBC/UKQCD18



Uncertainty dominated by matching terms α_s^2 and $\alpha_s \Lambda_{QCD}/m_b$.

HPQCD19, FNAL/MILC16, HPQCD09, ETM13



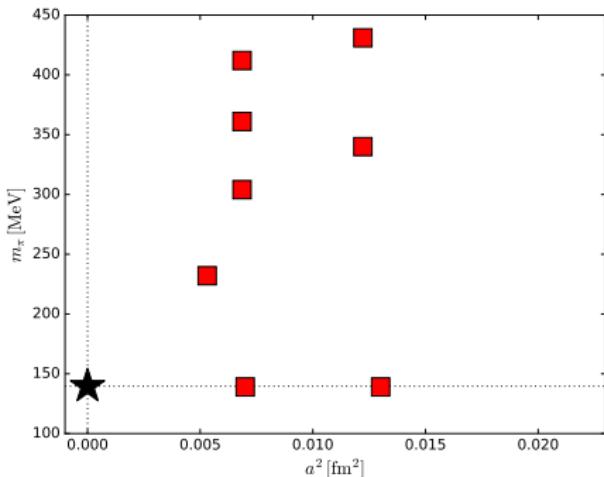
	$B_{B_s}^{(1)}$	$B_{B_s}^{(2)}$	$B_{B_s}^{(3)}$	$B_{B_s}^{(4)}$	$B_{B_s}^{(5)}$	$B_{B_s}^{(1)}/B_{B_d}^{(1)}$
lattice data	1.4	1.4	1.5	1.6	1.5	1.5
η_i^q	0.0	2.3	2.3	2.1	1.2	0.0
α_s^2 terms	2.1	2.9	5.2	1.9	1.5	0.1
$\alpha_s \Lambda_{QCD}/m_b$ terms	2.9	2.8	2.9	2.8	2.7	0.0
$(\alpha \Lambda_{QCD})^{2n}$ terms	1.8	1.9	2.3	1.5	1.8	0.1
m_f extrapolation	0.4	0.4	0.7	0.5	0.4	1.9
Total	4.3	5.3	7.0	4.6	4.1	2.5

Light and strange

- Unitary light quark mass
- Physical strange quark mass
- DWF parameters same between sea and valence

Heavy: charm and b(eyond)

- Möbius DWF
- $M_5 = 1.0$, $L_s = 12$
- Stout smeared (3 hits, $\rho = 0.1$)
- Many source-sink separations
- Range of heavy quark masses
 $m_c \lesssim m_h \lesssim m_b/2$



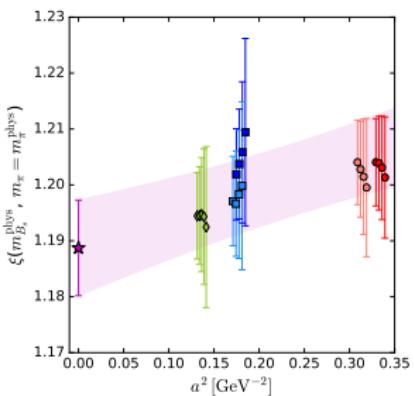
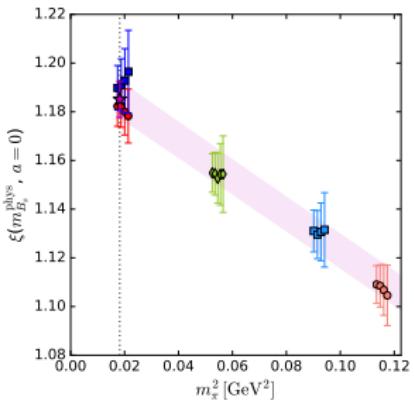
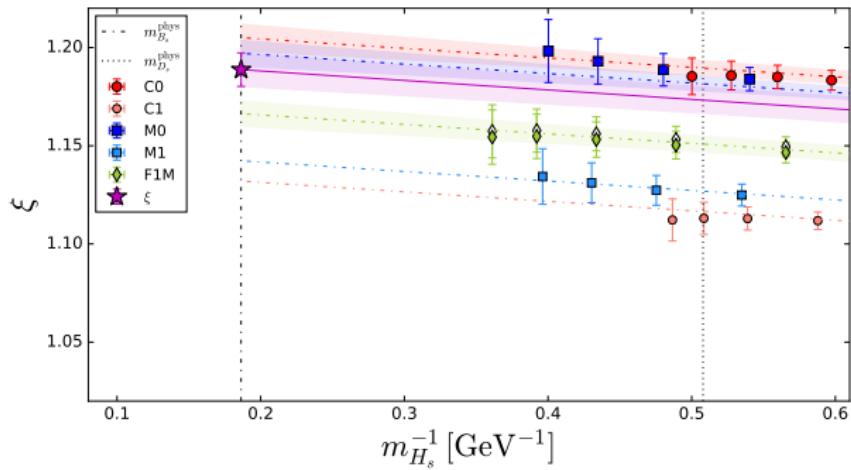
Ensembles

- $N_f = 2 + 1$ flavours, 3 lattice spacings, 2 physical pion mass ensembles

⇒ All Domain Wall Fermion mixed action set-up

RBC/UKQCD: strategy [1812.08791]

- Computation of $SU(3)$ breaking ratios f_{D_s}/f_D , f_{B_s}/f_B , B_{B_s}/B_{B_d} and ξ
- Renormalisation constants cancel
- Benign behaviour with $1/m_H$
- Simultaneous fit to m_π^2 , a^2 , $\frac{1}{m_H}$



RBC/UKQCD: error budget [%] [1812.08791]

	f_{D_s}/f_D	f_{B_s}/f_B	ξ	B_{B_s}/B_{B_d}
central	1.1740	1.1949	1.1939	0.9984
stat	0.43%	0.50%	0.56%	0.45%
fit chiral-CL	+0.31% -0.32%	+0.34% -0.54%	+0.38% -0.45%	+0.42% -0.01%
fit heavy mass	+0.07% -0.05%	+0.00% -0.82%	+0.00% -0.87%	+0.27% -0.22%
H.O. heavy	0.00%	0.47%	0.35%	0.21%
H.O. disc.	0.01%	0.01%	0.12%	0.17%
$m_u \neq m_d$	0.08%	0.07%	0.08%	0.01%
finite size	0.18%	0.18%	0.18%	0.18%
total systematic	+0.38% -0.38%	+0.61% -1.38%	+0.56% -1.37%	+0.66% -0.45%
total sys+stat	+0.58% -0.58%	+0.79% -1.47%	+0.80% -1.48%	+0.80% -0.63%

Uncertainty dominated by

- chiral-continuum fit
- heavy quark extrapolation
- estimates of higher order $1/m_H$ terms

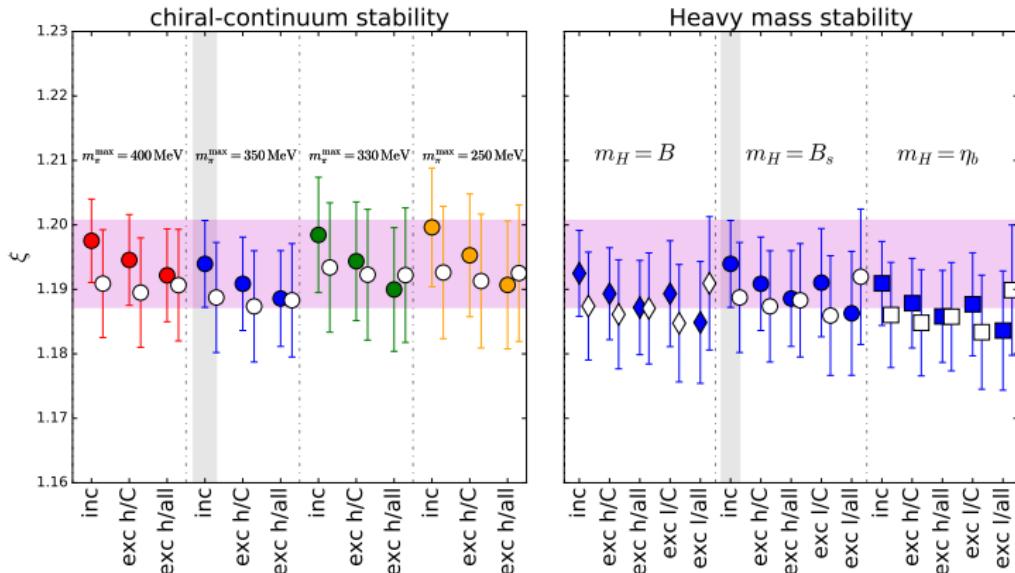
$$f_{D_s}/f_D = 1.1740(51)_{\text{stat}} \left({}^{+0.68}_{-0.68} \right)_{\text{sys}}$$

$$f_{B_s}/f_B = 1.1949(60)_{\text{stat}} \left({}^{+0.95}_{-1.75} \right)_{\text{sys}}$$

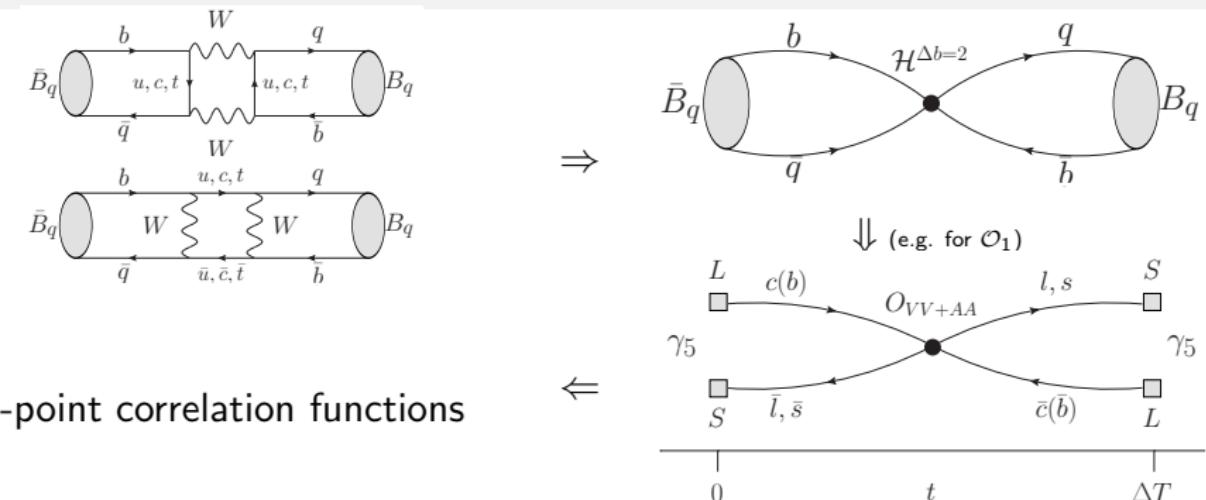
$$B_{B_s}/B_{B_d} = 0.9984(45)_{\text{stat}} \left({}^{+0.80}_{-0.63} \right)_{\text{sys}}$$

$$\xi = 1.1939(67)_{\text{stat}} \left({}^{+0.95}_{-1.77} \right)_{\text{sys}}$$

- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



Lattice strategy to simulate meson mixing



$$C_3^{ab,\mathcal{O}}(t; \Delta T) \sim \sum_i \sum_j \frac{Z_a^i Z_b^j \langle P_{(j)} | \mathcal{O} | \bar{P}_{(i)} \rangle}{4E_{(i)} E_{(j)}} e^{-E_{(j)} t} e^{-E_{(i)}(\Delta T - t)}$$

ΔT separation between source and sink

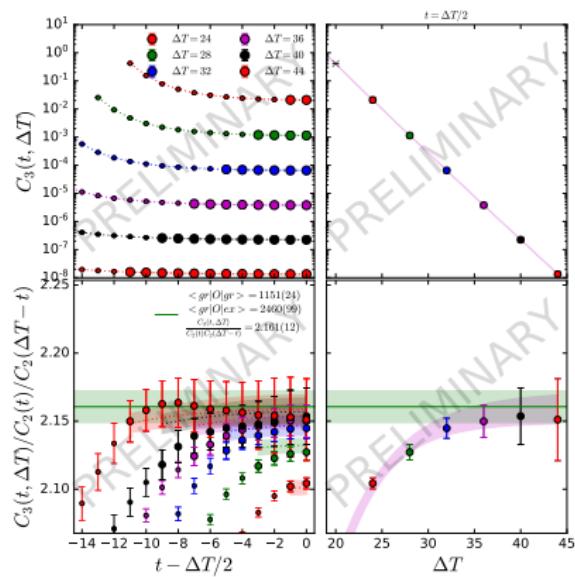
t position of operator \mathcal{O} between source and sink

From correlation functions to matrix elements

$$C_3^{ab,\mathcal{O}}(t; \Delta T) \sim \sum_i \sum_j \frac{Z_a^i Z_b^j \langle P_{(j)} | \mathcal{O} | \bar{P}_{(i)} \rangle}{4E_{(i)} E_{(j)}} e^{-E_{(j)} t} e^{-E_{(i)}(\Delta T - t)}$$

- Simultaneous fits of 2-point functions and 3-point functions for multiple ΔT 's to extract $\langle P_{\text{gr}} | \mathcal{O}_{(i)} | \bar{P}_{\text{gr}} \rangle$ with highest possible confidence and precision.
- Two separate independent analyses.
- Finalising investigation of stability of fits.

expl: heavy-light \mathcal{O}_2 for $a^{-1} \sim 3.5 \text{ GeV}$ →



Non-perturbative renormalisation on the lattice [hep-lat/9411010]

We want

$$\text{Amplitude} = C^{\overline{\text{MS}}}(\mu) \langle \mathcal{O} \rangle^{\overline{\text{MS}}}(\mu)$$

- Wilson coefficients (typically) computed in $\overline{\text{MS}}$ at some scale μ .
- Operators $\langle \mathcal{O} \rangle^{\text{bare}}(a)$ computed with lattice regulator a^{-1} .
- Renormalise $\langle \mathcal{O} \rangle^{\text{bare}}(a)$ at scale μ in regularisation independent (RI) scheme, by computing a non-pert. renormalisation factor $Z^{\text{RI}}(\mu, a)$.

$$\langle \mathcal{O} \rangle^{\text{RI}}(\mu) = \lim_{a^2 \rightarrow 0} Z^{\text{RI}}(\mu, a) \langle \mathcal{O} \rangle^{\text{bare}}(a)$$

- Match to preferred scheme (e.g. $\overline{\text{MS}}$) using P.T. at μ : $R^{\overline{\text{MS}} \leftarrow \text{RI}}(\mu)$
- If the operators mix: C and $\langle \mathcal{O} \rangle$ become vectors, R and Z matrices.

$$\text{Amplitude} = C_i^{\overline{\text{MS}}}(\mu) R_{ij}^{\overline{\text{MS}} \leftarrow \text{RI}}(\mu) \lim_{a \rightarrow 0} Z_{jk}^{\text{RI}}(\mu, a) \langle \mathcal{O}_k \rangle^{\text{bare}}(a)$$

- Chirally symmetric fermions $\Rightarrow R$ and Z are block diagonal.

NPR for Neutral Meson Mixing operators

- Analogous to the neutral kaon mixing case [1708.03552, 1812.04981]
- 5 operators $\mathcal{O}_1 = \mathcal{O}_{VV+AA}$, $\mathcal{O}_2 = \mathcal{O}_{VV-AA}$, $\mathcal{O}_3 = \mathcal{O}_{SS-PP}$,
 $\mathcal{O}_4 = \mathcal{O}_{SS+PP}$, $\mathcal{O}_5 = \mathcal{O}_{TT}$
- Block diagonal renormalisation matrix (up to $O(am_{\text{res}})$)

$$Z_{ij}^{RI}(\mu, a) = \begin{pmatrix} Z_{11} & 0 & 0 & 0 & 0 \\ 0 & Z_{22} & Z_{23} & 0 & 0 \\ 0 & Z_{32} & Z_{33} & 0 & 0 \\ 0 & 0 & 0 & Z_{44} & Z_{45} \\ 0 & 0 & 0 & Z_{54} & Z_{55} \end{pmatrix}$$

- Generalise idea from 1701.02644 for fully non-perturbative mixed action renormalisation to four quark operators

$$\frac{\mathcal{P}[\Lambda_A](II)\mathcal{P}[\Lambda_A](hh)}{(\mathcal{P}[\Lambda_A](Ih))^2} = \frac{(Z_A^{\text{lh}})^2}{Z_A^{II}Z_A^{hh}}$$

- Data production of the mixed action NPR in progress.