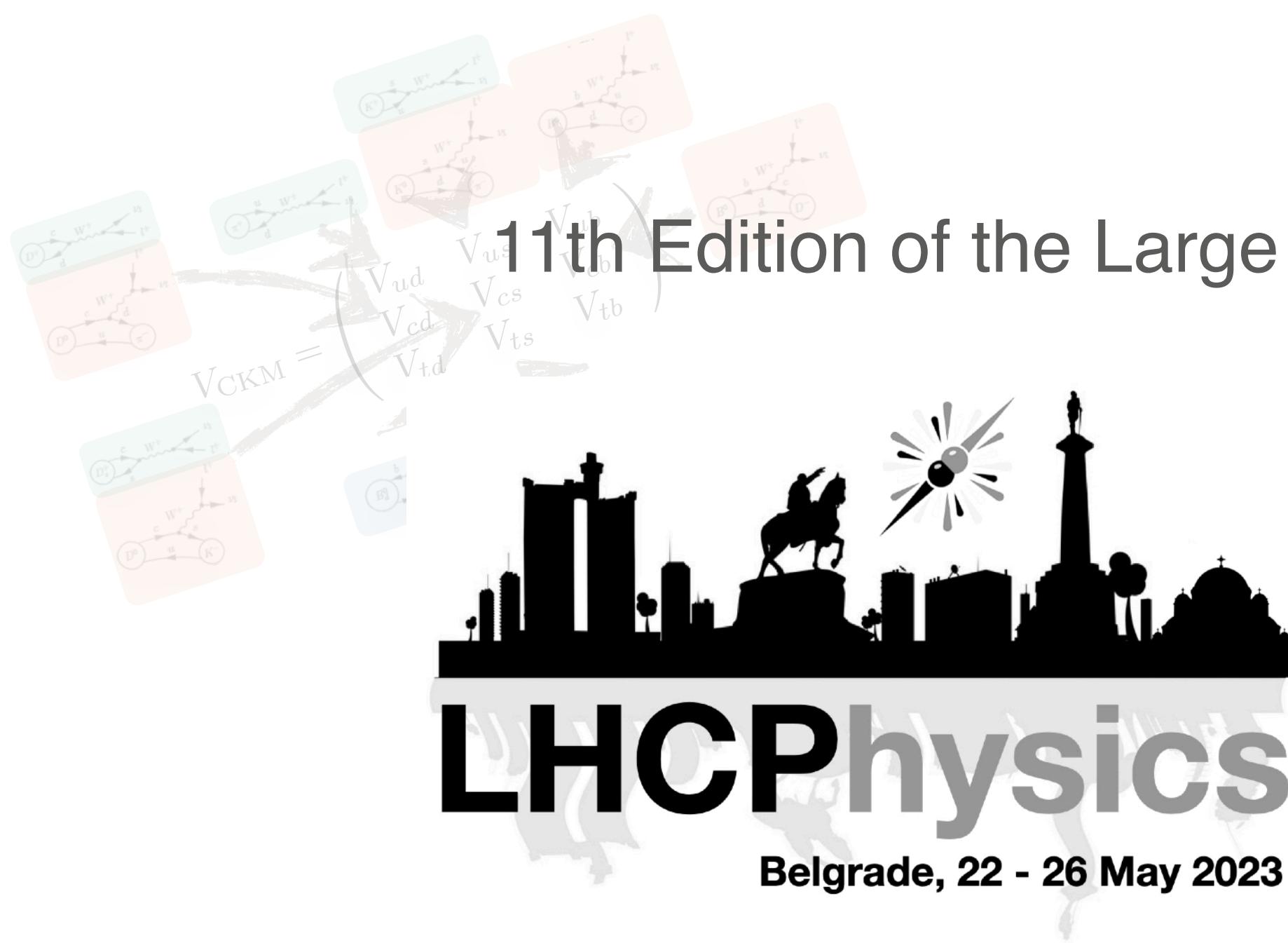


# B-semileptonic form-factors on the lattice



11th Edition of the Large Hadron Collider Physics Conference



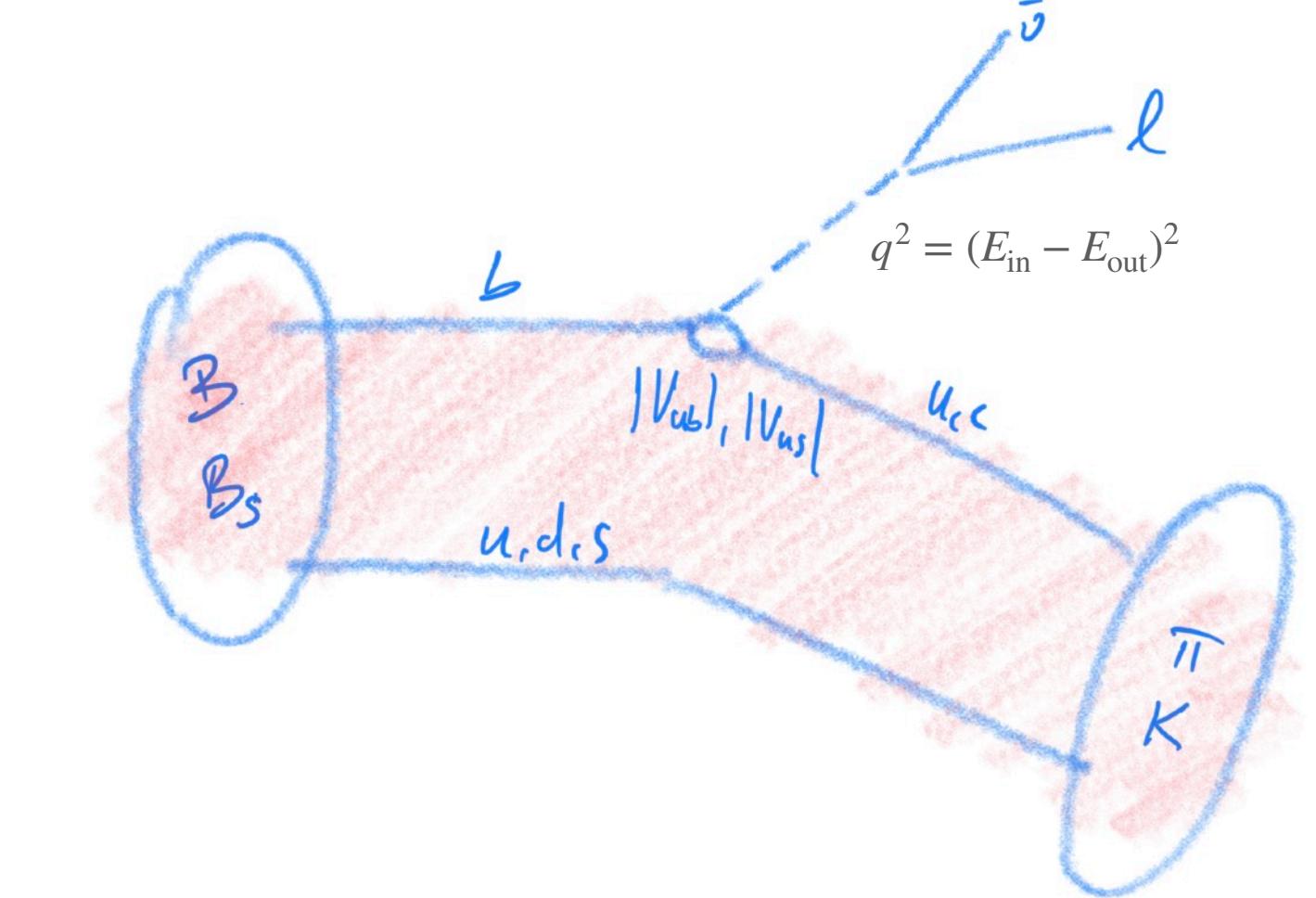
Andreas Jüttner



# three points

- exclusive decay form-factor calculations  $b \rightarrow u, b \rightarrow c$
- model- and truncation independent form-factor fitting
- inclusive decay rate calculations  $b \rightarrow c$

# Exclusive semileptonic meson decay



$$\frac{d\Gamma(B_s \rightarrow P \ell \nu)}{dq^2} = \eta_{\text{EW}} \frac{G_F^2 |V_{xb}|^2}{24\pi^3} \frac{(q^2 - m_\ell^2)^2 |\vec{k}|}{(q^2)^2} \left[ \left(1 + \frac{m_\ell^2}{2q^2}\right) \vec{k}^2 |f_+(q^2)|^2 + \frac{3m_\ell^2}{8q^2} \frac{(M_{B_s}^2 - M_P^2)^2}{M_{B_s}^2} |f_0(q^2)|^2 \right]$$

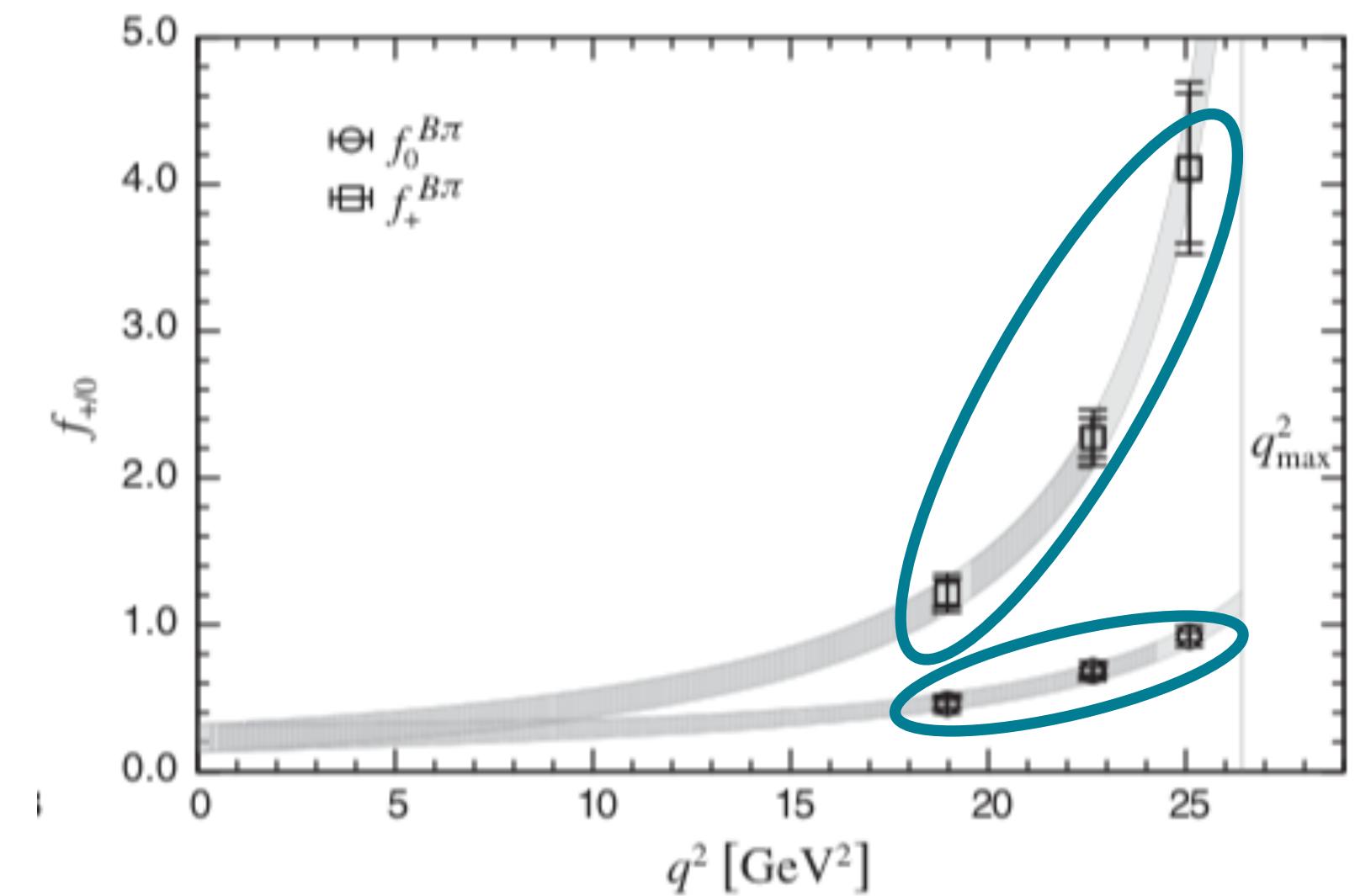
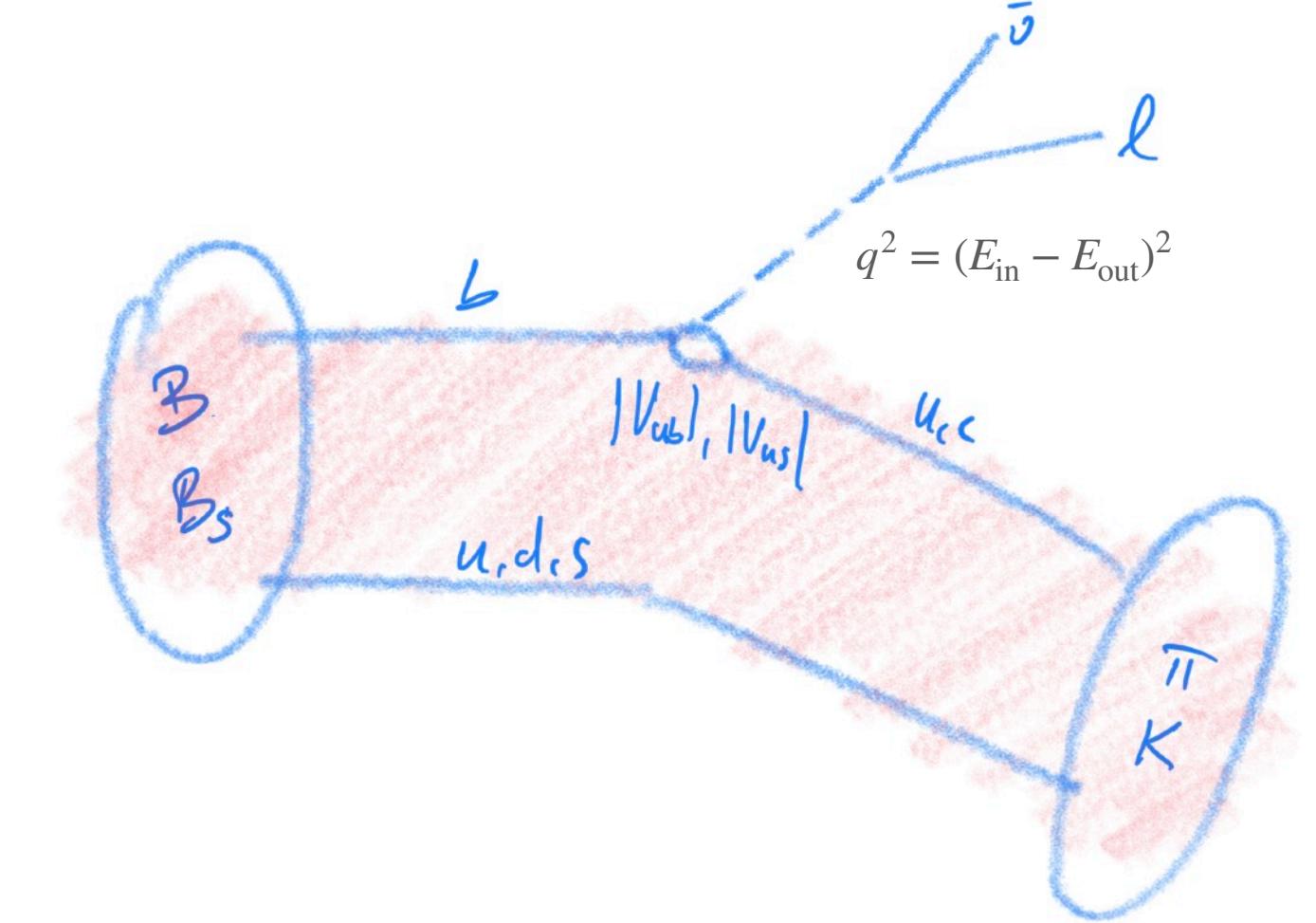
- form factors computed on lattice QCD at high precision (see [FLAG21](#))
- considered standard for tree-level decays
- heavily used for CKMology but also for lepton-flavour-universality tests
- only few collaborations competing
  
- estimating systematics in the last steps of the lattice-analysis can be rather challenging
- we (and others) have found instabilities in the extrapolation of lattice-form factors

$\langle \underline{\text{had}}' | H_w | \underline{\text{had}} \rangle_{\text{QCD}}$

# Exclusive semileptonic meson decay

Some novelties on the lattice:

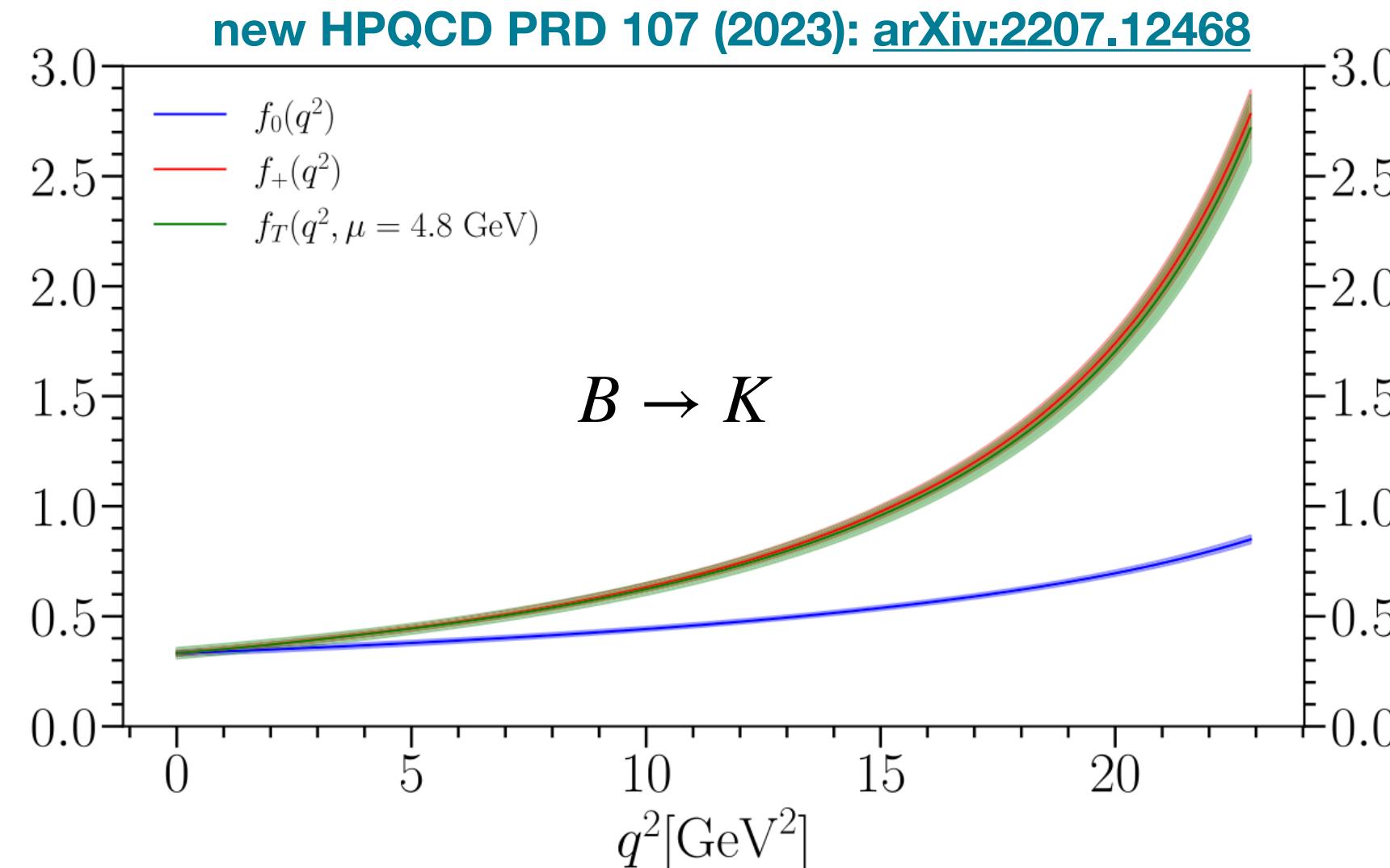
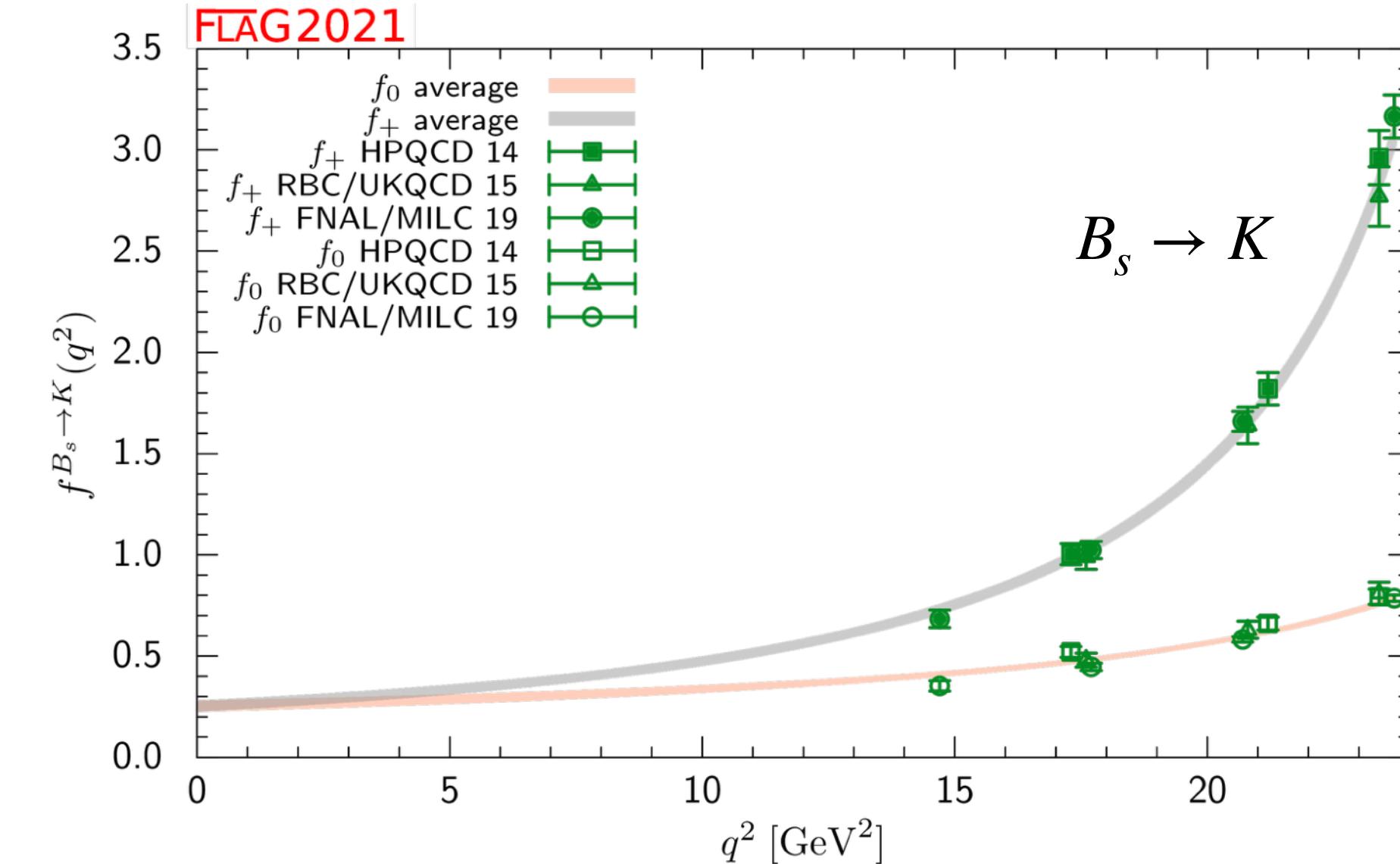
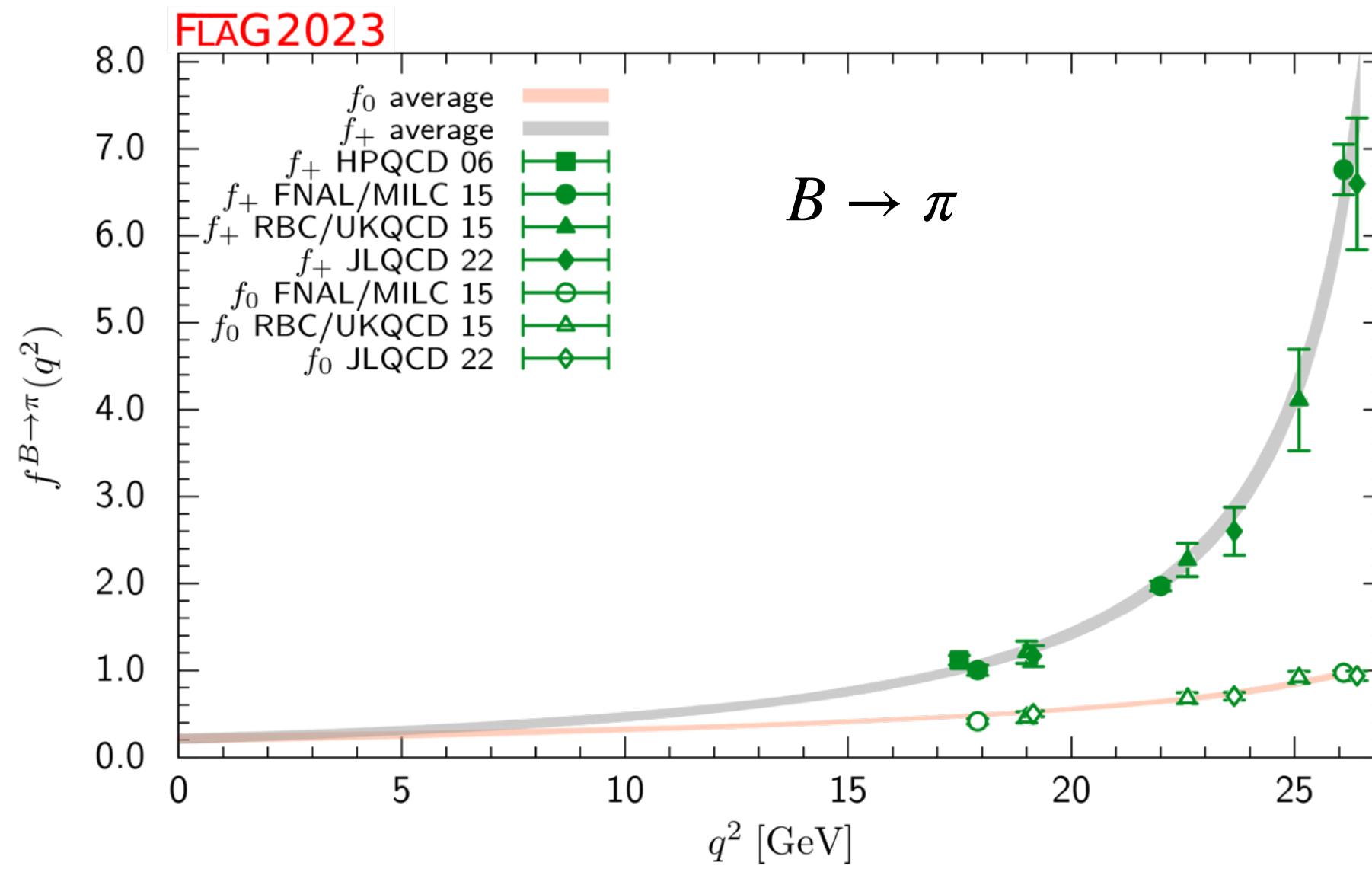
- **heavy-quark mass vs. lattice spacing (UV cutoff)**
  - so far use of effective field theory for heavy quark
  - work on fully relativistic lattice  $b$ -quark going on
  - challenge: control systematics due to discretisation
- **limited kinematic reach**
  - new finer lattice spacings allow to reach to lower  $q^2$
  - novel ideas for model-independent parameterisations of form factors



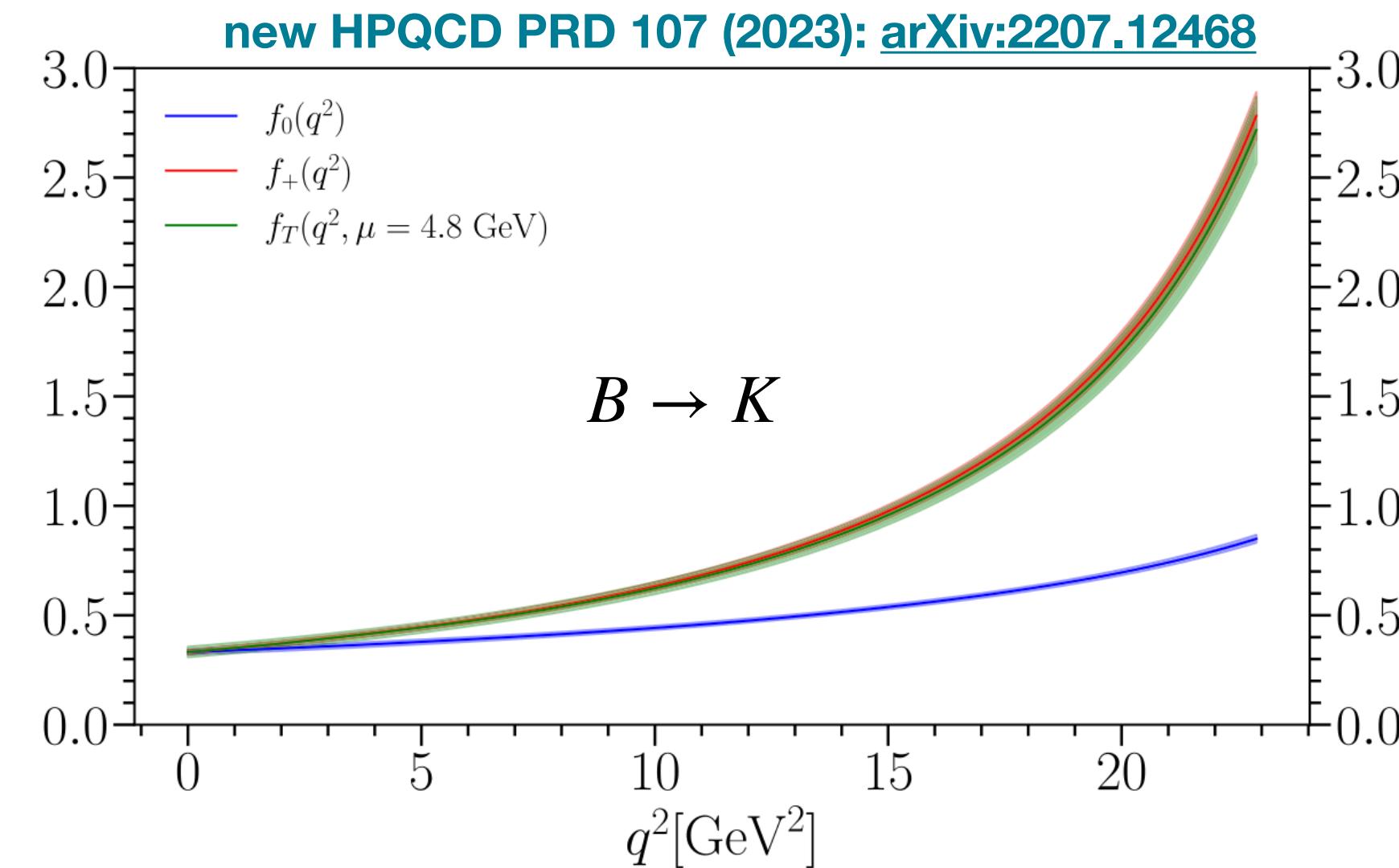
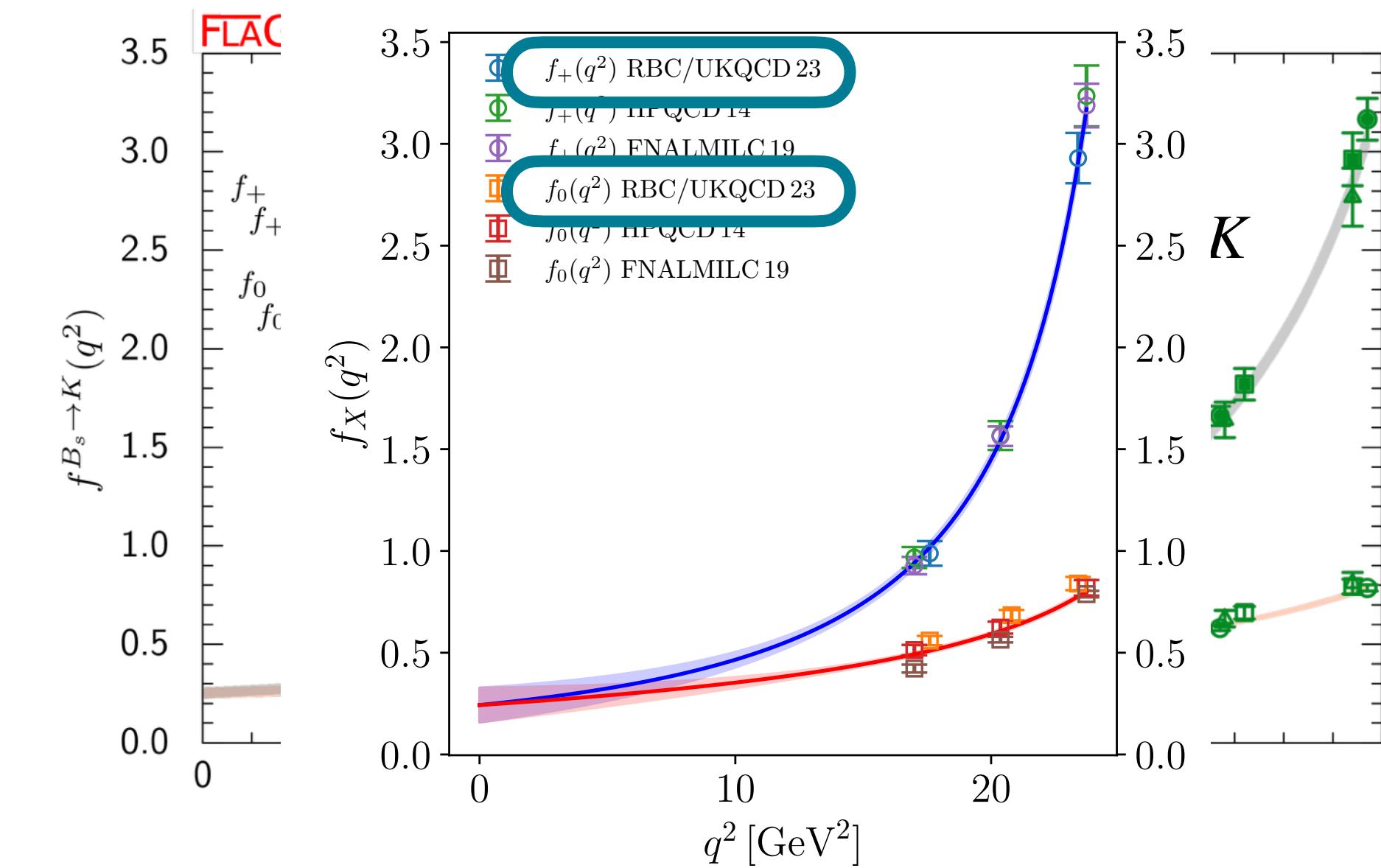
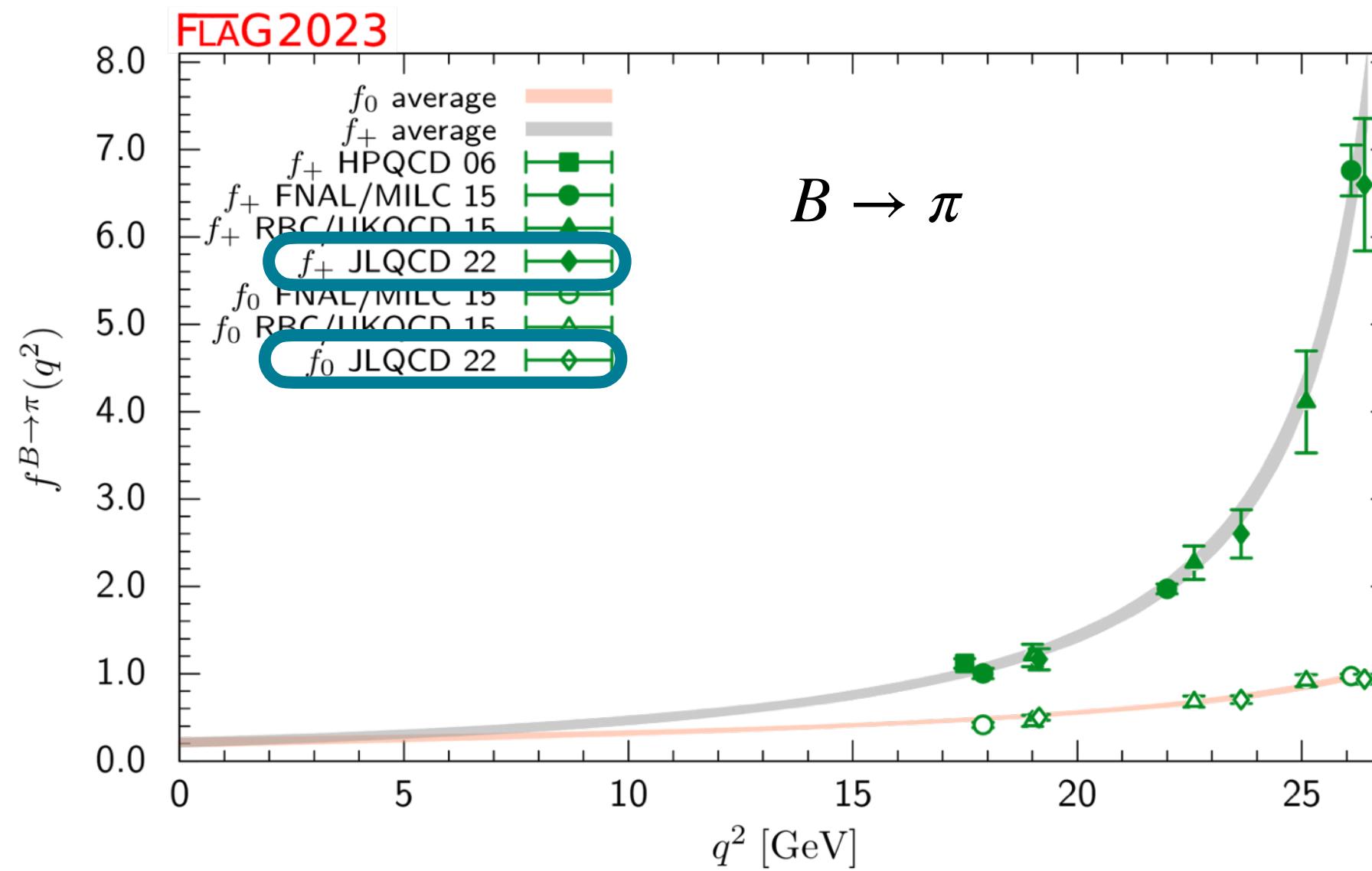
RBC/UKQCD PRD 91, 074510 (2015)2018

**exclusive  $b \rightarrow u$**

# $b \rightarrow u$ exclusive



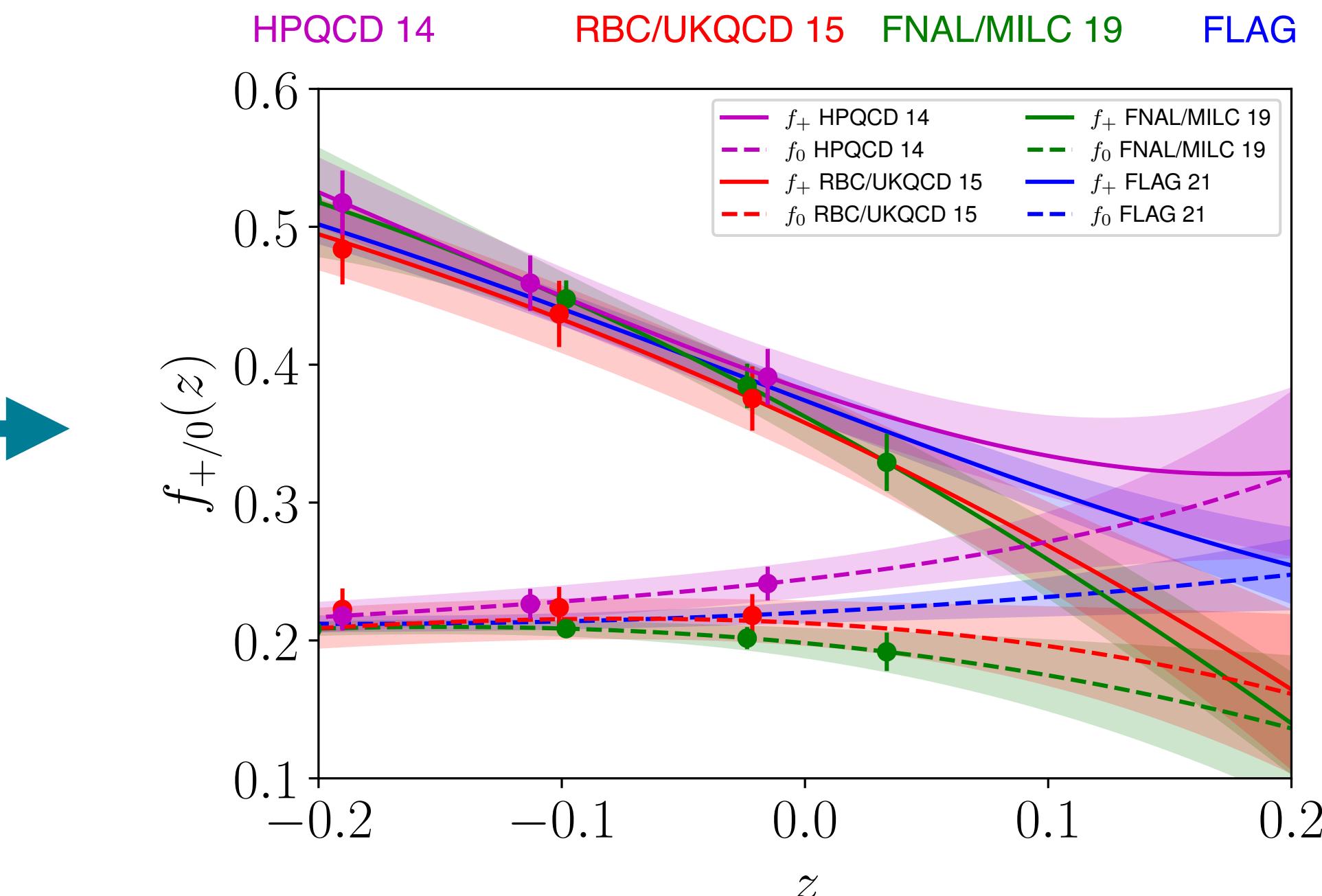
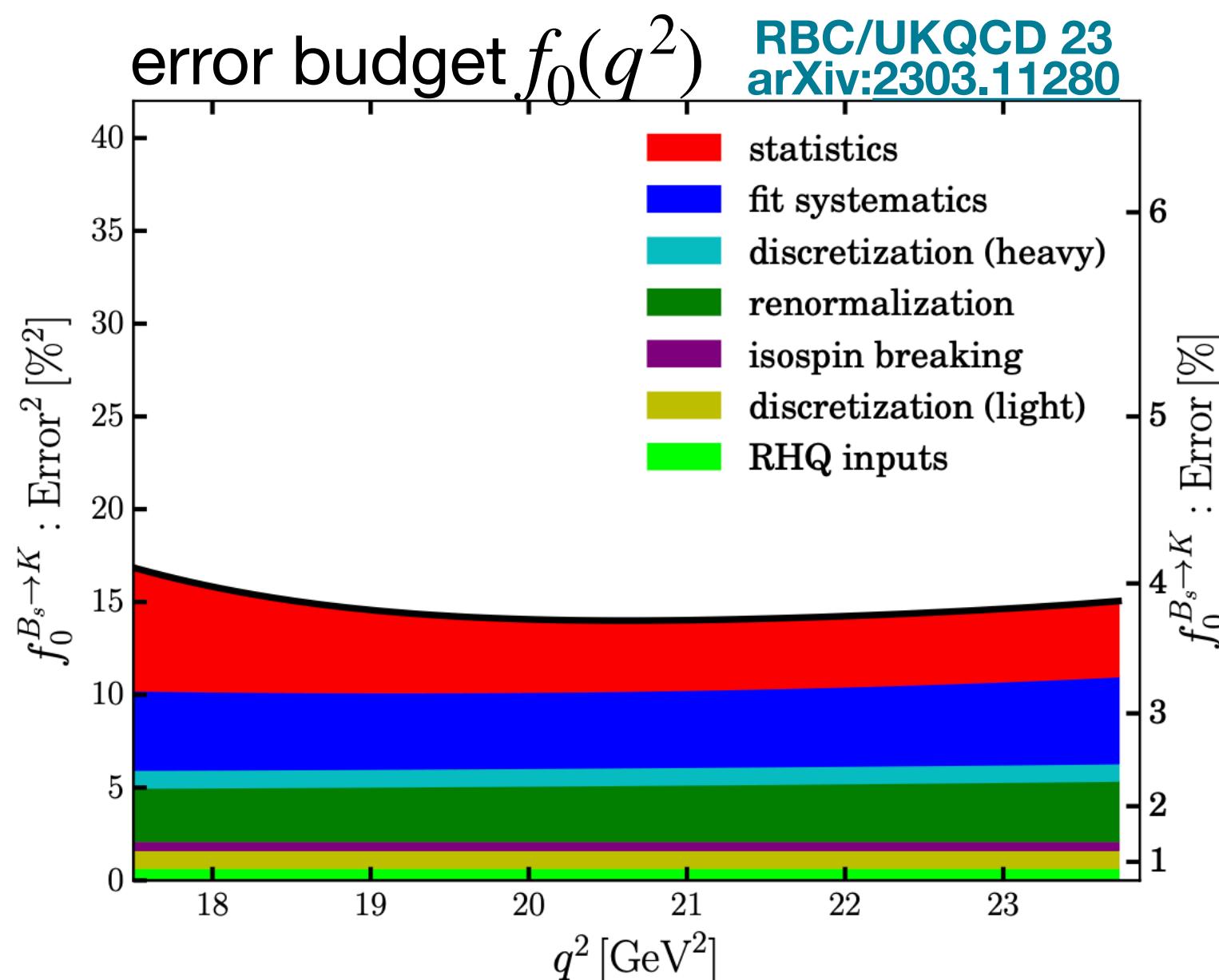
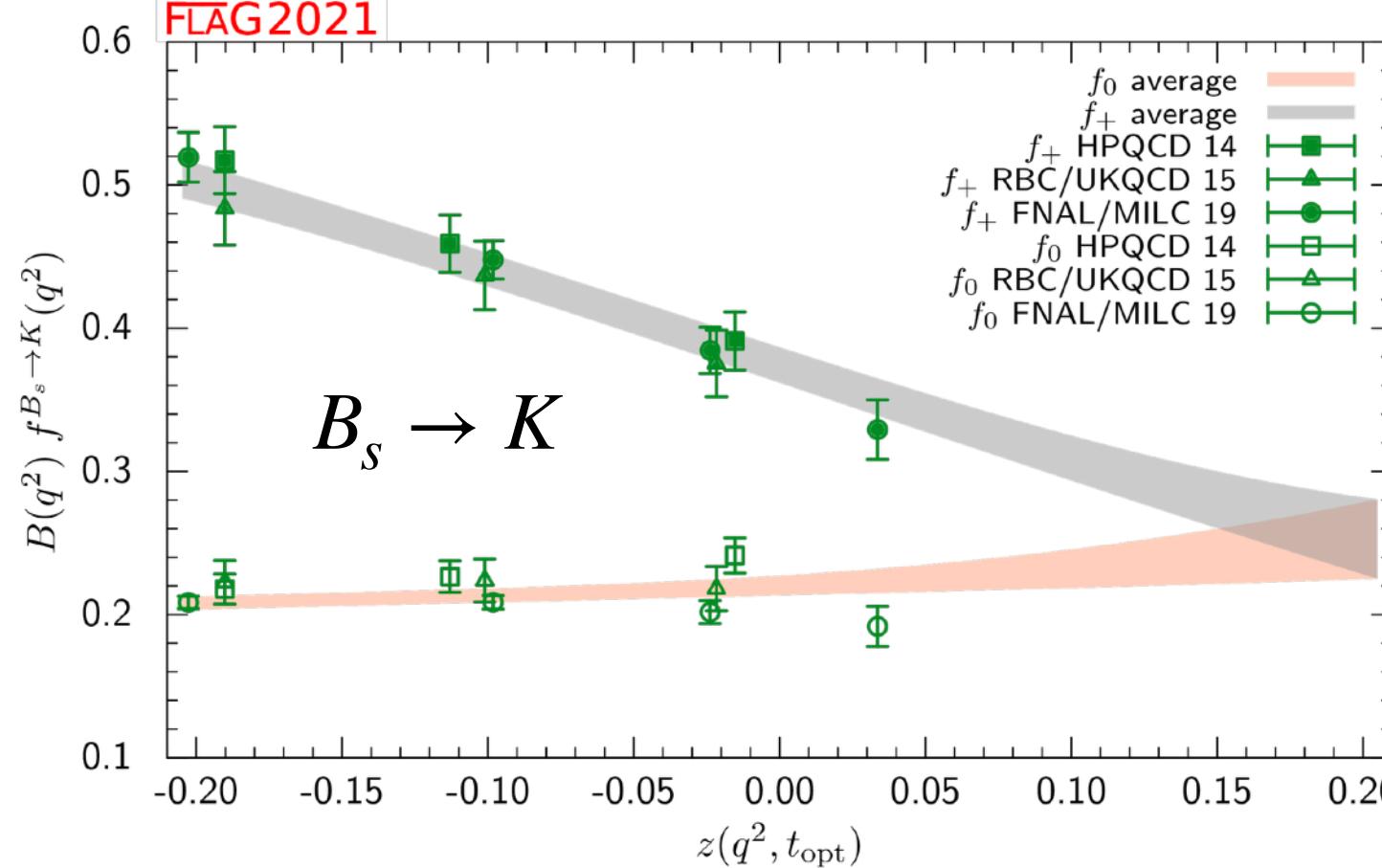
# $b \rightarrow u$ exclusive



New:

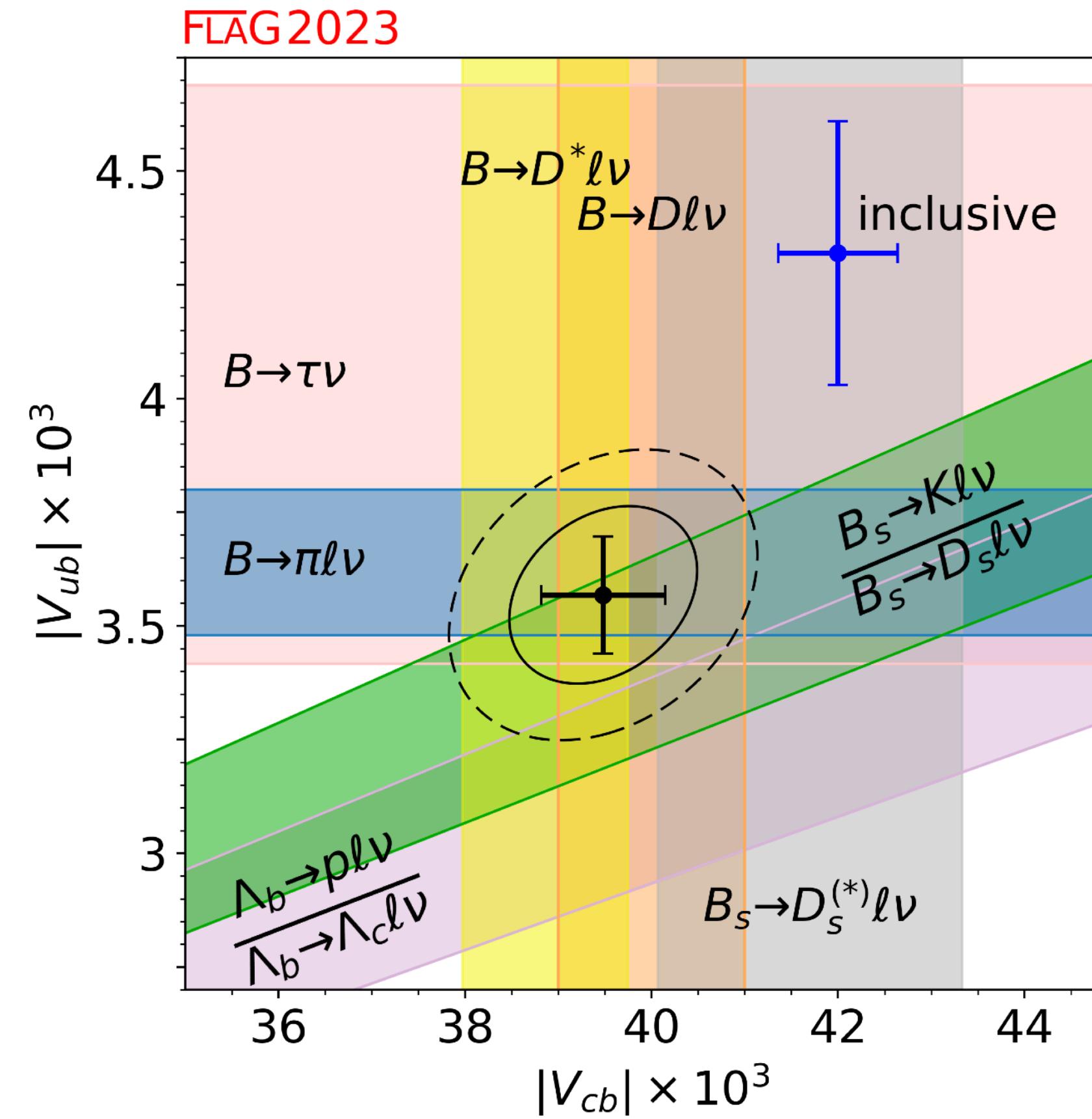
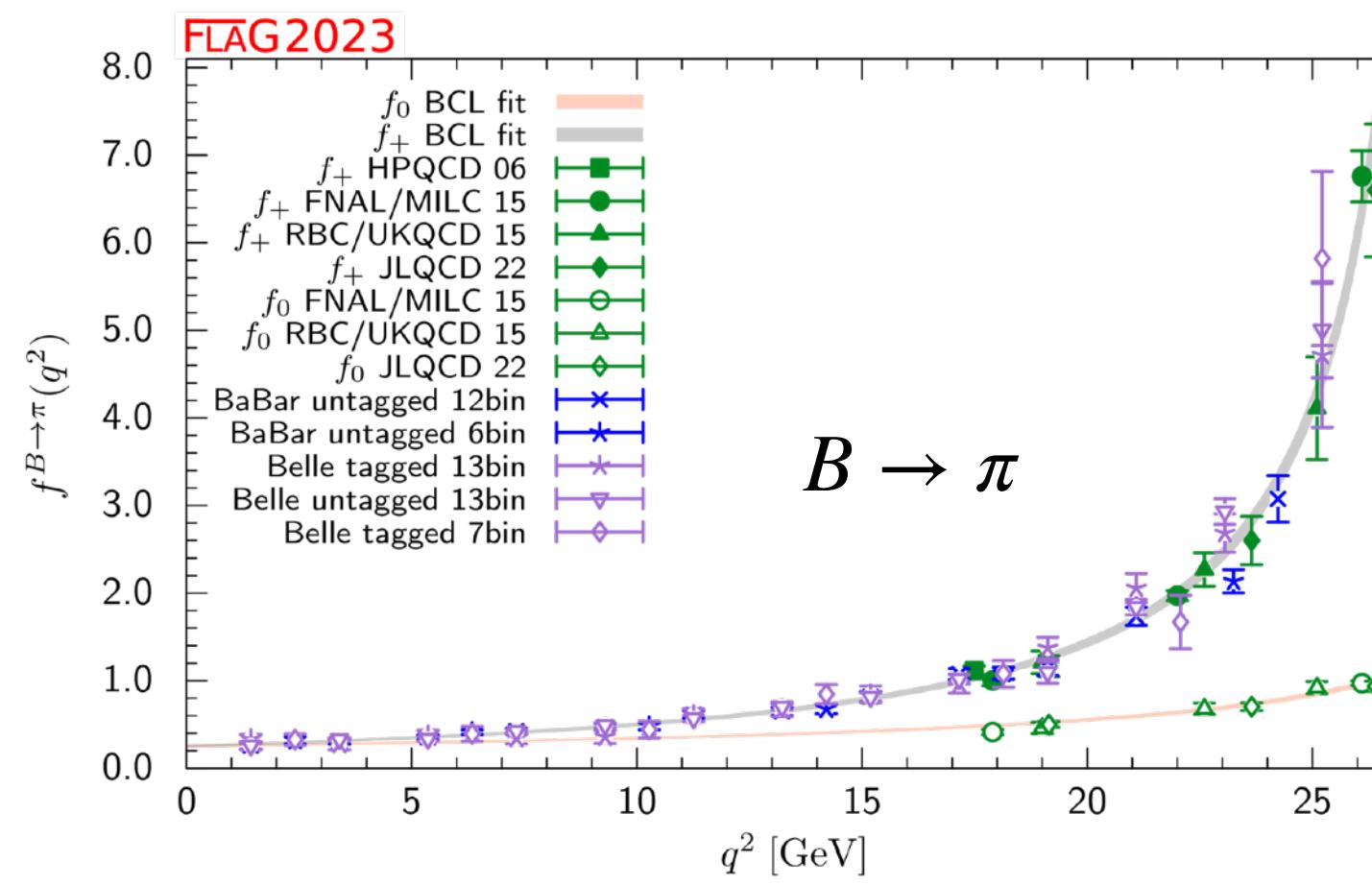
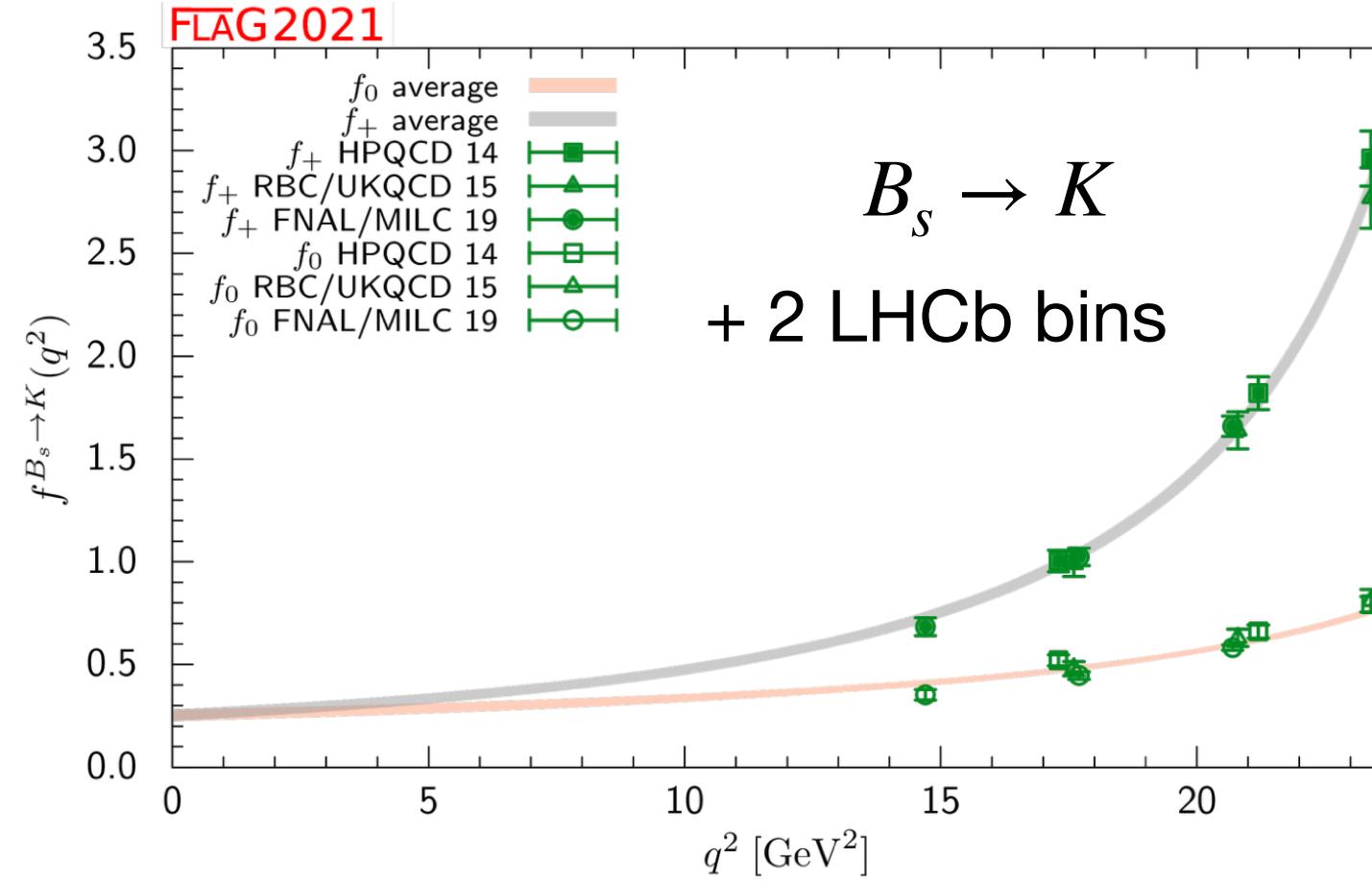
- $B \rightarrow \pi$  data by JLQCD 22 PRD [arXiv:2203.04938](#)
- $B_s \rightarrow K$  data by RBC/UKQCD 23 [arXiv:2303.11280](#)
- $B \rightarrow K$  data by HPQCD 23 PRD, [arXiv:2207.12468](#)

# $b \rightarrow u$ exclusive: work to do



- Lattice data sets show tension  $B \rightarrow \pi$ ,  $B_s \rightarrow K$  (combination requires PDG-inflation factor)
- by definition they should agree
- reasons yet to be understood (excited states Bär et al. arXiv:2210.06863, 2210.06857, chiral/cont extrapolation RBC/UKQCD arXiv:2303.11280, ...)

# $b \rightarrow u$ exclusive: $|V_{ub}|$



- $B \rightarrow \pi$  combined fit with experiment requires PDG inflation
- small tension inclusive vs exclusive

**exclusive**  $b \rightarrow c$

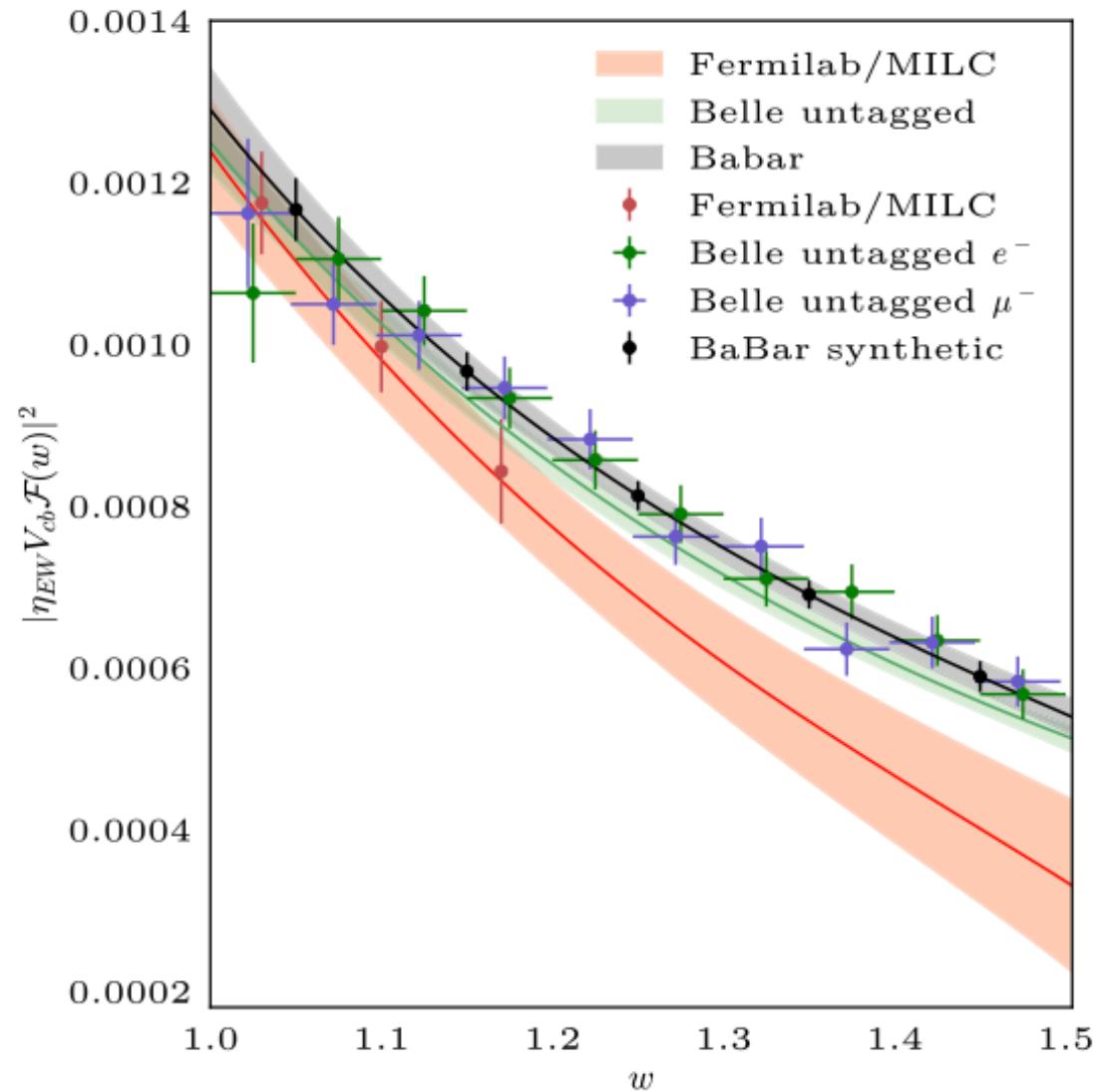
# A Puzzle in Flavour Physics?

## SL meson decay $B \rightarrow D^* \ell \nu$

### Novelties:

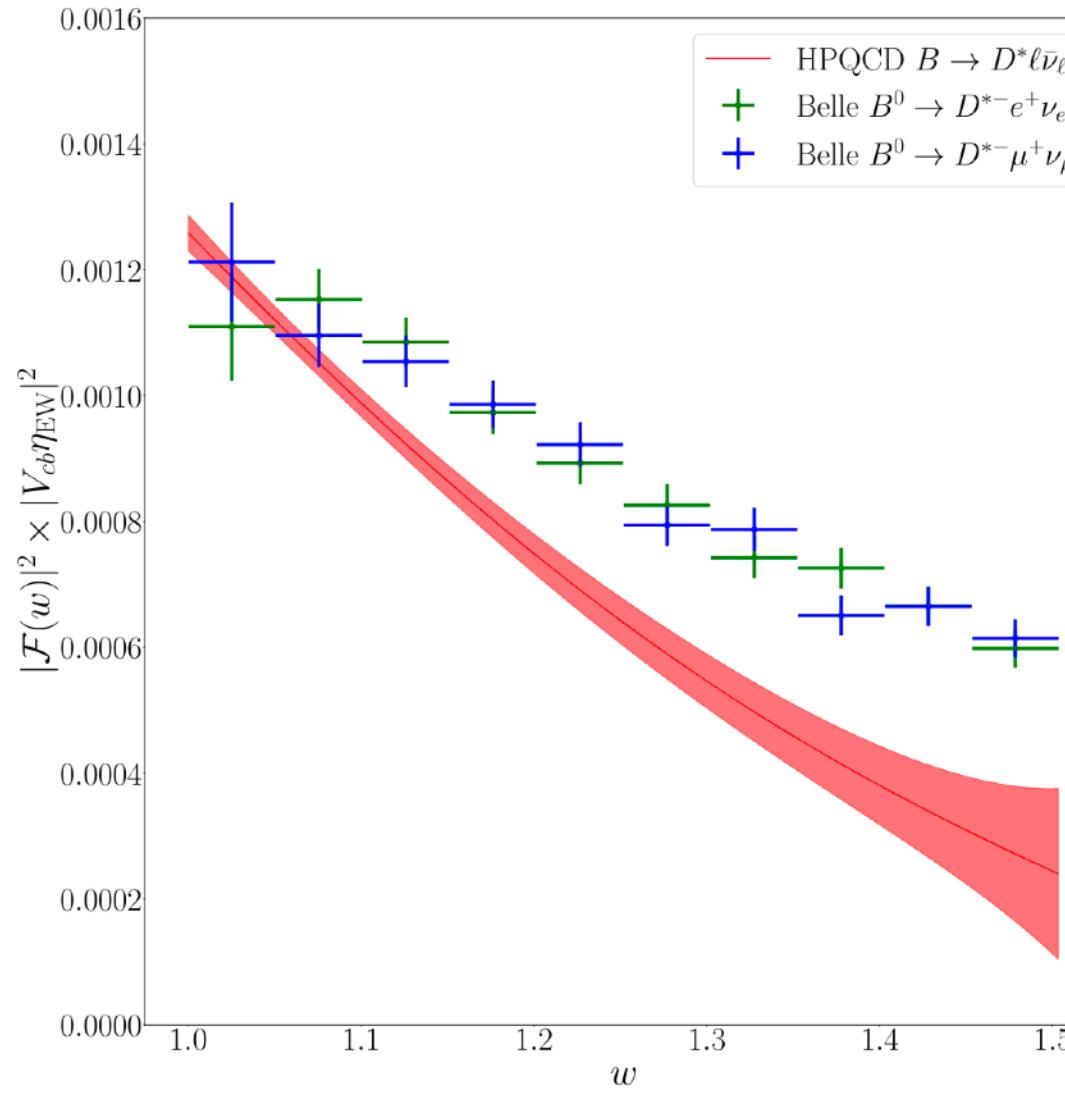
- form-factor shapes from the lattice FNAL/MILC, HPQCD
- JLQCD preliminary

FNAL/MILC EPJC (2022) [arXiv:2105.14019](https://arxiv.org/abs/2105.14019)



$$|V_{cb}| = (38.40 \pm 0.78) \times 10^{-3}$$
$$(\chi^2/N_{\text{dof}}/p) = (128, 84, 0.001)$$

HPQCD [arXiv:2304.03137](https://arxiv.org/abs/2304.03137)



$$|V_{cb}| = 39.31 \pm (74)$$

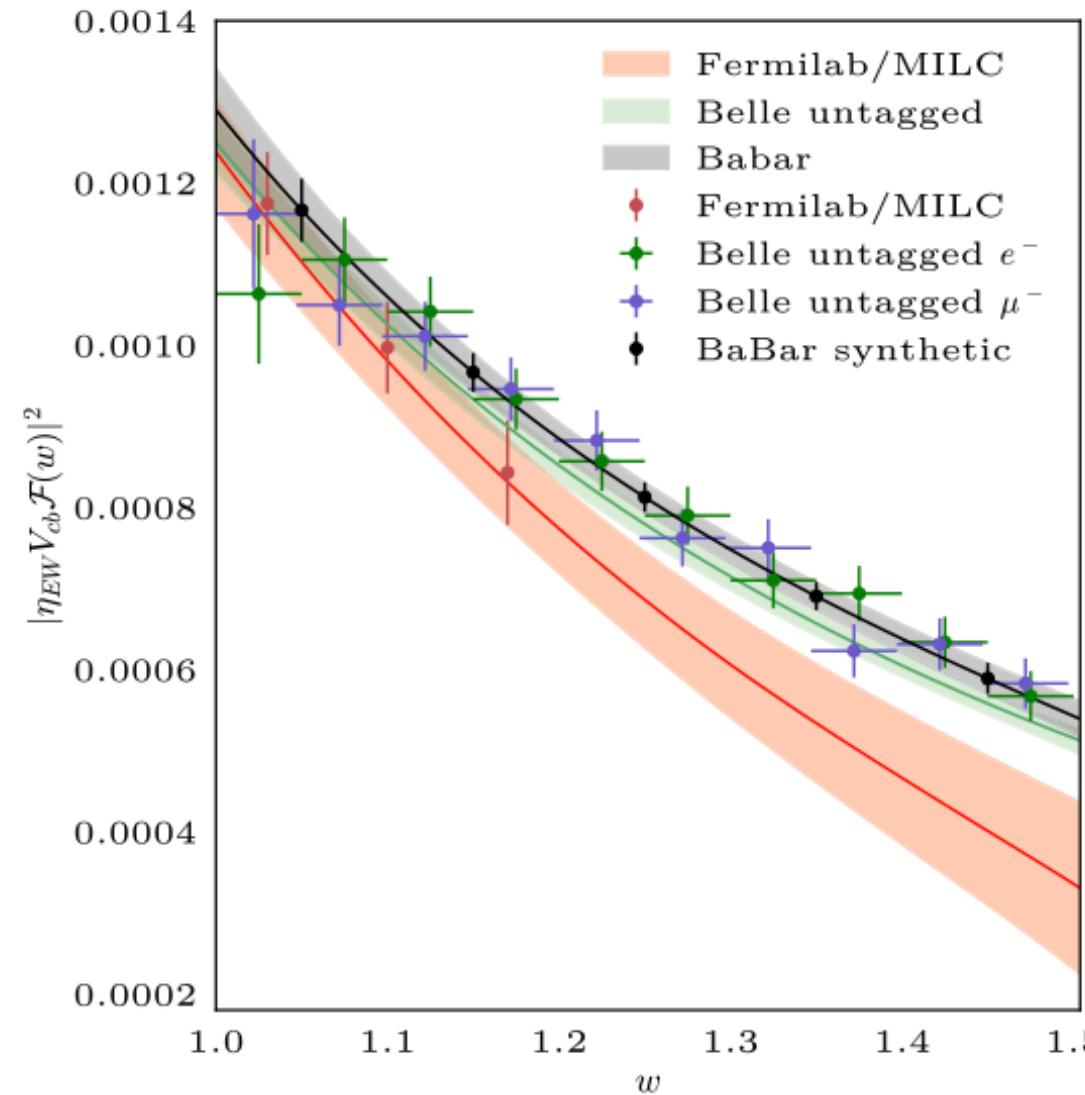
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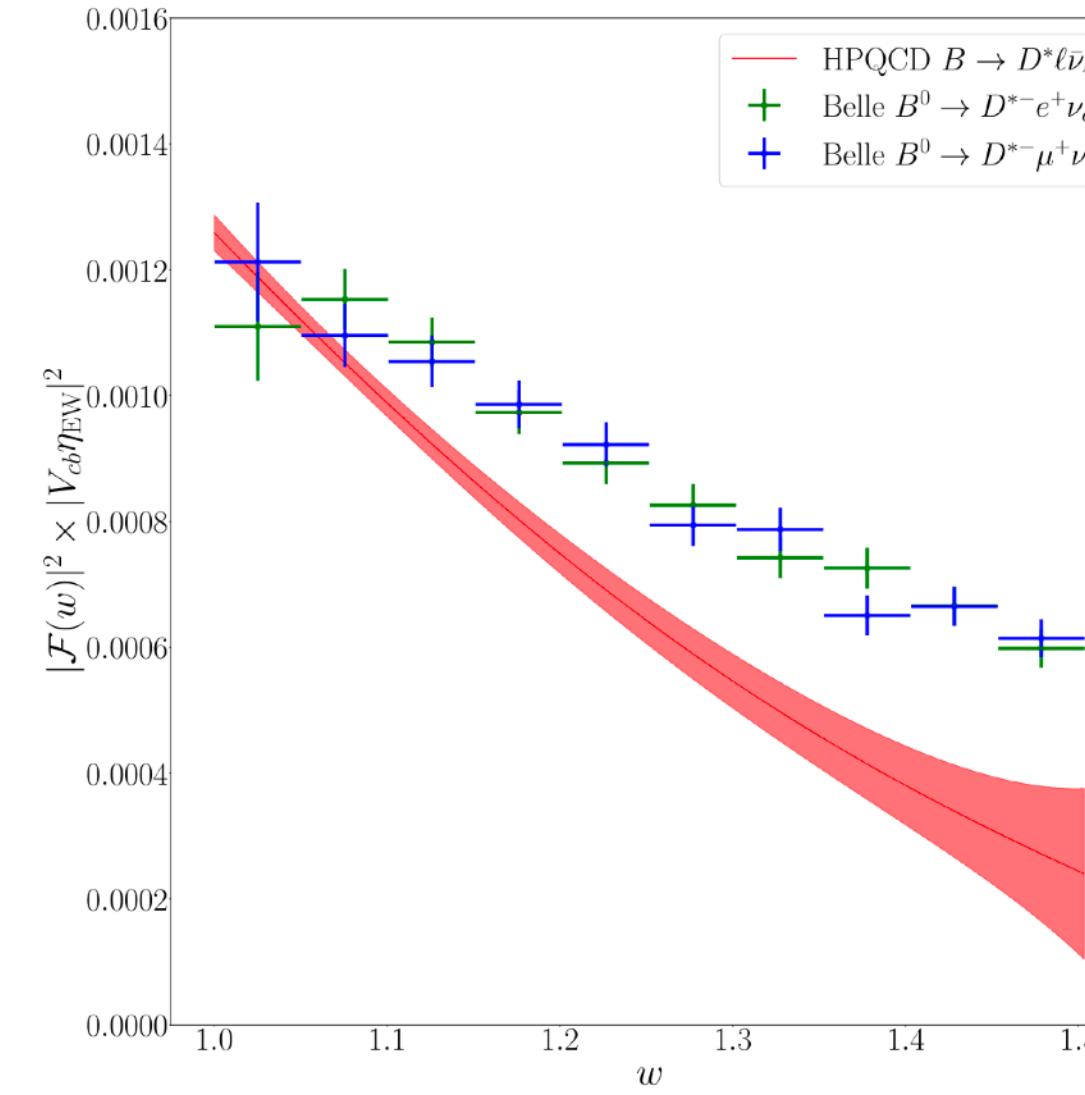
**FNAL/MILC EPJC (2022) [arXiv:2105.14019](#)**



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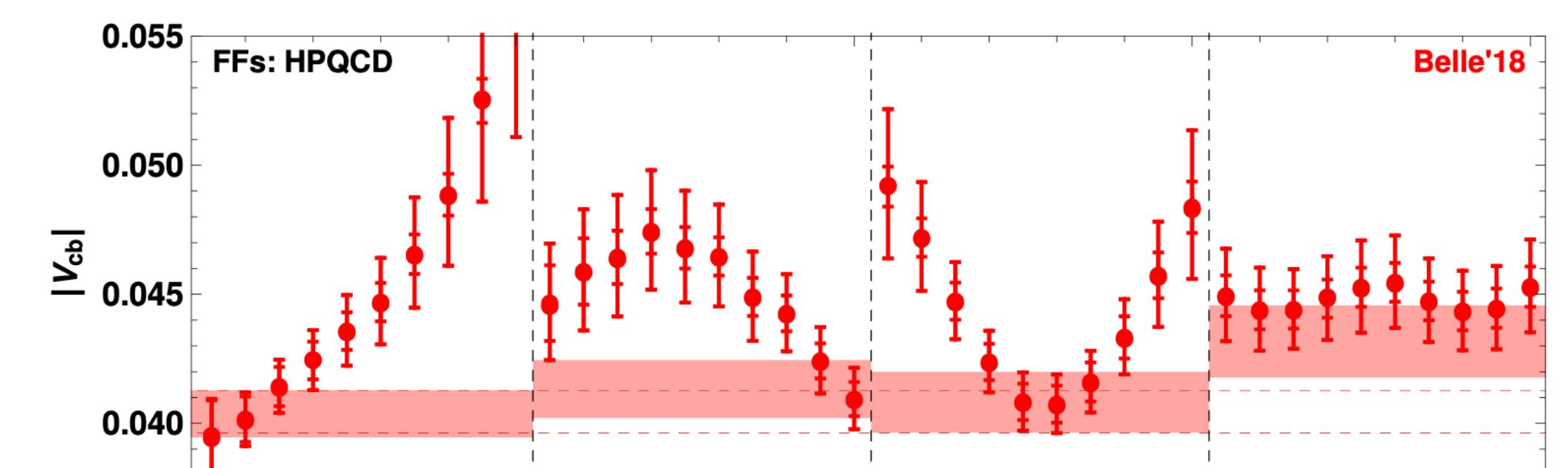
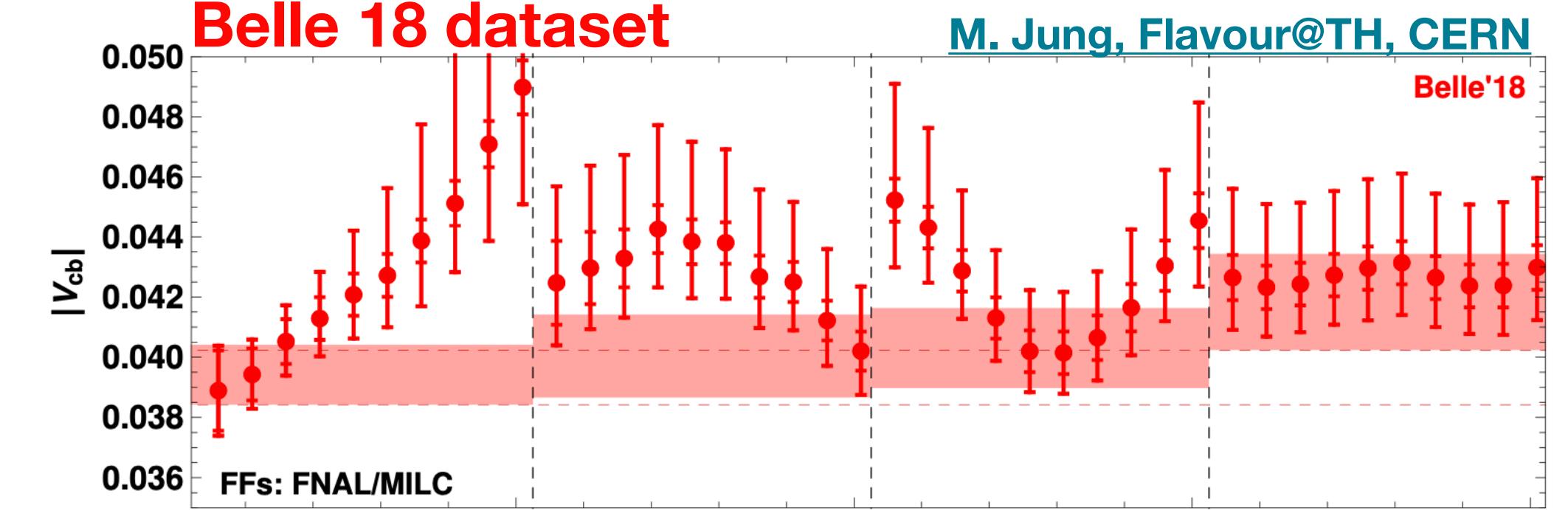
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**Belle 18 dataset**



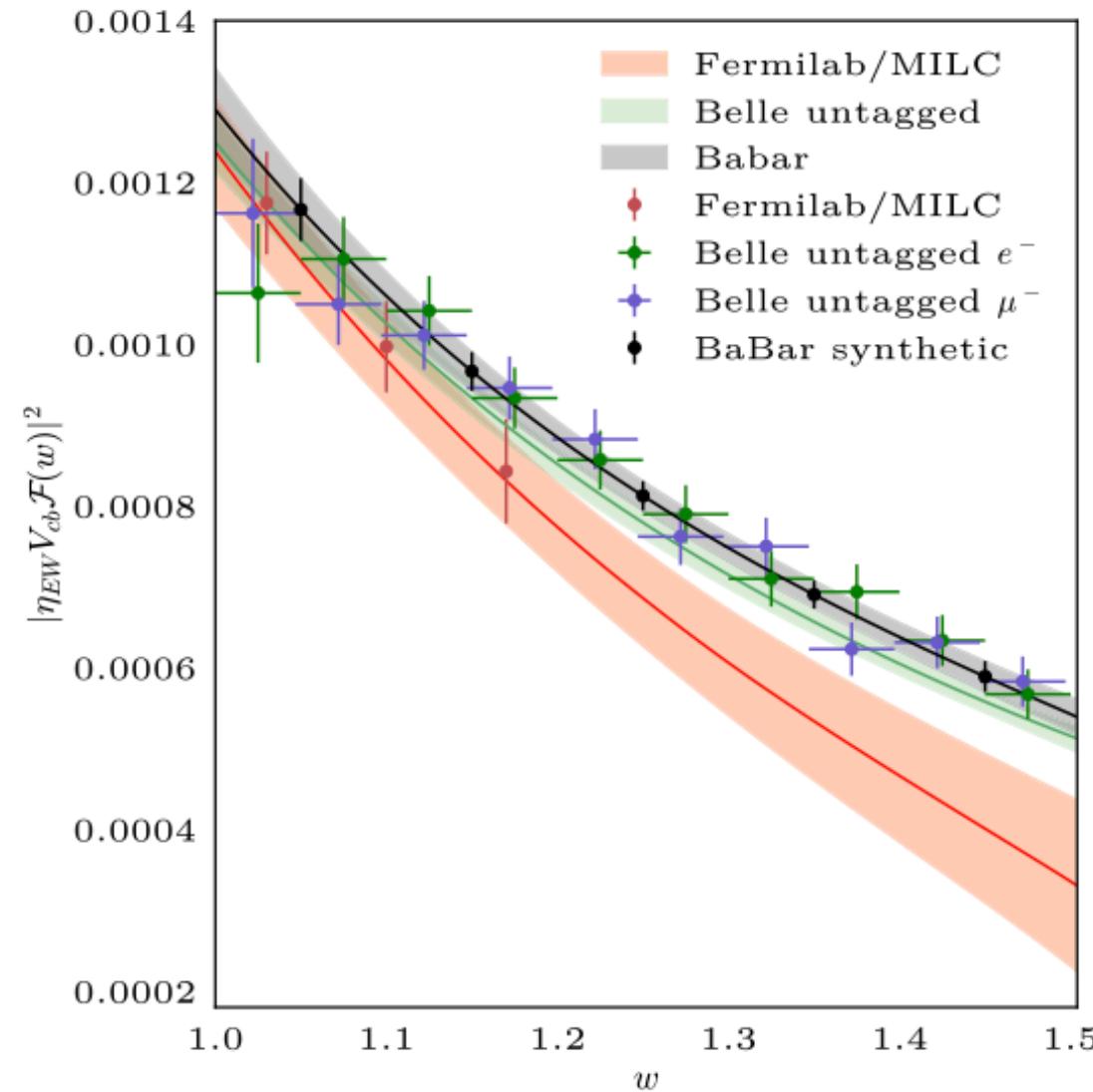
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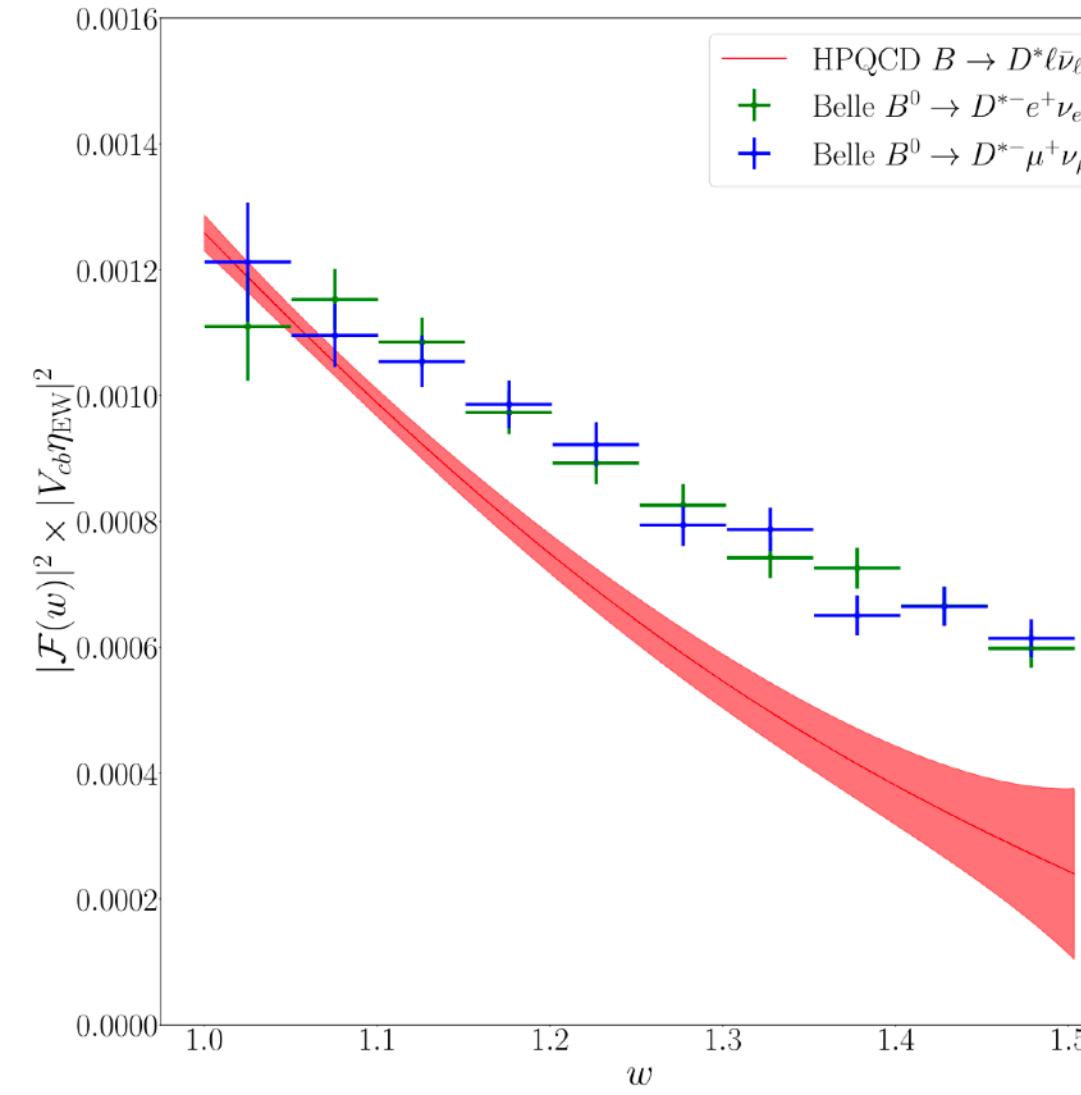
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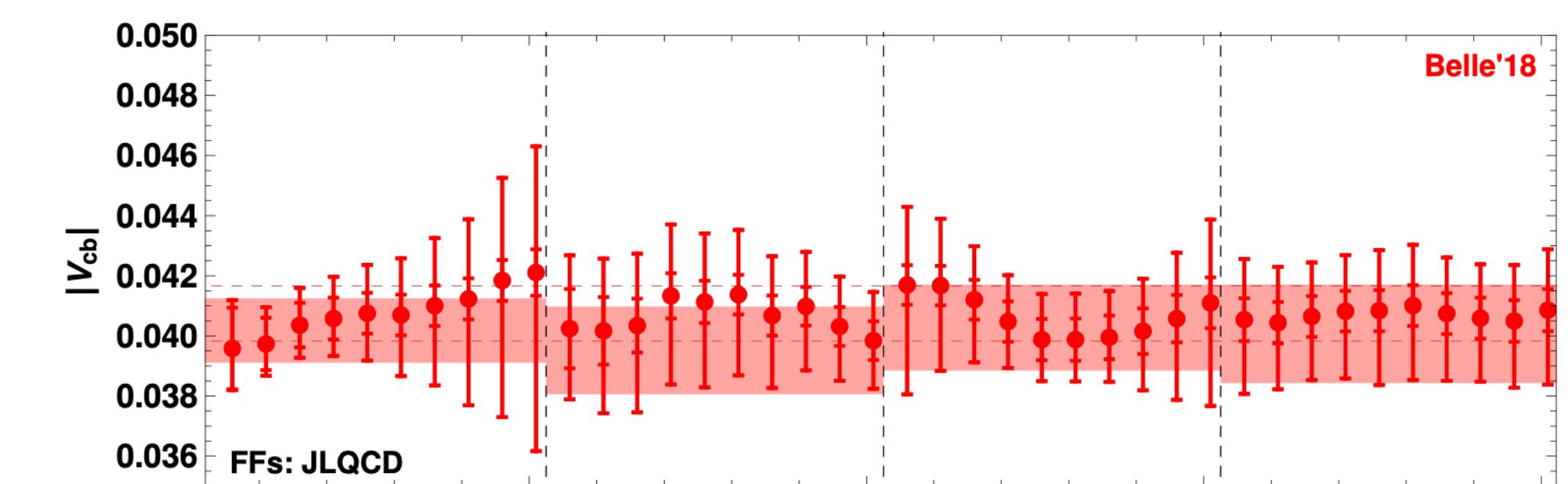
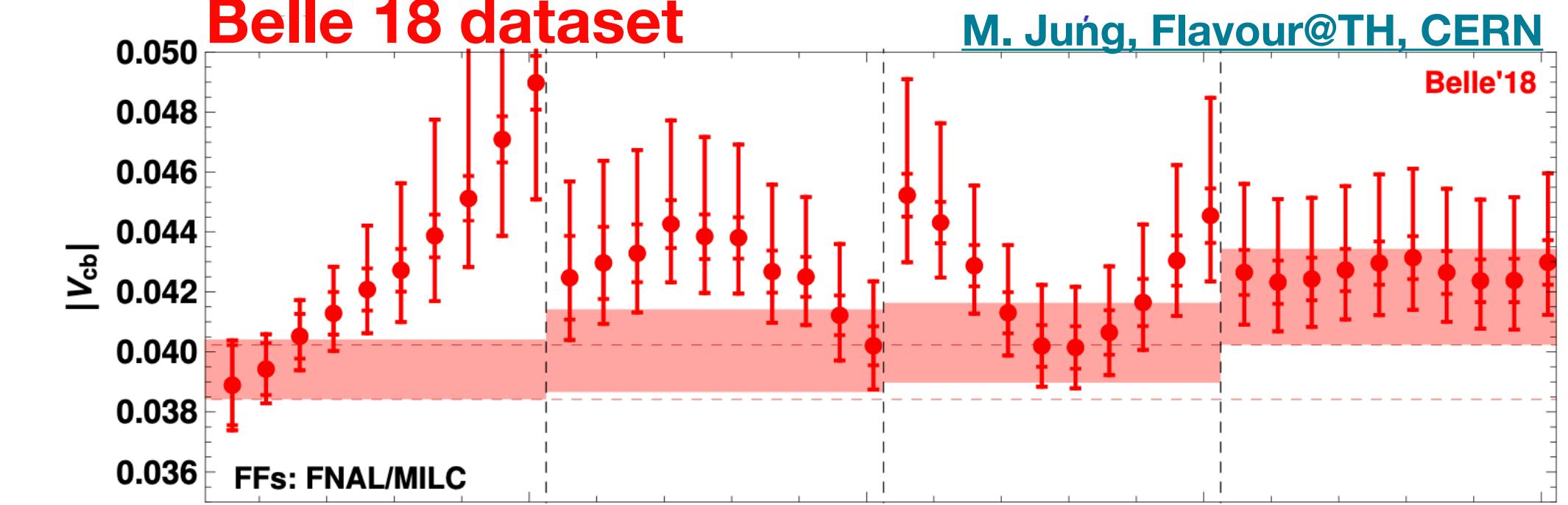
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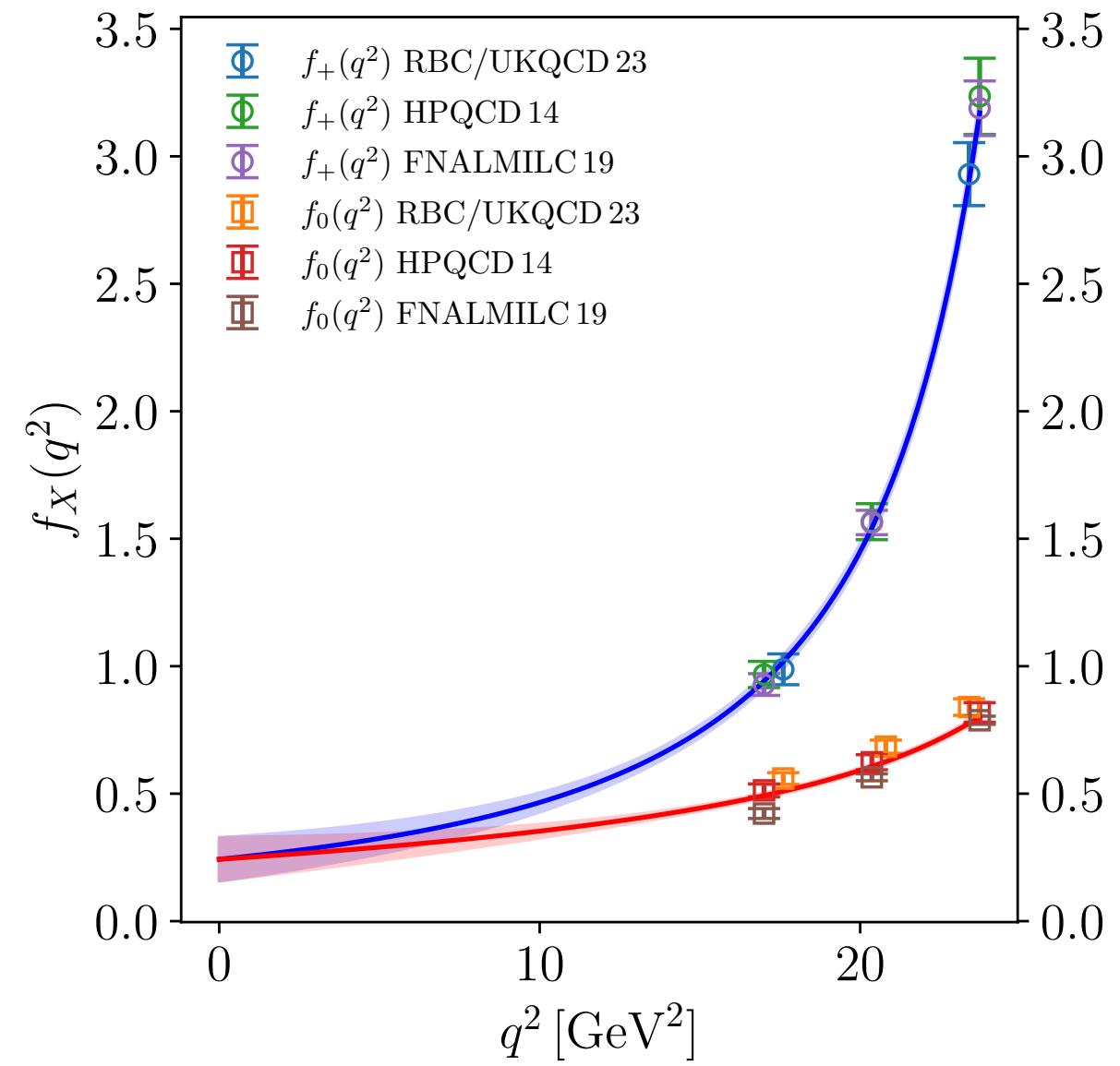
**new Belle analysis: arXiv:2301.07529**

data now publicly available:

<https://www.hepdata.net/record/ins1512299>

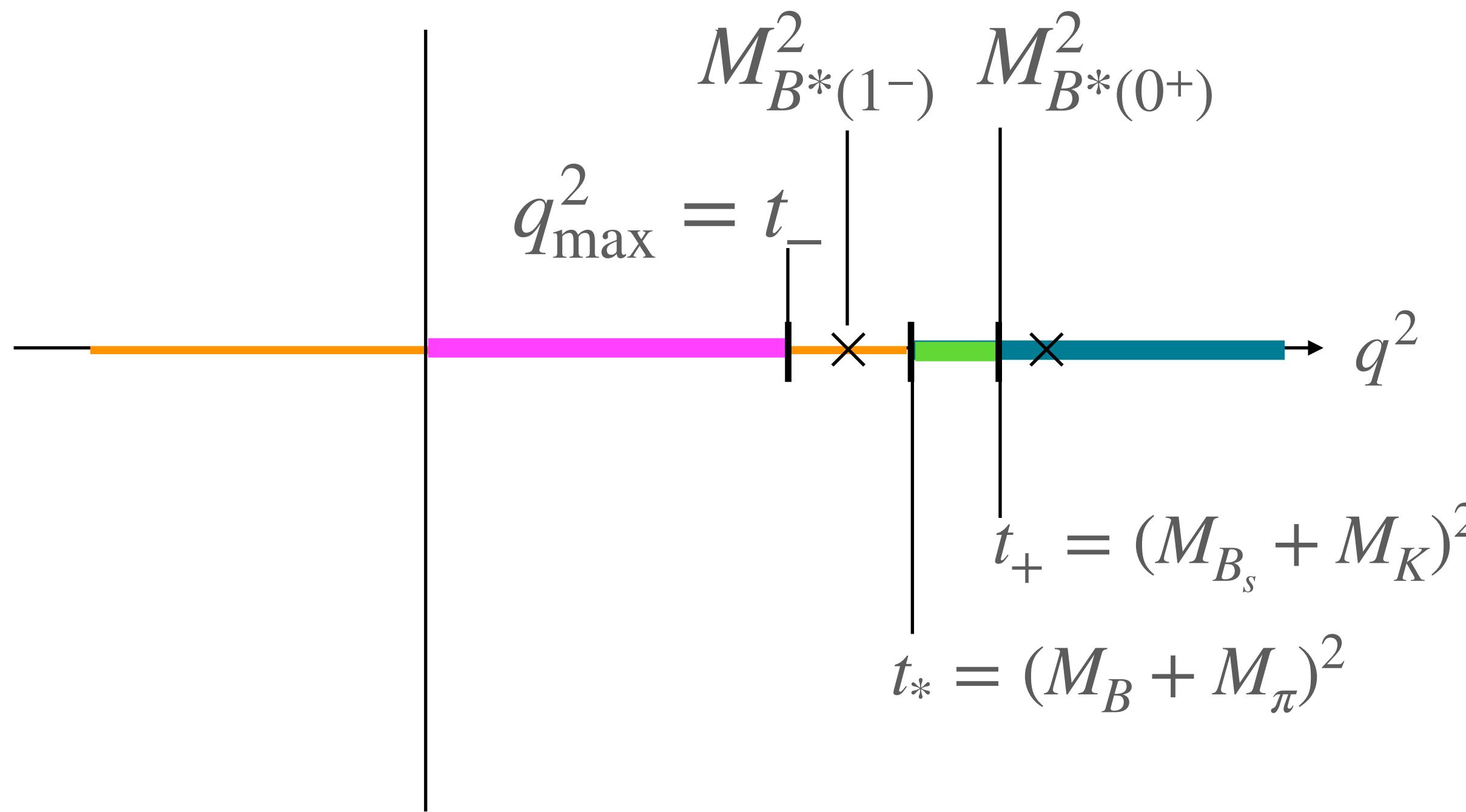
Today: Fedele arXiv:2305.15457

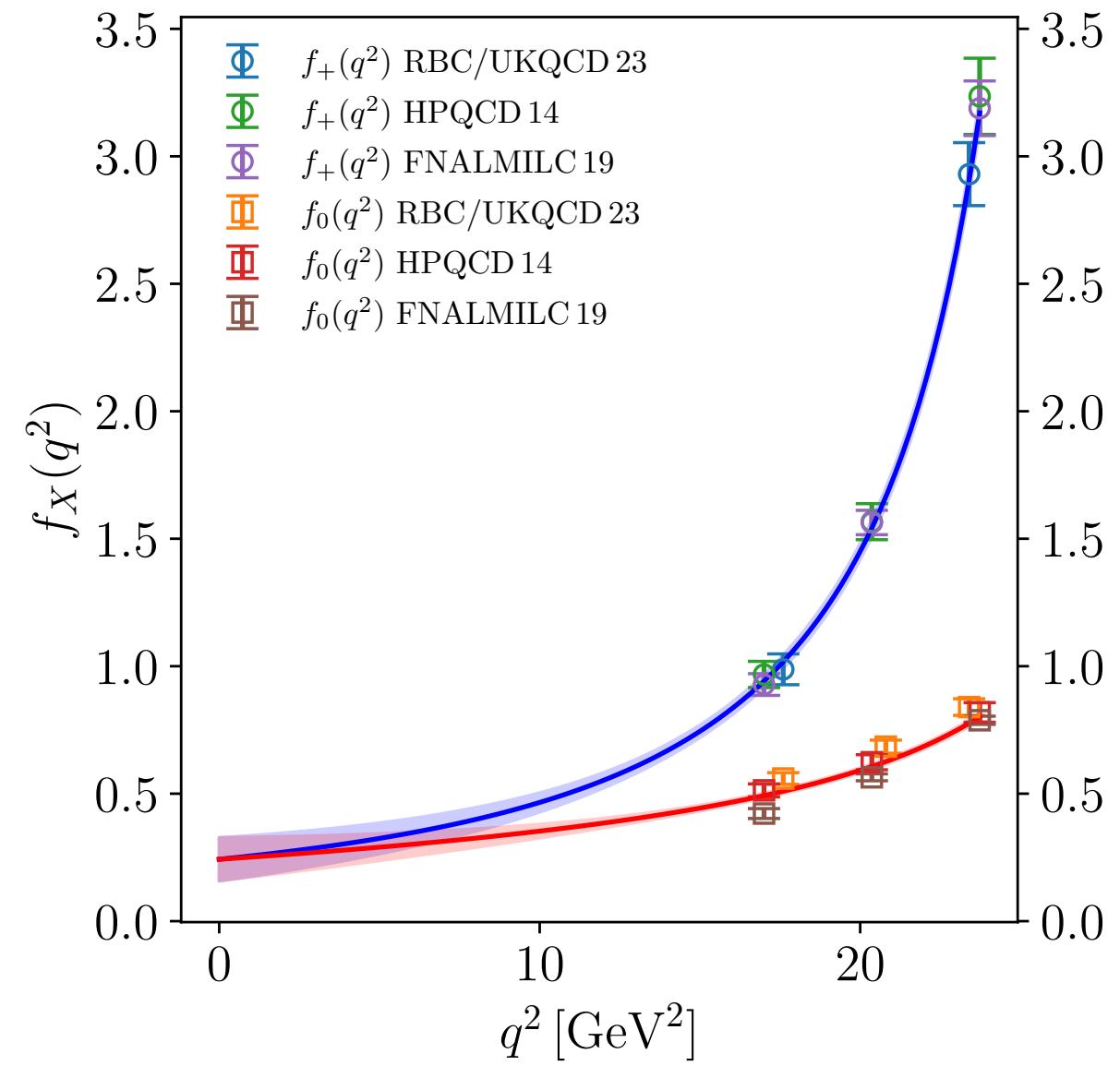
***z-fit***



# Example: $B_s \rightarrow Kl\nu$ kinematic extrapolation

Use unitarity/analyticity to devise model-independent  
and fast-converging polynomial expansion:

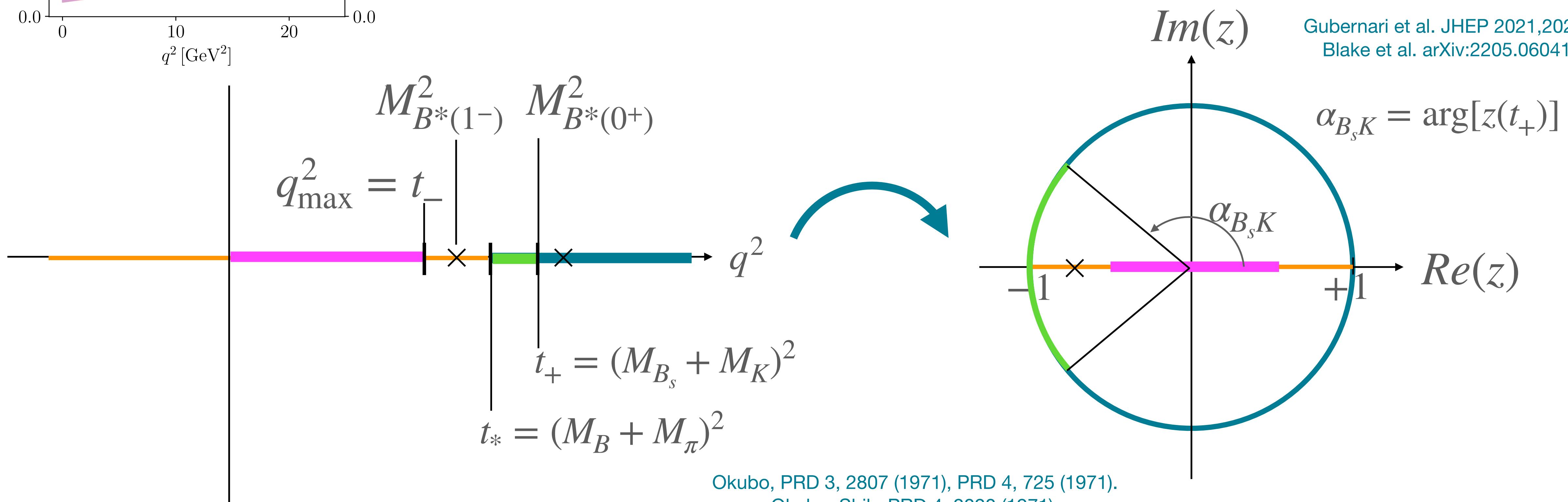




# Example: $B_s \rightarrow Kl\nu$ kinematic extrapolation

Use unitarity/analyticity to devise model-independent  
and fast-converging polynomial expansion:

Gubernari et al. JHEP 2021,2022,  
Blake et al. arXiv:2205.06041



# BGL expansion

$$f_X(q_i^2) = \frac{1}{B_X(q_i^2)\phi_X(q_i^2, t_0)} \sum_{n=0}^{K_X-1} a_{X,n} z(q_i^2)^n \quad X = +, 0$$
$$= Z_{XX,in} a_{X,n}$$

Boyd, Grinstein, Lebed, [PRL 74 \(1995\)](#)

Two constraints: **kinematic**  $f_+(0) = f_0(0)$  (eliminates one parameter in combined fit)

**unitarity constraint**  $\frac{1}{2\pi i} \oint_C \frac{dz}{z} \theta_{B_s K} |B_X(q^2)\phi_X(q^2, t_0)f_X(q^2)|^2 \leq 1$

$$a_{X,i} \langle z^i | z^j \rangle a_{X,j} \leq 1 \quad \text{Flynn, AJ, Tsang, } \text{arXiv:2303.11280}$$

# BGL fitting strategies

- Frequentist fit:**
- $N_{\text{dof}} = N_{\text{data}} - N_{\text{params}} \geq 1 \rightarrow$  in practice truncation of  $z$  expansion@low order
  - induced systematic difficult to estimate
- Bayesian fit:**
- **IDEA:** fit full  $z$  expansion (i.e. no truncation)
  - need regulator to control higher-order coefficients

# Model-independent fit

Flynn, AJ, Tsang, [arXiv:2303.11280](https://arxiv.org/abs/2303.11280)

Compute BGL parameters as expectation values  $\langle g(\mathbf{a}) \rangle = \mathcal{N} \int d\mathbf{a} g(\mathbf{a}) \pi(\mathbf{a} | \mathbf{f}, C_{\mathbf{f}}) \pi_{\mathbf{a}}$

where *probability for parameters given model and data*

$$\pi(\mathbf{a} | \mathbf{f}, C_{\mathbf{f}}) \propto \exp\left(-\frac{1}{2}\chi^2(\mathbf{a}, \mathbf{f})\right) \quad \text{where} \quad \chi^2(\mathbf{a}, \mathbf{f}) = (\mathbf{f} - Z\mathbf{a})^T C_{\mathbf{f}}^{-1} (\mathbf{f} - Z\mathbf{a})$$

where *prior knowledge just QFT*:

$$\pi_{\mathbf{a}} \propto \theta\left(1 - |\mathbf{a}_+|^2_{a_{B_sK}}\right) \theta\left(1 - |\mathbf{a}_0|^2_{a_{B_sK}}\right)$$

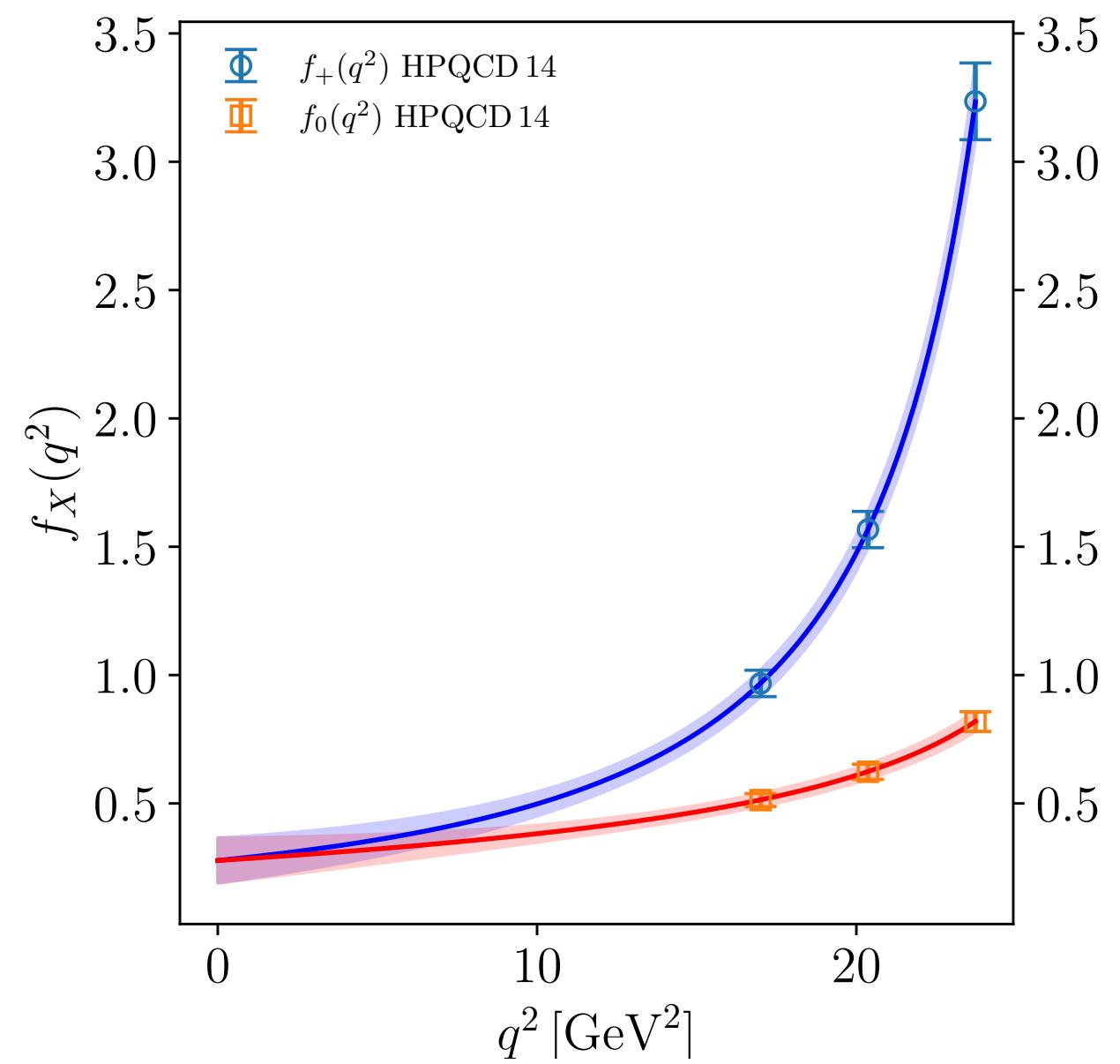
**In practice MC integration:** draw samples for  $\mathbf{a}$  from multivariate normal distribution and drop samples not compatible with unitarity

# Example: $B_s \rightarrow K\ell\nu$

HPQCD 14  
PRD 90 (2014) 054506

HPQCD 14 –  $\mathbf{a}_+$

$K_+$	$K_0$	$a_{+,0}$	$a_{+,1}$	$a_{+,2}$	$p$	$\chi^2/N_{\text{dof}}$	$N_{\text{dof}}$
2	2	0.0270(13)	-0.0792(50)	-	0.03	2.93	3
2	3	0.0273(13)	-0.0760(63)	-	0.02	4.06	2
3	2	0.0257(14)	-0.0805(50)	0.068(31)	0.15	1.89	2
3	3	0.0262(14)	-0.0727(64)	0.096(34)	0.97	0.00	1



HPQCD 14 –  $\mathbf{a}_0$

$K_+$	$K_0$	$a_{0,0}$	$a_{0,1}$	$a_{0,2}$	$p$	$\chi^2/N_{\text{dof}}$	$N_{\text{dof}}$
2	2	0.0883(44)	-0.250(17)	-	0.03	2.93	3
2	3	0.0880(44)	-0.242(19)	0.053(65)	0.02	4.06	2
3	2	0.0906(45)	-0.240(17)	-	0.15	1.89	2
3	3	0.0908(46)	-0.215(22)	0.138(71)	0.97	0.00	1

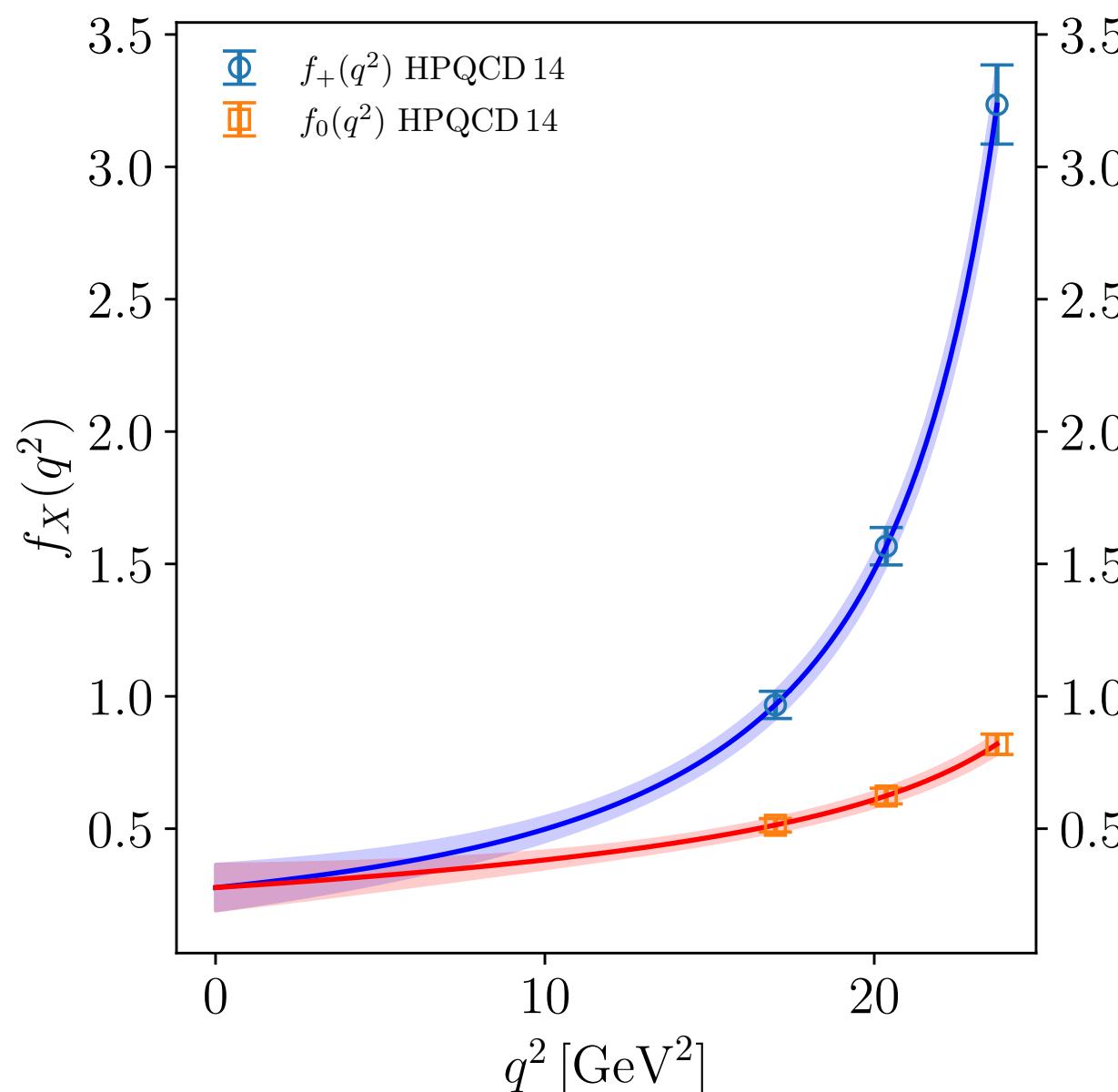
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[PRD 90 \(2014\) 054506](#)

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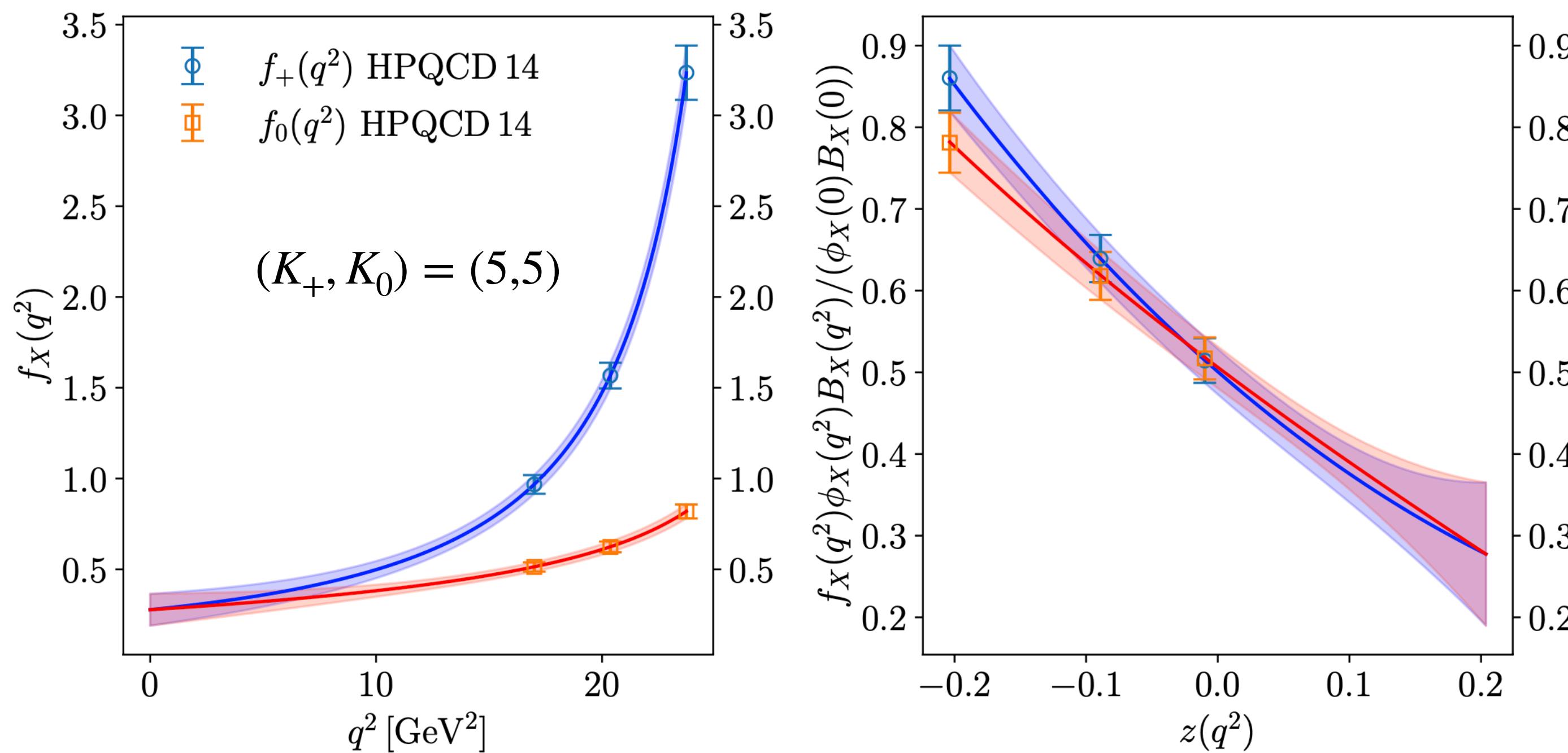
Flynn, AJ, Tsang, [arXiv:2303.11280](#)

$K_+$	$K_0$	$a_{+,0}$	$a_{+,1}$	$a_{+,2}$	$a_{+,3}$	$a_{+,4}$	$a_{+,5}$	$a_{+,6}$	$a_{+,7}$	$a_{+,8}$	$a_{+,9}$
2	2	0.0270(12)	-0.0792(49)	-	-	-	-	-	-	-	-
2	3	0.0273(13)	-0.0761(63)	-	-	-	-	-	-	-	-
3	2	0.0257(14)	-0.0805(49)	0.069(30)	-	-	-	-	-	-	-
3	3	0.0261(14)	-0.0728(64)	0.096(34)	-	-	-	-	-	-	-
3	4	0.0261(14)	-0.0728(76)	0.096(39)	-	-	-	-	-	-	-
4	3	0.0261(14)	-0.0729(68)	0.096(35)	0.008(90)	-	-	-	-	-	-
4	4	0.0261(14)	-0.0730(77)	0.091(62)	-0.02(20)	-	-	-	-	-	-
5	5	0.0262(15)	-0.0735(79)	0.084(67)	-0.03(19)	0.03(68)	-	-	-	-	-
6	6	0.0261(14)	-0.0735(79)	0.086(69)	-0.03(19)	-0.00(64)	0.01(65)	-	-	-	-
7	7	0.0262(14)	-0.0732(84)	0.088(69)	-0.02(18)	0.01(65)	0.02(73)	-0.03(70)	-	-	-
8	8	0.0261(14)	-0.0732(80)	0.089(72)	-0.02(18)	-0.00(66)	0.03(86)	-0.04(90)	0.03(73)	-	-
9	9	0.0261(14)	-0.0729(84)	0.095(75)	-0.02(19)	-0.04(68)	0.1(1.0)	-0.1(1.2)	0.1(1.1)	-0.06(79)	-
10	10	0.0261(14)	-0.0726(89)	0.101(79)	-0.01(20)	-0.09(73)	0.2(1.3)	-0.3(1.7)	0.2(1.8)	-0.2(1.4)	0.08(87)

HPQCD 14 –  $\mathbf{a}_0$

$K_+$	$K_0$	$a_{0,0}$	$a_{0,1}$	$a_{0,2}$	$a_{0,3}$	$a_{0,4}$	$a_{0,5}$	$a_{0,6}$	$a_{0,7}$	$a_{0,8}$	$a_{0,9}$
2	2	0.0883(44)	-0.250(17)	-	-	-	-	-	-	-	-
2	3	0.0880(44)	-0.243(19)	0.052(65)	-	-	-	-	-	-	-
3	2	0.0907(46)	-0.240(17)	-	-	-	-	-	-	-	-
3	3	0.0906(44)	-0.215(22)	0.137(73)	-	-	-	-	-	-	-
3	4	0.0907(47)	-0.215(22)	0.14(11)	-0.01(31)	-	-	-	-	-	-
4	3	0.0907(45)	-0.214(22)	0.139(72)	-	-	-	-	-	-	-
4	4	0.0907(46)	-0.215(25)	0.12(19)	-0.08(60)	-	-	-	-	-	-
5	5	0.0909(46)	-0.218(25)	0.10(19)	-0.12(55)	0.04(63)	-	-	-	-	-
6	6	0.0907(45)	-0.217(25)	0.10(19)	-0.11(53)	0.06(66)	-0.02(66)	-	-	-	-
7	7	0.0907(46)	-0.217(26)	0.11(20)	-0.08(51)	0.03(73)	0.03(81)	-0.04(70)	-	-	-
8	8	0.0908(46)	-0.217(25)	0.11(20)	-0.08(50)	-0.01(84)	0.1(1.0)	-0.09(96)	0.08(74)	-	-
9	9	0.0907(46)	-0.215(25)	0.13(22)	-0.05(50)	-0.06(95)	0.2(1.4)	-0.2(1.5)	0.1(1.2)	-0.05(82)	-
10	10	0.0907(46)	-0.214(27)	0.15(24)	-0.03(49)	-0.2(1.1)	0.4(1.8)	-0.5(2.2)	0.4(2.1)	-0.3(1.6)	0.13(90)

# Summary $z$ -fits



- new, clean, truncation and bias-free form-factor parameterisation
- works for any hadronic form factor
- use complementarity of frequentist and Bayesian analysis

Flynn, AJ, Tsang, [arXiv:2303.11280](https://arxiv.org/abs/2303.11280)

## Also have a look:

- Dispersive-matrix method, Di Carlo et al. PRD 2021, [arXiv:2105.02497](https://arxiv.org/abs/2105.02497)
- Self-consistency checks of  $z$  expansion, Simons, Gustafson, Meurice [arXiv:2304.13045](https://arxiv.org/abs/2304.13045)

# Code: BFF

Flynn, AJ, Tsang, arXiv:2303.11280

Python3 available via [github/Zenodo](#)

<https://github.com/andreasjuettner/BFF>

<https://zenodo.org/record/7799543#.ZEezTy8Ro80>

```
}

#####
# specify input for BGL fit
#####

input_dict = {
    'decay':      'Btopi',
    'Mi':        pc.mBphys,      # initial-state mass
    'Mo':        pc.mpiphys,    # final-state mass
    'sigma':     .5,            # sigma for prior in algorithm
    'Kp':         4,             # target Kp (BGL truncation) - can be changed later
    'K0':         4,             # target K0 (BGL truncation) - can be changed later
    'tstar':     '29.349570696829012', # value of t*
    't0':        'self.tstar - np.sqrt(self.tstar*(self.tstar-self.tm))', # definition of t0
    'chip':       pc.chip_Btopi, # susceptibility fp
    'chi0':       pc.chi0_Btopi, # susceptibility f0
    'mpolep':    [pc.mBstar],   # fplus pole
    'mpole0':    [],            # fzero pole (no pole for BstoK)
    'N' :          N,            # number of desired samples
    'outer_p':   [3./2,'48*np.pi',3,2], # specs for outer function fp
    'outer_0':   [3./2,'16*np.pi/(self.tp*self.tm)',1,1], # specs for outer function f0
    'seed':       123,           # RNG seed
```

```
input_data = {
    'RBCUKQCD 23 lat':
    {
        'data type':   'ff',
        'label':        'RBC/UKQCD 23',
        'Np':          2,
        'N0':          3,
        'qsqp':        np.array([17.60,23.40]),
        'qsq0':        np.array([17.60,20.80,23.40]),
        'fp':          fparray,
        'f0':          f0array,
        'Cff':         cov_array
    }
}
```

**inclusive  $b \rightarrow c$**

# $b \rightarrow c$ inclusive on the lattice

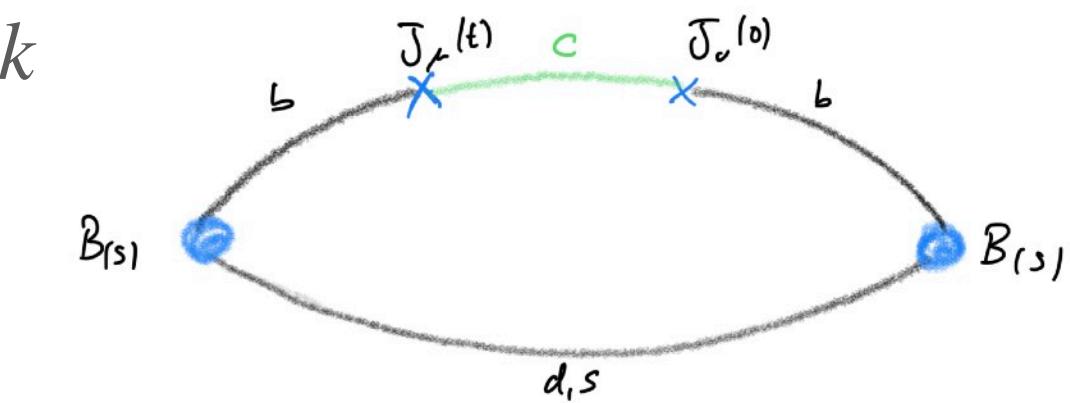
long-standing tension inclusive vs. exclusive CKM determination → find new ways to look at it

$$\Gamma(B_s \rightarrow X_c l \nu) = \frac{G_F^2 |V_{cb}|^2}{24\pi^3} \int_0^{q_{\max}^2} dq^2 \sqrt{q^2} \bar{X}(q)$$

Gambino, Hashimoto, PRL (2020) [arXiv:2005.13730](https://arxiv.org/abs/2005.13730)

Hansen, Meyer, Robaina, PRD (2017), [arXiv:1704.08993](https://arxiv.org/abs/1704.08993)

$$\bar{X}(q) = \sum_k c_{\mu\nu,k}(q) C_{\mu\nu}(ak, q)$$



# $b \rightarrow c$ inclusive on the lattice

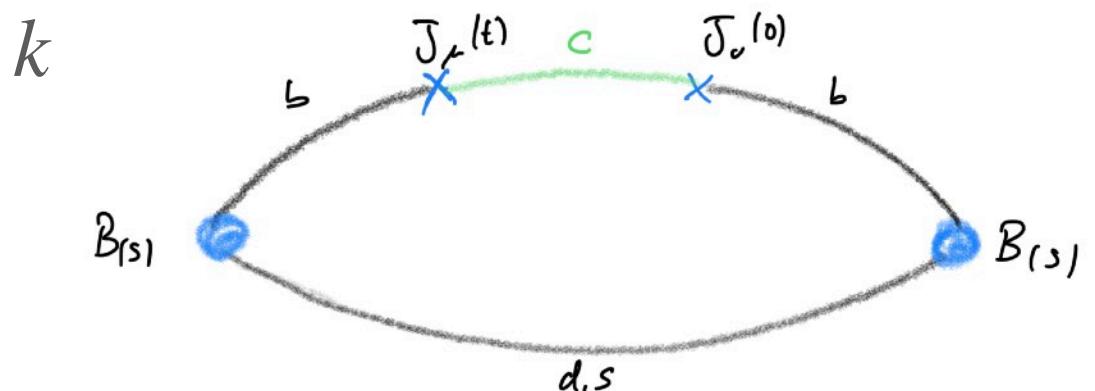
long-standing tension inclusive vs. exclusive CKM determination → find new ways to look at it

$$\Gamma(B_s \rightarrow X_c l \nu) = \frac{G_F^2 |V_{cb}|^2}{24\pi^3} \int_0^{q^2_{\max}} dq^2 \sqrt{q^2} \bar{X}(q)$$

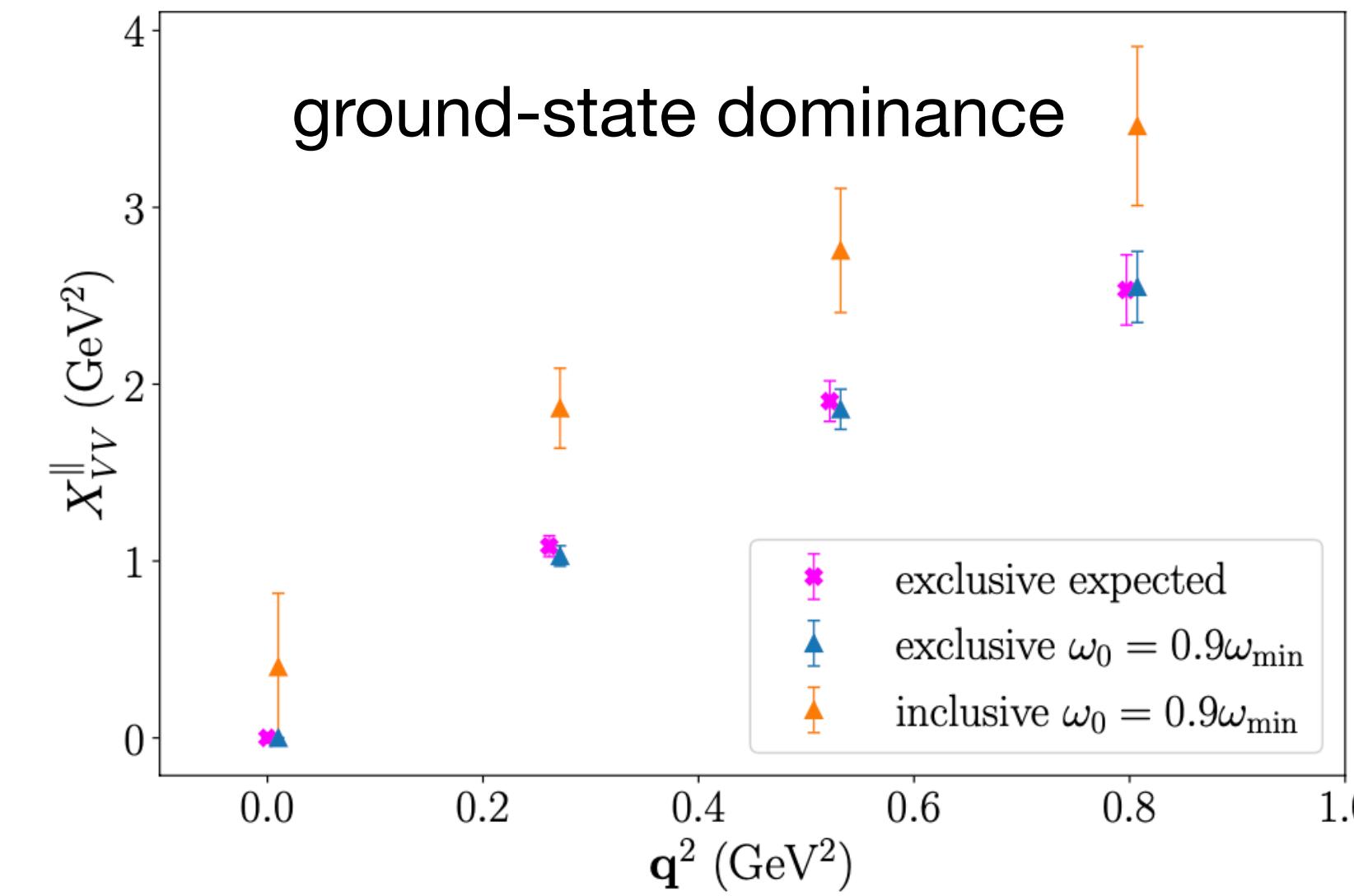
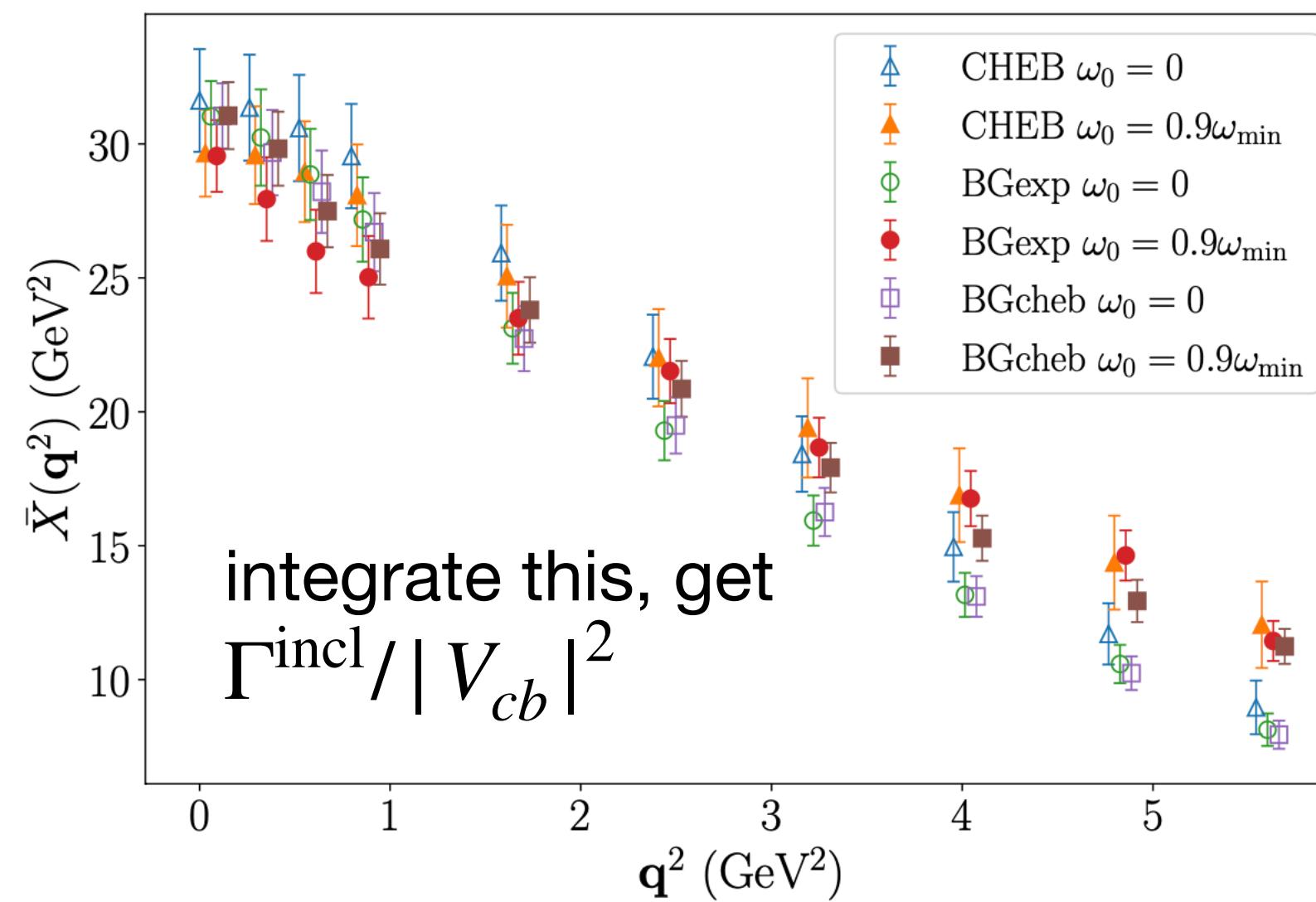
Gambino, Hashimoto, PRL (2020) arXiv:2005.13730

Hansen, Meyer, Robaina, PRD (2017), arXiv:1704.08993

$$\bar{X}(q) = \sum_k c_{\mu\nu,k}(q) C_{\mu\nu}(ak, q)$$



Barone, Hashimoto, AJ, Kaneko, Kellerman arXiv:2305.14092



- PRELIMINARY study
- concentrating on controlling systematics
- could be interesting to define new observables
- test continuum ↔ lattice
- Next: first comprehensive study with error budget and prediction

# Summary

- New results for exclusive  $b \rightarrow u$  and  $b \rightarrow c$  transitions
- Some tensions in lattice data require attention
- new ideas for truncation-independent form factor fit
- inclusive decays on the lattice exciting new area of research

**Other interesting work I would have liked to cover:**

$B \rightarrow \rho, B \rightarrow K^*$  beyond narrow-width approximation (Leskovec, Lattice22 talk)

# BACKUP

# Examples: $B_s \rightarrow K\ell\nu$

pheno results

$K_+$	$K_0$	$f(q^2 = 0)$	$R_{B_s \rightarrow K}^{\text{impr}}$	$R_{B_s \rightarrow K}$	$\frac{\Gamma^\tau}{ V_{ub} ^2} [\frac{1}{\text{ps}}]$	$\frac{\Gamma^\mu}{ V_{ub} ^2} [\frac{1}{\text{ps}}]$	$V_{\text{CKM}}^{\text{low}}$	$V_{\text{CKM}}^{\text{high}}$	$V_{\text{CKM}}^{\text{full}}$
2	2	0.208(25)	1.524(37)	0.727(25)	4.51(45)	6.23(76)	0.00383(47)	0.00352(35)	0.00363(37)
2	3	0.226(34)	1.511(41)	0.704(39)	4.67(49)	6.67(97)	0.00361(53)	0.00344(34)	0.00349(38)
3	2	0.233(27)	1.609(58)	0.733(27)	4.44(45)	6.08(77)	0.00368(45)	0.00367(37)	0.00367(38)
3	3	0.293(41)	1.592(57)	0.664(40)	4.84(51)	7.3(1.1)	0.00310(44)	0.00349(35)	0.00333(36)
3	4	0.293(56)	1.593(60)	0.667(59)	4.85(58)	7.4(1.4)	0.00313(55)	0.00349(37)	0.00338(40)
4	3	0.294(42)	1.594(60)	0.663(40)	4.85(52)	7.4(1.1)	0.00309(44)	0.00348(36)	0.00332(36)
4	4	0.285(92)	1.593(60)	0.677(88)	4.83(62)	7.3(1.7)	0.00328(86)	0.00350(38)	0.00346(42)
5	5	0.277(88)	1.595(62)	0.685(85)	4.81(62)	7.2(1.7)	0.00333(85)	0.00351(38)	0.00348(42)
6	6	0.277(88)	1.592(63)	0.685(86)	4.79(63)	7.2(1.7)	0.00335(88)	0.00350(38)	0.00348(43)
7	7	0.282(89)	1.592(60)	0.680(87)	4.82(64)	7.3(1.7)	0.00332(89)	0.00350(38)	0.00347(43)
8	8	0.283(88)	1.594(61)	0.679(85)	4.83(64)	7.3(1.7)	0.00330(85)	0.00351(37)	0.00347(41)
9	9	0.289(91)	1.594(62)	0.674(88)	4.85(64)	7.4(1.8)	0.00327(89)	0.00350(38)	0.00347(42)
10	10	0.293(95)	1.593(60)	0.670(91)	4.87(67)	7.5(1.9)	0.00325(92)	0.00349(38)	0.00346(42)

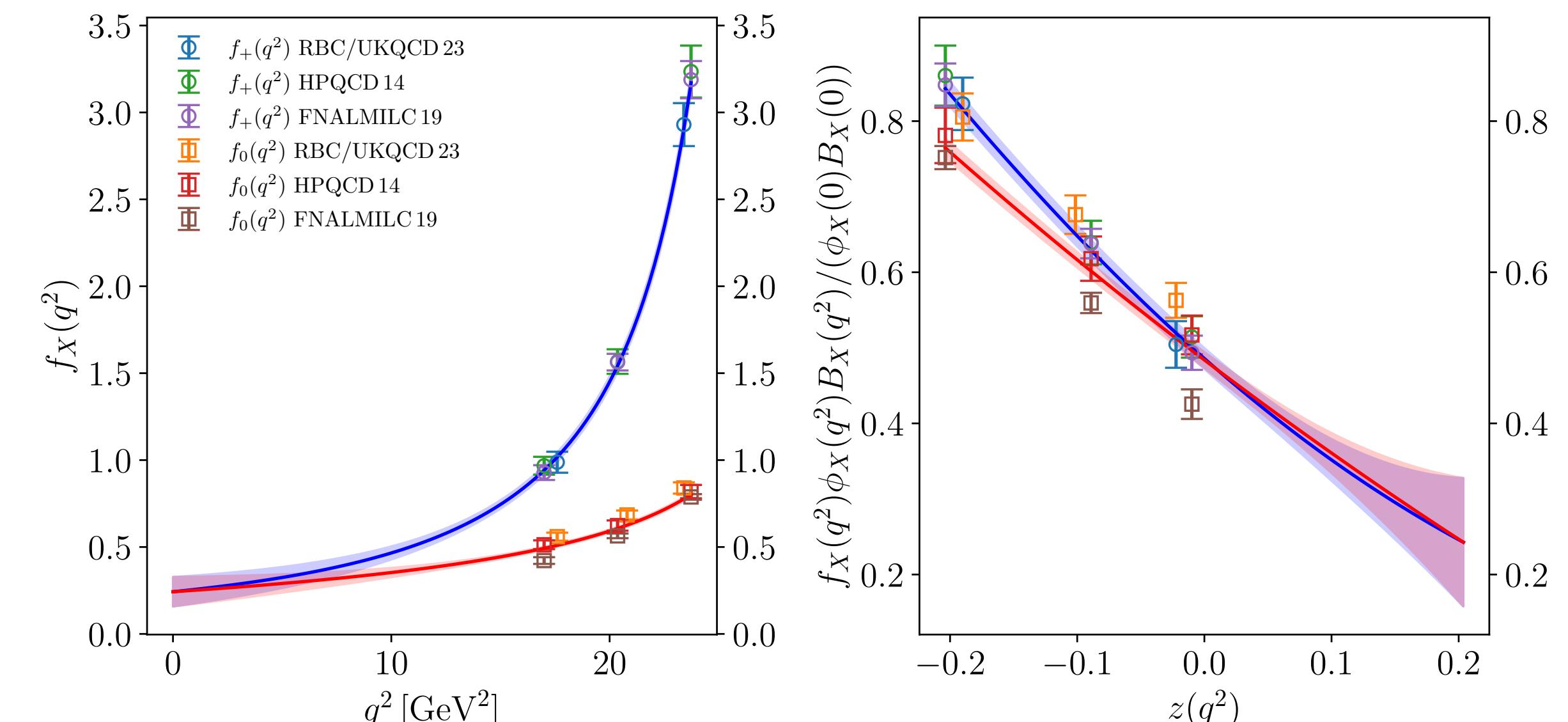
$K_+$	$K_0$	$I[\mathcal{A}_{\text{FB}}^\tau] [\frac{1}{\text{ps}}]$	$I[\mathcal{A}_{\text{FB}}^\mu] [\frac{1}{\text{ps}}]$	$\bar{\mathcal{A}}_{\text{FB}}^\tau$	$\bar{\mathcal{A}}_{\text{FB}}^\mu$	$I[\mathcal{A}_{\text{pol}}^\tau] [\frac{1}{\text{ps}}]$	$I[\mathcal{A}_{\text{pol}}^\mu] [\frac{1}{\text{ps}}]$	$\bar{\mathcal{A}}_{\text{pol}}^\tau$	$\bar{\mathcal{A}}_{\text{pol}}^\mu$
2	2	1.22(13)	0.0278(51)	0.2708(37)	0.00443(34)	0.74(15)	6.15(75)	0.164(29)	0.98767(96)
2	3	1.26(14)	0.0314(70)	0.2709(38)	0.00465(44)	0.81(18)	6.59(96)	0.173(31)	0.9872(12)
3	2	1.23(13)	0.0319(59)	0.2780(43)	0.00524(51)	0.46(19)	5.99(76)	0.103(40)	0.9852(15)
3	3	1.36(15)	0.045(10)	0.2814(48)	0.00612(66)	0.53(20)	7.2(1.1)	0.110(40)	0.9830(18)
3	4	1.37(17)	0.046(14)	0.2814(50)	0.00611(83)	0.53(22)	7.3(1.3)	0.109(41)	0.9830(22)
4	3	1.37(15)	0.046(10)	0.2815(50)	0.00616(71)	0.53(22)	7.2(1.1)	0.109(42)	0.9829(20)
4	4	1.36(19)	0.046(21)	0.2810(69)	0.0060(15)	0.53(21)	7.2(1.7)	0.109(42)	0.9834(41)
5	5	1.35(19)	0.044(20)	0.2806(67)	0.0058(15)	0.53(22)	7.1(1.6)	0.109(44)	0.9837(39)
6	6	1.35(20)	0.044(20)	0.2803(69)	0.0058(15)	0.53(22)	7.1(1.7)	0.111(44)	0.9838(39)
7	7	1.35(20)	0.045(20)	0.2806(69)	0.0059(15)	0.53(21)	7.2(1.7)	0.111(43)	0.9835(39)
8	8	1.36(20)	0.045(20)	0.2808(69)	0.0059(15)	0.53(22)	7.2(1.7)	0.109(44)	0.9835(39)
9	9	1.36(20)	0.047(21)	0.2812(71)	0.0060(15)	0.53(22)	7.3(1.7)	0.109(44)	0.9832(40)
10	10	1.37(21)	0.048(23)	0.2815(72)	0.0061(15)	0.53(22)	7.4(1.8)	0.109(43)	0.9831(41)

# $B_s \rightarrow K\ell\nu$ — continued

Easy to combine independent or correlated data sets:

$K_+$	$K_0$	$a_{+,0}$	$a_{+,1}$	$a_{+,2}$	$a_{+,3}$	$a_{+,4}$	$p$	$\chi^2/N_{\text{dof}}$	$N_{\text{dof}}$
2	2	0.02641(58)	-0.0824(26)	-	-	-	0.00	5.15	14
2	3	0.02668(68)	-0.0811(31)	-	-	-	0.00	5.50	13
3	2	0.02477(68)	-0.0829(26)	0.054(12)	-	-	0.00	3.95	13
3	3	0.02534(73)	-0.0792(31)	0.062(12)	-	-	0.00	3.89	12
3	4	0.02534(73)	-0.0781(34)	0.067(14)	-	-	0.00	4.19	11
4	3	0.02535(73)	-0.0776(38)	0.074(20)	0.023(30)	-	0.00	4.19	11
4	4	0.02592(97)	-0.033(50)	0.69(69)	2.1(2.3)	-	0.00	4.53	10
5	5	0.0266(10)	0.052(65)	2.21(97)	11.1(5.6)	17.2(15.1)	0.00	5.04	8

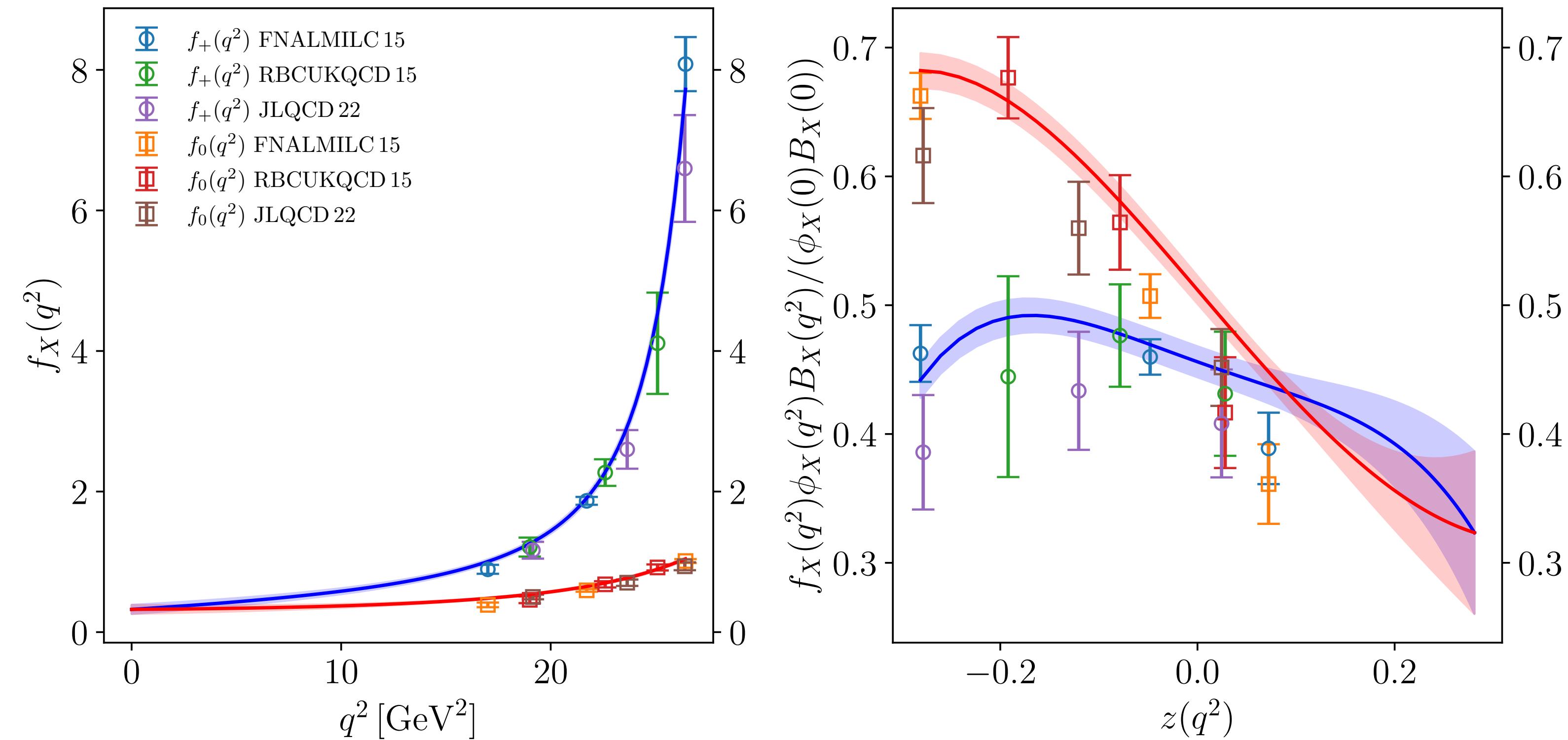
$K_+$	$K_0$	$a_{0,0}$	$a_{0,1}$	$a_{0,2}$	$a_{0,3}$	$a_{0,4}$	$p$	$\chi^2/N_{\text{dof}}$	$N_{\text{dof}}$
2	2	0.0854(17)	-0.2565(75)	-	-	-	0.00	5.15	14
2	3	0.0856(18)	-0.2527(91)	0.021(27)	-	-	0.00	5.50	13
3	2	0.0858(18)	-0.2501(77)	-	-	-	0.00	3.95	13
3	3	0.0864(18)	-0.2379(95)	0.061(28)	-	-	0.00	3.89	12
3	4	0.0869(19)	-0.231(13)	0.067(29)	-0.08(10)	-	0.00	4.19	11
4	3	0.0869(19)	-0.229(15)	0.091(48)	-	-	0.00	4.19	11
4	4	0.0887(27)	-0.08(17)	2.2(2.4)	7.0(7.9)	-	0.00	4.53	10
5	5	0.0887(28)	0.07(20)	6.1(3.3)	41.5(19.0)	93.3(44.0)	0.00	5.04	8



Bayesian and frequentist provide complementary information — consider both simultaneously!

Conclusion: World lattice data for  $B_s \rightarrow K\ell\nu$  is in quite bad shape...

# $B \rightarrow \pi \ell \nu$

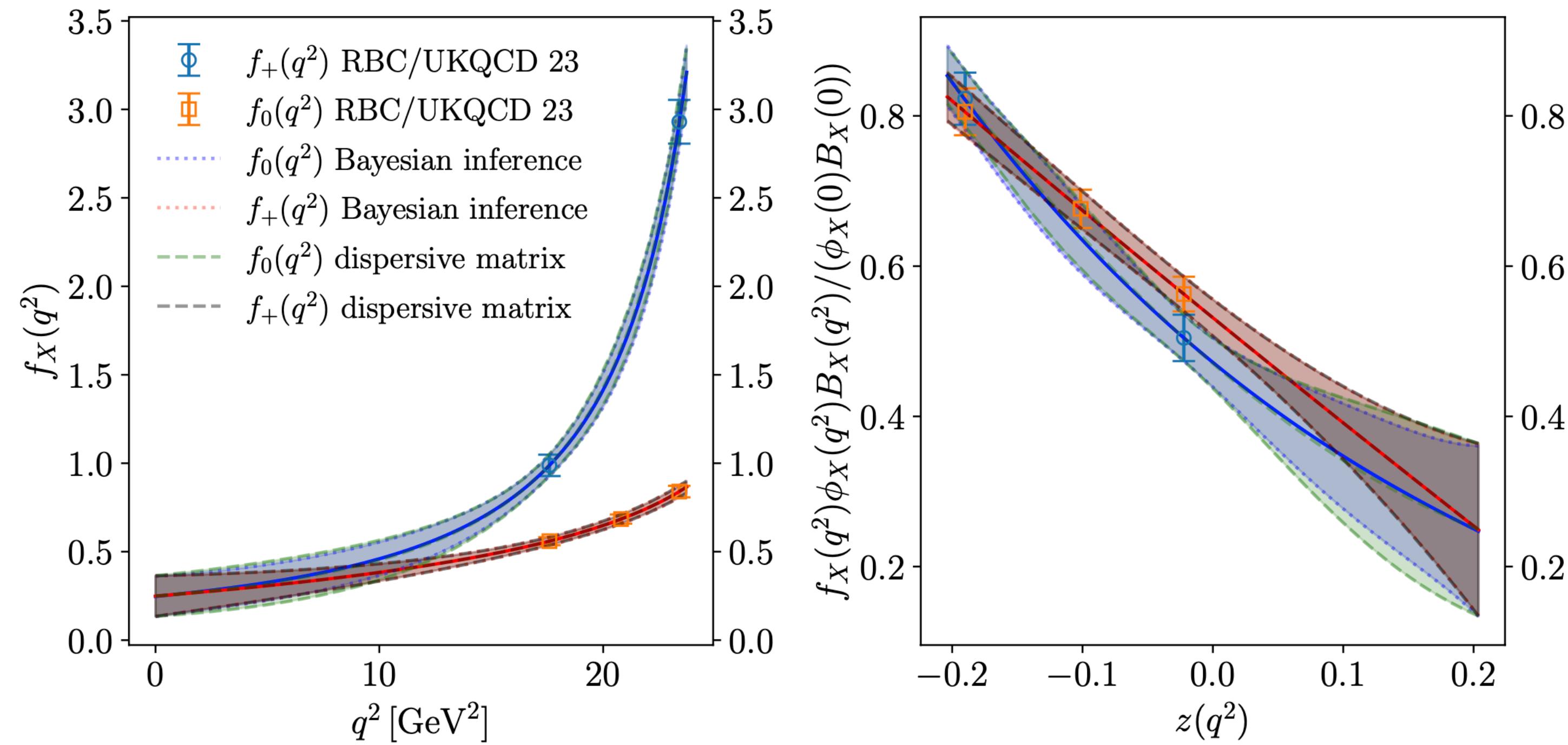


Bayesian and frequentist provide complementary information – consider both simultaneously!

Conclusion: World lattice data for  $B \rightarrow \pi \ell \nu$  is in quite bad shape...

# Relation to dispersive-matrix method?

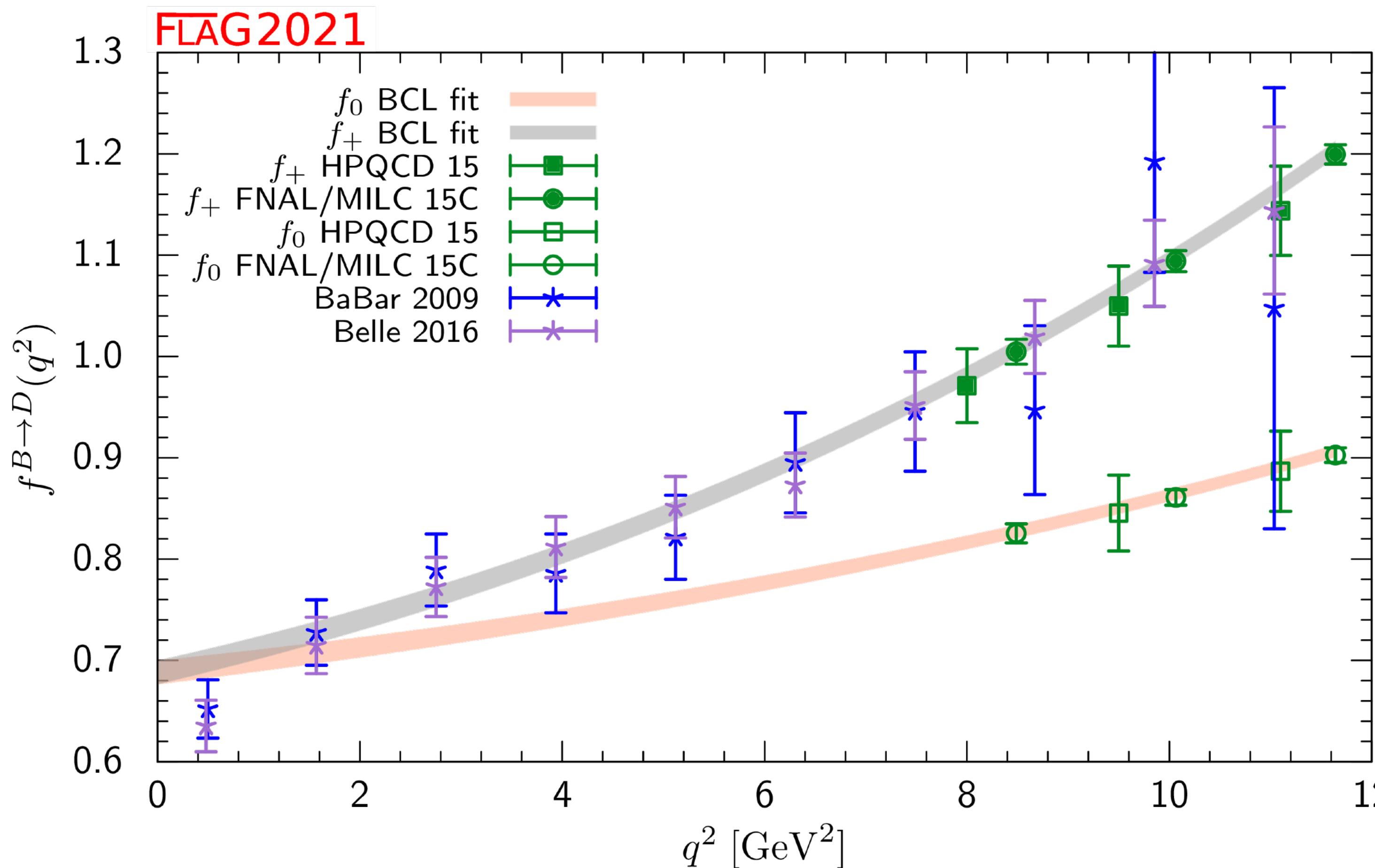
Di Carlo, Martinelli, Naviglio et al. PRD 104 (2021) 054502

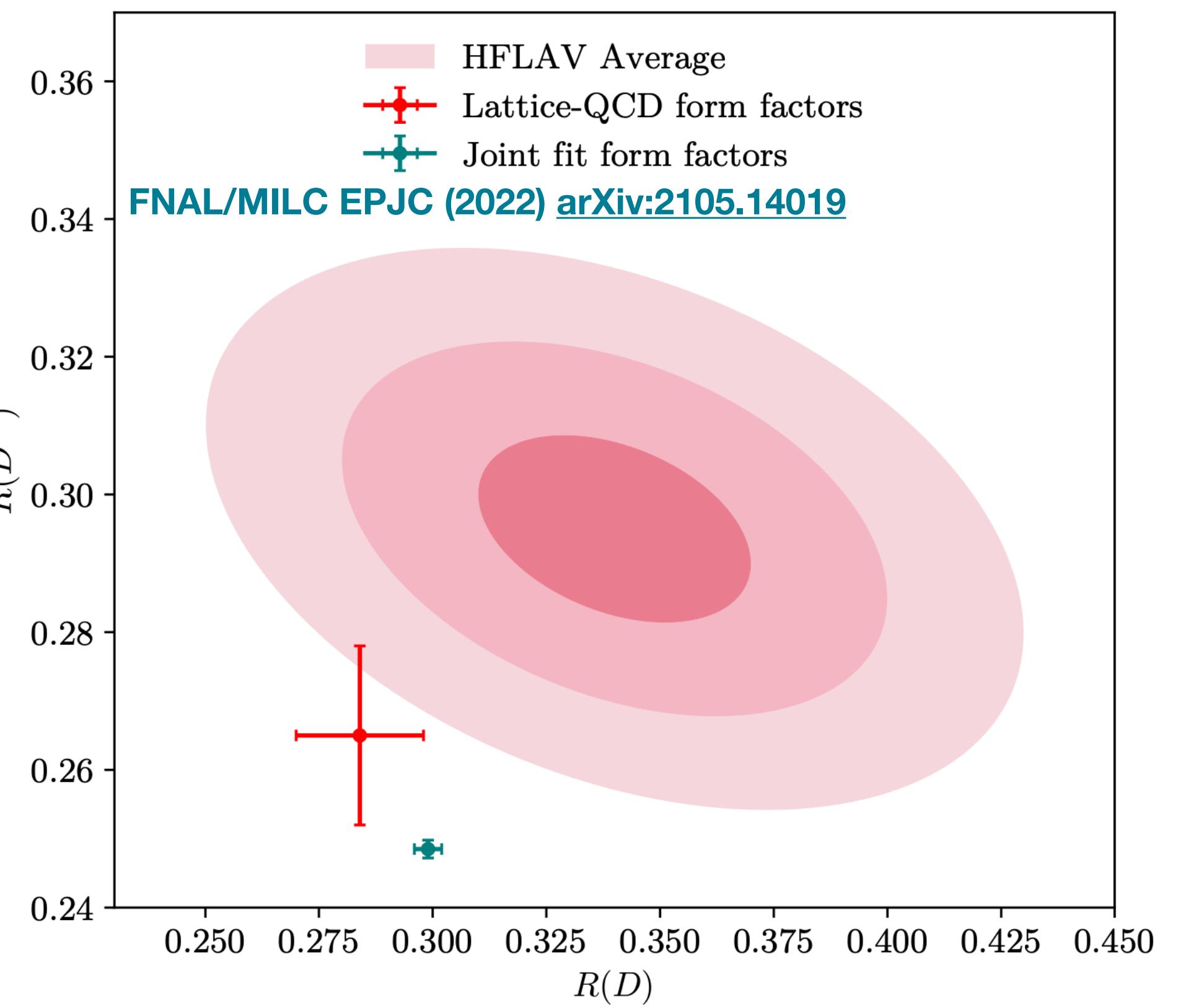


The methods produce essentially the same results. Clear practical advantages of Bayesian inference:

- kinematical constraints exactly and cleanly implemented
- simultaneous fit over various (correlated) data sets possible
- clean statistical underpinning

# $B \rightarrow D$





**FNAL/MILC 0.265(13)  
HPQCD 0.279(13)**